# Selected topics in hadronic and nuclear physics 

# October 10, 2021 Symposium on QCD and Nuclei 

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## Outline

$>$ Introduction
$>$ Proton charge radius
$>$ SIDIS@SoLID, tensor charge, and EDM
$>$ Double polarized photodisintegration of ${ }^{3} \mathrm{He}$ and Gerasimov-Drell-Hearn (GDH) sum rule
$>$ Summary



## Collaboration

Spokespersons: Haiyan Gao, John R. Calarco Students: Benjamin Clasie, Chris Crawford, Jason Seely And the BLAST collaboration.

## Projects

Reports
Software
Target
Nucleon Form Factors
The proton charge and current radii are fundamental quantities in physics. Precise determination of the proton charge radius is extremely important to the understanding of the proton structure in terms of quark and gluon degrees of freedom of Quantum Chromodynamics. It is also essential for high-precision tests of Quantum Electrodynamics using the hydrogen Lamb shift measurements. We propose a new precision measurement of the proton charge radius using a laser-driven polarized hydrogen internal gas target and the BLAST detector in the South Hall Ring at the MIT-Bates Laboratory. This measurement will fully utilize the unique features of the polarized internal target, polarized electron beam in the storage ring, and the BLAST large acceptance spectrometer. This measurement is expected to provide the most precise information on the proton charge radius from electron scattering, which will have significant impact on tests of QCD and QED.

The Proton Radius Experiment has been conditionally approved to run at the BATES Linear Accelerator. See the full proposal (also in Postscript).
maintainer: Chris Crawford 617-253-6734 last modified: 2001 April 23


## Proton Charge Radius and the puzzle

- Proton charge radius:

1. An important quantity for proton
2. Important for understanding how QCD works
3. An important physics input to the bound state QED calculation, affects muonic H Lamb shift $\left(2 \mathrm{~S}_{1 / 2}-2 \mathrm{P}_{1 / 2}\right)$ by as much as $2 \%$, and critical in determining the Rydberg constant


RP, Gilman, Miller, Pachucki, Annu. Rev. Nucl. Part. Sci. 63, 175 (2013)

2. Hydrogen spectroscopy (atomic physics)
> Ordinary hydrogen
$>$ Muonic hydrogen

$$
\sqrt{<r^{2}>}=\sqrt{-\left.6 \frac{d G\left(q^{2}\right)}{d q^{2}}\right|_{q^{2}=0}}
$$

$>$ Important point: the proton radius measured in lepton scattering defined the same as in atomic spectroscopy (G.A. Miller, 2019)

## Electron-proton elastic scattering

- Unpolarized elastic e-p cross section (Rosenbluth separation)

$$
\begin{aligned}
& \frac{d \sigma}{d \Omega}=\frac{\alpha^{2} \cos ^{2} \frac{\theta}{2}}{4 E^{2} \sin ^{4} \frac{\theta}{2}} \frac{E^{\prime}}{E}\left(\frac{G_{E}^{p 2}+\tau G_{M}^{p}{ }^{2}}{1+\tau}+2 \tau G_{M}^{p}{ }^{2} \tan ^{2} \frac{\theta}{2}\right) \\
&=\sigma_{M} f_{\text {rec }}^{-1}\left(A+B \tan ^{2} \frac{\theta}{2}\right) \\
& 4 M^{2}
\end{aligned}
$$



- Recoil proton polarization measurement (pol beam only)

$$
\frac{G_{E}^{p}}{G_{M}^{p}}=-\frac{P_{t}}{P_{l}} \frac{E+E^{\prime}}{2 M} \tan \frac{\theta}{2}
$$

- Asymmetry (super-ratio) measurement (pol beam and pol target)

$$
R_{A}=\frac{A_{1}}{A_{2}}=\frac{a_{1}-b_{1} \cdot G_{E}^{p} / G_{M}^{p}}{a_{2}-b_{2} \cdot G_{E}^{p} / G_{M}^{p}}
$$


$A_{\text {exp }}=P_{b} P_{t} \frac{-2 \tau v_{T^{\prime}} \cos \theta^{*} G_{M}^{p}{ }^{2}+2 \sqrt{2 \tau(1+\tau)} v_{T L^{\prime}} \sin \theta^{*} \cos \phi^{*} G_{M}^{p} G_{E}^{p}}{(1+\tau) v_{L} G_{E}^{p}{ }^{2}+2 \tau v_{T} G_{M}^{p}{ }^{2}}$
C. Crawford et al, PRL98, 052301 (2007)


