

# Is Our Universe the Remnant of Chiral Anomaly in Axion-Inflation?

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# Cosmic Inflation

A period of exponential expansion of space shortly after the Big Bang



Guth Phys. Rev. D23 (1981) Linde Phys. Lett. B 108 (1982)



# What caused inflation?

A scalar field "slow-rolling" toward its true vacuum provides a simple model for inflation.





## Puzzles of SM & Cosmology

- I) Particle physics of Inflation
- II) Origin of matter asymmetry
- III) Origin of Neutrino mass
- IV) Particle nature of DM

Puzzles of Standard Model of Particle Physics (SM) & Cosmology Which need Physics Beyond SM Inflation

Observab

### Puzzles of SM & Cosmology

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Puzzles of Standard Model of Particle Physics (SM) & Cosmology Which need Physics Beyond SM

Curious cosmological coincidences  $\eta_B \simeq 0.3 P_{\zeta}$  and  $\Omega_{DM} \simeq 5\Omega_B$ !

$$\eta_B = \frac{\mathbf{n}_B - \mathbf{n}_{\bar{B}}}{\mathbf{n}_{\gamma}} \approx 6 \times 10^{-10}$$

Baryon to Photon Ratio Today

$$p_{\zeta} = \frac{1}{2\epsilon} \left( \frac{1}{2\pi} \frac{H}{M_{pl}} \right)^2 \approx 2 \times 10^{-9}$$

Inflation

Curvature Power Spectrum in Inflation

### Puzzles of SM & Cosmology

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- 1. Ad hoc parity violation
- 2. Accidental B-L global symmetry
- 3. Vacuum Stability problem

SM as a particle physics model also faces some conceptual issues

mirror

 $SU(2)_{R}$ 

 $SU(2)_L$ 



• Observations are in perfect agreement with Inflation.

- The Particle Physics of Inflation is still unknown.
- The Standard models of inflation are based on Scalars.

Inflaton: a scalar field beyond the SM. A well-motivated candidate: Axion!





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Inflaton: a scalar field beyond the SM. A well-motivated candidate: Axion!

• Primordial Gravitational Waves (PGW):

Vacuum fluctuations: unpolarized, red-tilted, and nearly Gaussian.







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- They are building blocks of particle physics, SM & beyond.
- They are naturally coupled with axion like particles.
- What do they do in Axion-inflation?!







Collaborators:

E. Komatsu, K. Lozanov, L. Mirzagholi , I. Wolfson, M. Sheikh-Jabbari, J. Soda

#### Colleagues:

P. Adshead, E. Martinec, M. Peloso, E. Dimastrogiovanni, T. Fajita, M. Wyman, E. Sfakianakis, M. Fasiello, R. Caldwell, C. Devulder, Y. Watanabe, ...





#### I) Axion-inflation and gauge fields (non-Abelian) Collaborators:

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II) Embedding Axion-inflation in LR symmetric model A.M. JHEP 2021, 113 (2021)





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 A.M. JHEP 2021, 113 (2021)
 A Common Origin for Inflation,
 Baryogenesis & Dark Matter



























## Gauge fields given by Yang-Mills dilutes like radiation $A_{\mu} \sim 1/a$ $\dot{\rho}_A + 4H\rho_A = 0$

(Axion fields are naturally coupled to gauge fields.)

 $\frac{\lambda}{8f} F \tilde{F} \varphi^{\bullet} \text{Axion}$  $\dot{\rho}_A + 4H \rho_A = \frac{\lambda}{f} \dot{\phi} \text{E. B}$ (Axion generates gauge fields!





### Gauge fields given by Yang-Mills dilutes like radiation $A_{\mu} \sim 1/a$

Spatial isotropy & homogeneity U(1) vacuum  $A_{\mu}$ 

 $A_i = Q(t)\delta_i^3$ 

Gauge fields coupled to inflaton are generated in inflation.

 $\frac{\lambda}{8f} F \tilde{F} \phi^{\prime}$  Axion

(Axion fields are naturally coupled to gauge fields.)

SU(2) vacuum  $A_{\mu} = A^{a}_{\mu} T_{a}$   $[T_{a}, T_{b}] = i \varepsilon^{abc} T_{c}$ Spatially isotropic  $A^{a}_{i} = Q(t)\delta^{a}_{i}$ 

**\.M.** & Sheikh-Jabbari, 2011

so(3) & su(2) are isomorphic

## How SU(2) restores isotropy?

Let us work in temporal gauge,  $A_0 = 0$ .

