Selected topics in hadronic and nuclear physics

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Outline

- ➤ Introduction
- Proton charge radius
- ➢ SIDIS@SoLID, tensor charge, and EDM
- Double polarized photodisintegration of ³He and Gerasimov-Drell-Hearn (GDH) sum rule
- ➢ Summary



Spokespersons: <u>Haiyan Gao</u>, John R. Calarco Students: <u>Benjamin Clasie</u>, Chris Crawford, Jason Seely

And the **BLAST** collaboration.



Nucleon Form Factors The proton charge and current radii are fundamental quantities in physics. Precise determination of the proton charge radius is extremely important to the understanding of the proton structure in terms of quark and gluon degrees of freedom of Quantum Chromodynamics. It is also essential for high-precision tests of Quantum Electrodynamics using the hydrogen Lamb shift measurements. We propose a new precision measurement of the proton charge radius using a laser-driven polarized hydrogen internal gas target and the BLAST detector in the South Hall Ring at the MIT-Bates Laboratory. This measurement will fully utilize the unique features of the polarized internal target, polarized electron beam in the storage ring, and the BLAST large acceptance spectrometer. This measurement is expected to provide the most precise information on the proton charge radius from electron scattering, which will have significant impact on tests of QCD and QED.

The Proton Radius Experiment has been conditionally approved to run at the BATES Linear Accelerator. See the full proposal (also in Postscript).

maintainer: Chris Crawford 617-253-6734 last modified: 2001 April 23

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Proton Charge Radius and the puzzle

- Proton charge radius:
 - 1. An important quantity for proton
 - 2. Important for understanding how QCD works
 - 3. An important physics input to the bound state QED calculation, affects muonic H Lamb shift $(2S_{1/2} 2P_{1/2})$ by as much as 2%, and critical in determining the Rydberg constant
- Methods to measure the proton charge radius:
 - 1. Lepton-proton elastic scattering (nuclear physics)
 - ➢ ep elastic scattering (Mainz-A1, PRad,..)
 - > μp elastic scattering (MUSE, AMBER)
 - 2. Hydrogen spectroscopy (atomic physics)
 - Ordinary hydrogen
 - Muonic hydrogen
- Important point: the proton radius measured in lepton Δ scattering defined the same as in atomic spectroscopy (G.A. Miller, 2019)







$$E = -4\pi\alpha G'^{p}_{E}(0)|\psi_{n0}(0)|^{2}\delta_{l0}$$
$$= 4\pi\alpha \frac{r_{p}^{2}}{6}|\psi_{n0}(0)|^{2}\delta_{l0}.$$

Electron-proton elastic scattering

• **Unpolarized elastic e-p cross section** (*Rosenbluth separation*)

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 \cos^2 \frac{\theta}{2}}{4E^2 \sin^4 \frac{\theta}{2}} \frac{E'}{E} \left(\frac{G_E^{p\ 2} + \tau G_M^{p\ 2}}{1 + \tau} + 2\tau G_M^{p\ 2} \tan^2 \frac{\theta}{2} \right)$$

$$= \sigma_M f_{rec}^{-1} \left(A + B \tan^2 \frac{\theta}{2} \right)$$

$$\tau = \frac{Q^2}{4M^2}$$
One-photon-exchange

• Recoil proton polarization measurement (pol beam only)

$$\frac{G_E^p}{G_M^p} = -\frac{P_t}{P_l} \frac{E+E'}{2M} \tan \frac{\theta}{2}$$

• Asymmetry (super-ratio) measurement (pol beam and pol target)

$$R_{A} = \frac{A_{1}}{A_{2}} = \frac{a_{1} - b_{1} \cdot G_{E}^{p} / G_{M}^{p}}{a_{2} - b_{2} \cdot G_{E}^{p} / G_{M}^{p}}$$



$$A_{exp} = P_b P_t \frac{-2\tau v_{T'} \cos \theta^* G_M^{p-2} + 2\sqrt{2\tau(1+\tau)} v_{TL'} \sin \theta^* \cos \phi^* G_M^p G_E^p}{(1+\tau) v_L G_E^{p-2} + 2\tau v_T G_M^{p-2}}$$

C. Crawford et al, PRL98, 052301 (2007)



The absolute frequency of H energy levels has been measured with an accuracy of 1.4 part in 10^{14} via comparison with an atomic cesium fountain clock as a primary frequency standard.

Yields Rydberg constant R_{∞} (one of the most precisely known constants)

Comparing measurements to QED calculations that include corrections for the finite size of the proton can provide very precise value of the rms proton charge radius **Proton charge radius effect on the muonic hydrogen Lamb shift is 2%**

Muonic hydrogen Lamb shift at PSI (2010, 2013)



Electron-proton Scattering – Mainz A1 experiment

Three spectrometer facility of the A1 collaboration:



- Large amount of overlapping data sets
- Cross section measurement
- Statistical error $\leq 0.2\%$
- Luminosity monitoring with spectrometer
- Q² = 0.004 1.0 (GeV/c)² result: r_p =0.879(5)_{stat}(4)_{sys}(2)_{mod}(4)_{group}
 - J. Bernauer, PRL 105, 242001 (2010)

Measurements @ Mainz



5-7 σ higher than muonic hydrogen result !

(J. Bernauer)

JLab Recoil Proton Polarization Experiment



The situation on the Proton Charge Radius in 2013



This proton charge radius puzzle triggered intensive experimental and theoretical efforts worldwide in the last decade or so



 $R\infty = 10\ 973\ 731.568\ 53(14)\ m^{-1}, r_p = 0.877(13)\ fm$ Fleurbaey *et al.* PRL 120, 183001 (2018)

> Parthey *et al.*, PRL 107, 203001 (2011) Matveev *et al.* PRL 110, 230801 (2013)

The Proton Charge Radius Puzzle in 2018



H spectroscopy (2017): H spectroscopy (2018):

0.8335 ± 0.0095 fm (A. Beyer et al. Science 358(2017) 6359) 0.877 ± 0.013 fm (H. Fleurbaey et al. PRL.120(2018) 183001)

Not shown: ep scattering (ISR, 2017): $0.810 \pm 0.035_{stat.} \pm 0.074_{svst.} \pm 0.003$ (delta_a, delta_b) (Mihovilovic PLB 771 (2017); 0.878 ± 0.011_{stat.} ±0.031_{svst.} ±0.002_{mod}. (Mihovilovic 2021))

The PRad Experiment in Hall B at JLab



- High resolution, large acceptance, hybrid HyCal calorimeter (PbWO₄ and Pb-Glass)
- Windowless H₂ gas flow target
- Simultaneous detection of elastic and Moller electrons
- Q² range of 2x10⁻⁴ 0.06 GeV²
- XY veto counters replaced by GEM detector
- Vacuum chamber

Spokespersons: A. Gasparian (contact), H. Gao, D. Dutta, M. Khandaker

> I. Larin, Y. Zhang *et al.*, Science 6490, 506 (2020)



Elastic ep Cross Sections

- Differential cross section v.s. Q^2 , with 2.2 and 1.1 GeV data
- Statistical uncertainties: $\sim 0.15\%$ for 2.2 GeV, $\sim 0.2\%$ for 1.1 GeV per point
- Systematic uncertainties: 0.3%~1.1% for 2.2 GeV, 0.3%~0.5% for 1.1 GeV (shown as shadow area)



Systematic uncertainties shown as bands

Xiong et al., Nature 575, 147–150 (2019)

Proton Electric Form Factor G'_E (Normalized)

• n_1 and n_2 obtained by fitting PRad G_E to

 $\begin{cases} n_1 f(Q^2), \text{ for 1GeV data} \\ n_2 f(Q^2), \text{ for 2GeV data} \end{cases}$

- G'_E as normalized electric Form factor:
- PRad fit shown as $f(Q^2)$

 $\begin{cases} G_E/n_1, \text{ for 1GeV data} \\ G_E/n_2, \text{ for 2GeV data} \end{cases}$

Using rational (1,1) $f(Q^{2}) = \frac{1 + p_{1}Q^{2}}{1 + p_{2}Q^{2}}$

Yan et al., PRC 98, 025204 (2018)

 r_p = 0.831 +/- 0.007 (stat.) +/- 0.012 (syst.) fm



Proton radius at the time of PRad publication

- PRad result r_p: 0.831 +/- 0.0127 fm, *Xiong et al., Nature 575, 147–150 (2019)*
- H Lamb Shift: 0.833 +/- 0.010 fm, Bezginov et al., Science 365, 1007-1012 (2019)
- CODATA 2018 value of r_p : 0.8414 +/- 0.0019 fm, *E. Tiesinga et al.*, *RMP 93*, 025010(2021)



CODATA has also shifted the value of the Rydberg constant.