U(1) vacuum  $A_{\mu}$  $A_i = Q(t)\delta_i^3$ 



Rotation

 $A_i \xrightarrow{\mathrm{SO}(3)} R_{ij} A$ 

A.M. and M. M. Sheikh-Jabbari, 2011

## <u>How SU(2) restores isotropy?</u>

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SU(2) VEV,  $A_{\mu} = A^{a}_{\mu} T_{a}$   $A^{a}_{i} = Q(t)\delta^{a}_{i}$ A.M. and M. M. Sheikh-Jabbari, 2011 Rotation Rotation  $A^{a}_{i} \rightarrow R_{ij}A^{a}_{j}$   $A^{a}_{i} \rightarrow R_{ab}A^{b}_{j}$   $A^{a}_{i} \rightarrow A^{a}_{j}$   $A^{a}_{i} \rightarrow R_{ab}A^{b}_{j}$   $A^{a}_{i} \rightarrow R_{ab}A^{b}_{j}$  $A^{a}_{i} \rightarrow R_{ab}A^{b}_{j}$ 

Rotation

 $A_i \xrightarrow{SO(3)} R_{ij} A_j$ 

- Gauge-flation A. M., & Sheikh-Jabbari, 2011  $S_{Gf} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{4}F^2 + \frac{\kappa}{384}(F\tilde{F})^2 \right)$
- Chromo-natural P. Adshead, M. Wyman, 2012

$$S_{Cn} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{2} \left( (\partial_\mu \varphi)^2 - \mu^4 \left( 1 + \cos(\frac{\varphi}{f}) \right) \right) - \frac{1}{4} F^2 - \frac{\lambda}{8f} \varphi F \tilde{F} \right)$$

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$$Natural inflation$$
Friction

K. Freese, J. A. Frieman and A. V. Olinto 1990

# <u>SU(2)-Axion Model Building</u>

- Gauge-flation A. M., & Sheikh-Jabbari, 2011  $S_{Gf} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{4}F^2 + \frac{\kappa}{384}(F\tilde{F})^2 \right)^4$
- Chromo-natural P. Adshead, M. Wyman, 2012

#### Ruled-out by the data

R. Namba, E. Dimastrogiovanni, M. Peloso 2013 P. Adshead, E. Martinec, M. Wyman 2013

> + Theoretical issue: <u>Very large  $\lambda \sim 100!$ </u>

D. Baumann & L. McAllister 2014

$$S_{Cn} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{2} \left( (\partial_\mu \varphi)^2 - \mu^4 \left( 1 + \cos(\frac{\varphi}{f}) \right) \right) - \frac{1}{4} F^2 - \frac{\lambda}{8f} \varphi F \tilde{F} \right)$$

Inspired by them, several different models with SU(2) fields have been proposed & studied.

### An incomplete list of Different Realizations of the SU(2)-Axion Inflation:

- 1. A. M. and M. M. Sheikh-Jabbari, Phys. Rev. D 84:043515, 2011 [arXiv:1102.1513]
- 2. P. Adshead, M. Wyman, Phys. Rev. Lett.(2012) [*arXiv:1202.2366*]
- 3. **A. M.** JHEP 07 (2016) 104 [arXiv:1604.03327]
- 4. C. M. Nieto and Y. Rodriguez Mod. Phys. Lett. A31 (2016) [arXiv:1602.07197]
- 5. E. Dimastrogiovanni, M. Fasiello, and T. Fujita JCAP 1701 (2017) [arXiv:1608.04216]
- 6. P. Adshead, E. Martinec, E. I. Sfakianakis, and M. Wyman JHEP 12 (2016) 137 [arXiv:1609.04025]

....

- 7. P. Adshead and E. I. Sfakianakis JHEP 08 (2017) 130 [arXiv:1705.03024]
- 8. R. R. Caldwell and C. Devulder Phys. Rev. D97 (2018) [arXiv:1706.03765]
- 9. E. McDonough, S. Alexander, JCAP11 (2018) 030 [arXiv:1806.05684]
- 10. L. Mirzagholi, E. Komatsu, K. D. Lozanov, and Y. Watanabe, [arXiv:2003.04350]
- 11. Y. Watanabe, E. Komatsu, [arXiv:2004.04350]
- 12. J. Holland, I. Zavala, G. Tasinato, [arXiv:2009.00653]

13.