More from ordinary hydrogen spectroscopy



Proton radius from ordinary and muonic H spectroscopy



Experiment	Type	Transition(s)	$\sqrt{< r_{Ep}^2 >} \ ({\rm fm})$	$r_{\infty} (\mathrm{m}^{-1})$
Pohl 2010	μH	$2S_{1/2}^{F=1} - 2P_{3/2}^{F=2}$	0.84184(67)	
Antognini 2013	μH	$2S_{1/2}^{F=1} - 2P_{3/2}^{F=2}$	0.84087(39)	
		$2S_{1/2}^{F=0} - 2P_{3/2}^{F=1}$		
Beyer 2017	Н	2S - 4P	0.8335(95)	$10 \ 973 \ 731.568 \ 076 \ (96)$
		with $(1S - 2S)$		
Fleurbaey 2018	H	1S - 3S	0.877(13)	$10\ 973\ 731.568\ 53(14)$
		with $(1S - 2S)$	21	
Bezginov 2019	H	$2S_{1/2} - 2P_{1/2}$	0.833(10)	
Grinin 2020	H	1S - 3S	0.8482(38)	$10\ 973\ 731.568\ 226(38)$
		with $(1S - 2S)$		

(Re)analyses of e-p scattering data



Gao and Vanderhaeghen, arXiv:2105.00571 (to appear in RMP)

e-p scattering: magnetic spectrometer and calorimetric method





PRad-II Experimental Setup (Side View)



Projections for PRad-II



Ongoing and Future Experiments





Orbital motion - Nucleon Structure from 1D to 3D

→ Nucleon Spin
→ Quark Spin





Separation of Collins, Sivers and Pretzelosity through angular dependence

SIDIS SSAs depend on 4-D variables (x, Q², z, P_T) and small asymmetries demand large acceptance + high luminosity allowing for measuring symmetries in 4-D binning with precision!

 $A_{UT}(\phi_h, \phi_S) = \frac{1}{P_{t nol}} \frac{N^{+} - N^{+}}{N^{+} + N^{+}}$

Leading twist formulism (higher-twist terms can be included)

$$=A_{UT}^{Collins}\sin(\phi_h+\phi_S)+A_{UT}^{Pretzelosity}\sin(3\phi_h-\phi_S)+A_{UT}^{Sivers}\sin(\phi_h-\phi_S)$$

$$\begin{bmatrix}Collins\\T\end{bmatrix} \propto \langle \sin(\phi_h + \phi_S) \rangle_{UT} \propto h_1 \otimes H_1^{\perp} \checkmark$$

 $\frac{Pretzelosity}{UT} \propto \langle \sin(3\phi_h - \phi_S) \rangle_{UT} \propto h_{1T}^{\perp} \otimes H_1^{\perp} \checkmark$

Collins fragmentation function from e⁺e⁻ collisions

 $(2\pi \text{ azimuthal coverage})$

$$A_{UT}^{Sivers} \propto \langle \sin(\phi_h - \phi_S) \rangle_{UT} \propto f_{1T}^{\perp} \otimes D_1 \longleftarrow \text{Unpolarized fragmentation} function$$

Haiyan Gao

SoLID@12-GeV JLab: QCD at the intensity frontier

SoLID will maximize the science return of the 12-GeV CEBAF upgrade by combining...

High Luminosity 10³⁷⁻³⁹/cm²/s [>100x CLAS12][>1000x EIC]

Large Acceptance Full azimuthal ϕ coverage

Research at **SoLID** will have the *unique* capability to explore the QCD landscape while complementing the research of other key facilities

- Pushing the phase space in the search of new physics and of hadronic physics
- 3D momentum imaging of a relativistic strongly interacting confined system (<u>nucleon spin</u>)
- Superior sensitivity to the differential electro- and photo-production cross section of J/ψ near threshold (proton mass)

Synergizing with the pillars of EIC science (proton spin and mass) through high-luminosity valence quark tomography and precision J/ψ production near threshold



Haiyan Gao

SIDIS with polarized "neutron" and proton @ SoLID



- E12-10-006:Single Spin Asymmetries on Transversely Polarized ³He @ 90 daysRating ASpokespersons: J.P. Chen, H. Gao (contact), J.C. Peng, X. Qian
- E12-11-007:Single and Double Spin Asymmetries on Longitudinally Polarized ³He @ 35 daysRating ASpokespersons: J.P. Chen (contact), J. Huang, W.B. Yan

E12-11-108:Single Spin Asymmetries on Transversely Polarized Proton @ 120 daysRating ASpokespersons: J.P. Chen, H. Gao (contact), X.M. Li, Z.-E. Meziani

Run group experiments approved for TMDs, GPDs, and spin

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QCD intensity frontier with SoLID: large-acceptance & high luminosity



BigBite

Quantum leap: 4-D binning for the first time!

SoLID-SIDIS program: Large acceptance, Full azimuthal coverage + High luminosity

- 4-D mapping of asymmetries with precision $\Delta z = 0.05$, $\Delta P_T = 0.2$ GeV, $\Delta Q^2 = 1$ GeV², x bin sizes vary with median bin size 0.02 (statistical uncertainty for each bin: $\delta A \le 0.02$)
- Constrain models and forms of TMDs, Tensor charge, ...
- Lattice QCD, QCD dynamics, models



• More than 1400 bins in x, Q², P_T and z for 11/8.8 GeV beam.

e $3 He \approx 90\%$ -90% -1.5% -8% *X. Qian et al., PRL107, 072003(2011)*Haiyan Gao

 π

Polarized

³He Target

HRS,

Transversity and Tensor Charge

Transversity distribution

 h_1 (Collinear & TMD)

- Chiral-odd, unique for the quarks
- No mixing with gluons, simpler evolution effect
- A transverse counterpart to longitudinal spin g₁, difference shows the relativistic effect
- Zeroth moment gives tensor charge:

$$\left\langle \mathbf{P}, \mathbf{S} | \overline{\psi}_q i \sigma^{\mu\nu} \psi_q | \mathbf{P}, \mathbf{S} \right\rangle = g_T^q \overline{u}(\mathbf{P}, \mathbf{S}) i \sigma^{\mu\nu} u(\mathbf{P}, \mathbf{S})$$
$$g_T^q = \int_0^1 \left[h_1^q(x) - h_1^{\overline{q}}(x) \right] dx$$

- A fundamental QCD quantity, valence quarks dominate
- Precisely calculated on the lattice
- Difference from nucleon axial charge is due to relativity
- High luminosity-large acceptance allows for high-precision test of LQCD predictions

FLAG review 2019: 1902.08191 Relative uncertainty 4% (u), 7% (d)



JAM20: arxiv:2002.08384





Constraint on quark EDMs with combined proton and neutron EDMs

	d _u upper limit	d _d upper limit
Current g _T + current EDMs	1.27×10 ⁻²⁴ <i>e</i> cm	1.17×10 ⁻²⁴ <i>e</i> cm
SoLID g_T + current EDMs	6.72×10 ⁻²⁵ <i>e</i> cm	1.07×10 ⁻²⁴ <i>e</i> cm
SoLID g _T + future EDMs	1.20×10 ⁻²⁷ <i>e</i> cm	7.18×10 ⁻²⁸ <i>e</i> cm

Include 10% isospin symmetry breaking uncertainty Sensitivity to new physics

$$d_q \sim e m_q / (4 \pi \Lambda^2)$$

Three orders of magnitude improvement on quark EDM limit

Current quark EDM limit: $10^{-24}e$ cm

Future quark EDM limit: $10^{-27}e$ cm

Image credit: D. Pitonyak

Probe to $30 \sim 40$ times higher scale





 $30 \sim 40 \text{ TeV}$

H. Gao, T. Liu, Z. Zhao, PRD 97, 074018 (2018)

TENSOR CHARGE AT THE EIC AND JLAB



JAM20: Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato, Phys.Rev.D 102 (2020)



L. Gamberg, Z. Kang, D. Pitonyak, A. Prokudin, N. Sato Phys.Lett.B 816 (2021)

EIC data will allow to have g_T extraction at the precision at the level of lattice QCD calculations

 JLab 12 data will allow to have complementary information on tensor charge to test the consistency of the extraction and expand the kinematical region

A. Prokudin@PANIC 2021

Thank you for your attention!

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