A. M., SU(2)R –axion inflation [arXiv:2012.11516]

Oksana larygina, Evangelos I. Sfakianakis, [arXiv:2105.06972]

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- Gauge-flation A. M., & Sheikh-Jabbari, 2011  $S_{Gf} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{4}F^2 + \frac{\kappa}{384}(F\tilde{F})^2 \right)^{-1}$
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#### SU(2)-Axion inflation has a very rich phenomenology:

- o A new mechanism for generation of Primordial Gravitational Waves
- o All Sakharov conditions are satisfied in inflation: a new baryogenesis mechanism
- o Particle Production in inflation by Schwinger effect and chiral anomaly

P. Adshead et. al 2013 Dimastrogiovanni et. al 2013 **A. M.** et. al, 2013

> **A. M. 2014 & A.M. 2016** R. Caldwell et. al 2017

K. Lozanov, A. M, E. Komatsu 2017,
L. Mirzagholi, A. M, K. Lozanov 2019,
Domcke et al 2019, A.M. 2019

- Gauge-flation A. M., & Sheikh-Jabbari, 2011  $S_{Gf} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{4}F^2 + \frac{\kappa}{384}(F\tilde{F})^2 \right)$
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• Minimal Scenario of SU(2)-axion inflation A.M., 2016 f<0.1 Mpl & λ<0.1

$$S_{AM} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{2} \left( (\partial_\mu \varphi)^2 - V(\varphi) \right) - \frac{1}{4} F^2 - \frac{\lambda}{8f} \varphi F \tilde{F} \right)$$

Axion Monodromy or any mechanism that gives a flat potential



#### <u>New Tensorial mode in SU(2) Gauge Field</u> For $\delta_C > 0$ • $\delta A_i^a = (B_+(t,k)e_{ij}^+(\vec{k}) + B_-(t,k)e_{ij}^-(\vec{k}))\delta_i^a$ Short tachyonic growth of $B_+$ $B_{\pm}^{\prime\prime} + \left[k^2 + \delta_C k \mathcal{H} + \frac{m^2}{H^2} \mathcal{H}^2 - \frac{a^{\prime\prime}}{a}\right] B_{\pm} \approx 0$ $n_B \sim \frac{H^3}{6\pi^2} \delta_C^3 e^{\frac{(2-\sqrt{2})\pi}{2}\delta_C}$ effective frequency Particle Production $(\delta_C \text{ and } \frac{m^2}{\mu^2} \text{ are positive, given by BG})$ A. M. and E. Komatsu, 2018 Vacuum structure $B_+/a$ Axion field $\langle \varphi \rangle$ $\prime$ ( $\delta_C > 0$ ) --- B\_/a Slow-roll A -20 Slow-roll A<sub>P</sub> Parity -40 0.01 0.1 k 1 10 $(\delta_C < 0)$ аН A. M., 2016

# <u>Gauge Field sources Primordial GWs</u>

- $\delta A_i^a = (B_+(t,k)e_{ij}^+(\vec{k}) + B_-(t,k)e_{ij}^-(\vec{k}))\delta_j^a$
- The field equation:  $B_{\pm}^{\prime\prime} + [k^2 + \delta_C k \mathcal{H} + \frac{m^2}{H^2} \mathcal{H}^2 \frac{a^{\prime\prime}}{a}] B_{\pm} \approx 0$



• That sourced the GWs  $h_{\pm}^{\prime\prime} + [k^2 - \frac{a^{\prime\prime}}{a}] h_{\pm} = \mathcal{H}^2 \Pi_{\pm}[B_{\pm}]$ 



Gravitational waves have two uncorrelated terms



$$h_{\pm} = h_{\pm}^{vac} + h_{\pm}^{s}$$

VacuumSourced byGWs $B_{\pm}$ unpolarizedPolarized $h^{vac}_{+} = h^{vac}_{-}$  $h^{s}_{+} \neq h^{s}_{-}$ 





# Novel Observable Signature: CMB

• The sourced tensor modes is Highly non-Gaussian.  $F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu} - ig \left[A_{\mu}, A_{\nu}\right]_{Self-interaction}$ 

Agrawal, Fujita, Komatsu 2018

 That can be probe with future CMB missions., e.g. Litebird and CMB-S4!



#### Equilateral Shape

Maresuke Shiraishi, Front. Astron. Space Sci. 2019

## Novel Observable Signature: Beyond CMB

Detection of this background is an excellent target for all GW experiments across at least 21 decades in frequencies.



P. Campeti, E. Komatsu, D. Poletti, C. Baccigalupi 2020

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P. Campeti, E. Komatsu, D. Poletti, C. Baccigalupi 2020

J. Garcia-Bellido, H. Murayama, and G. White 2021



II) Embedding axion-inflation in Left-Right Symmetric Models

(How to Connect Inflaton to BSM?)



Left-Right Symmetric Model (LRSM)

**A.M.** JHEP 2021, 113 (2021)



• An SU(2) gauge extension of SM with 3 Right-handed Neutrinos coupled to it.



J. C. Pati and A. Salam, Phys. Rev. D 10, 275-289 (1974) R. N. Mohapatra and J. C. Pati, Phys. Rev. D 11, 2558 (1975) G. Senjanovic and R. N. Mohapatra, Phys. Rev. D 12, 1502 (1975)

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- 1. Ad hoc parity violation
- 2. Massive Neutrinos
- 3. Accidental B-L global symmetry
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### Interesting features of LRSM:

# How to Connect it to the SM?

Let us Extend SM Gauge Symmetry by an  $SU(2)_R$  and couple it to Axion Inflaton!

- Minimal Scenario of SU(2)-axion inflation A.M., 2016 f<0.1 Mpl &  $\lambda$ <0.1

$$S_{AM} = \int d^4x \sqrt{-g} \left( -\frac{R}{2} - \frac{1}{4}F^2 - \frac{1}{2} \left( (\partial_\mu \varphi)^2 - V(\varphi) \right) - \frac{\lambda}{8f} \frac{\varphi F \tilde{F}}{4} \right)$$

Axion Monodromy or any mechanism that gives a flat potential

**A. M.** arXiv: 2012.11516 **A.M.** arXiv:2103.14611

### Gauge field Production in Axion-Inflation

• All Gauge fields are diluted by inflation & unimportant , BUT  $SU(2)_R$ :

Axion inflaton arphi

Gauge field (active in inflation)

 $W_R$ 

 $W_R$ 

### Gauge field Production in Axion-Inflation

Axion inflaton  $arphi \otimes$ 

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Gauge field (active in inflation)

 $W_R$ 

 $W_R$ 

•  $W_R^i = B_+(t,k)e_i^+(\vec{k}) + B_-(t,k)e_i^-(\vec{k})$ 

Field Eq.  $B_{\pm}^{\prime\prime} + [k^2 + \xi k\mathcal{H}] B_{\pm} \approx 0$ 

 $(\xi = \frac{2\lambda \partial_t \varphi}{fH})$  effective frequency

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 $W_R$ 

 $W_R$ 

Axion inflaton  $arphi \otimes$ ---

 $n_{WR} \sim \frac{H^3}{6\pi^2} \xi^3 e^{\frac{(2-\sqrt{2})\pi}{2}\xi}$ 

Particle Production



•  $W_R^i = B_+(t,k)e_i^+(\vec{k}) + B_-(t,k)e_i^-(\vec{k})$ Field Eq.  $B_{\pm}^{\prime\prime} + [k^2 + \xi k\mathcal{H}]B_{\pm} \approx 0$ 

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# Summary & Conclusions

### Gauge fields are expected to contribute in physics of axion-inflation.

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Compelling Consequences:

This Set-up is a complete BSM that can solve I-IV:

Particle physics of Inflation
 Origin of matter asymmetry
 Origin of Neutrino mass
 Particle nature of DM

of Particle Cosmology

PUZZLES O



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It provides a deep connection between inflation, baryogenesis & DM, So naturally explains cosmological coincidences  $\eta_B \simeq 0.3 P_{\zeta}$  and  $\Omega_{DM} \simeq 5\Omega_B$ !



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 $m_{\nu}$ 

Inflation

DA

DIOJ

Baryogenest

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It comes with a cosmological smoking gun on Primordial GW.