The absolute frequency of $H$ energy levels has been measured with an accuracy of 1.4 part in $10^{14}$ via comparison with an atomic cesium fountain clock as a primary frequency standard.
Yields Rydberg constant $\mathrm{R}_{\infty}$ (one of the most precisely known constants)

Comparing measurements to QED calculations that include corrections for the finite size of the proton can provide very precise value of the rms proton charge radius Proton charge radius effect on the muonic hydrogen Lamb shift is $\mathbf{2 \%}$

## Muonic hydrogen Lamb shift at PSI $(2010,2013)$



## Electron-proton Scattering - Mainz A1 experiment

Three spectrometer facility of the A1 collaboration:


- Large amount of overlapping data sets
- Cross section measurement
- Statistical error $\leq 0.2 \%$
- Luminosity monitoring with spectrometer
- $Q^{2}=0.004-1.0(\mathrm{GeV} / \mathrm{c})^{2}$
result: $r_{p}=0.879(5)_{\text {stat }}(4)_{\text {sys }}(2)_{\text {mod }}(4)_{\text {group }}$
J. Bernauer, PRL 105, 242001 (2010)

Measurements @ Mainz


5-7 $\sigma$ higher than muonic hydrogen result!

## JLab Recoil Proton Polarization Experiment



The situation on the Proton Charge Radius in 2013


This proton charge radius puzzle triggered intensive experimental and theoretical efforts worldwide in the last decade or so

## Ordinary hydrogen spectroscopy


$\mathrm{R} \infty=10973731.568076(96) \mathrm{m}^{-1}, \mathrm{r}_{\mathrm{p}}=0.8335(95) \mathrm{fm}$
Beyer et al., Science 358, 79 (2017)

$\mathrm{R} \infty=10973731.568$ 53(14) $\mathrm{m}^{-1}, \mathrm{r}_{\mathrm{p}}=0.877(13) \mathrm{fm}$
Fleurbaey et al. PRL 120, 183001 (2018)

## The Proton Charge Radius Puzzle in 2018



Electron scattering: $\quad 0.879 \pm 0.011 \mathrm{fm}$ (CODATA 2014)
Muon spectroscopy: $\quad 0.8409 \pm 0.0004 \mathrm{fm}$ (CREMA 2010, 2013)
H spectroscopy (2017): $0.8335 \pm 0.0095 \mathrm{fm}(\mathrm{A}$. Beyer et al. Science 358(2017) 6359)
H spectroscopy (2018): $\quad 0.877 \pm 0.013 \mathrm{fm}$ (H. Fleurbaey et al. PRL.120(2018) 183001)

Not shown: ep scattering (ISR, 2017): $0.810 \pm 0.035_{\text {stat. }} \pm 0.074_{\text {syst. }} \pm 0.003$ (delta_a, delta_b) (Mihovilovic PLB 771 (2017); $0.878 \pm 0.011_{\text {stat. }} \pm 0.031_{\text {syst. }} \pm 0.002_{\text {mod }}$. (Mihovilovic 2021))

## The PRad Experiment in Hall B at JLab



- High resolution, large acceptance, hybrid HyCal calorimeter ( $\mathrm{PbWO}_{4}$ and Pb -Glass)
- Windowless $\mathbf{H}_{2}$ gas flow target
- Simultaneous detection of elastic and Moller electrons
- $Q^{2}$ range of $2 \times 10^{-4}-0.06 \mathrm{GeV}^{2}$
- XY - veto counters replaced by GEM detector
- Vacuum chamber

Spokespersons: A. Gasparian (contact), H. Gao, D. Dutta, M. Khandaker
I. Larin, Y. Zhang et al., Science 6490, 506 (2020)

## Elastic ep Cross Sections

- Differential cross section v.s. $\mathrm{Q}^{2}$, with 2.2 and 1.1 GeV data
- Statistical uncertainties: $\sim 0.15 \%$ for $2.2 \mathrm{GeV}, \sim 0.2 \%$ for 1.1 GeV per point
- Systematic uncertainties: $0.3 \% \sim 1.1 \%$ for $2.2 \mathrm{GeV}, 0.3 \% \sim 0.5 \%$ for 1.1 GeV (shown as shadow area)


Systematic uncertainties shown as bands
Xiong et al., Nature 575, 147-150 (2019)

## Proton Electric Form Factor G $_{E}$ (Normalized)

- $n_{1}$ and $n_{2}$ obtained by fitting PRad $\mathrm{G}_{\mathrm{E}}$ to

$$
\left\{\begin{array}{l}
n_{1} f\left(Q^{2}\right), \text { for } 1 \mathrm{GeV} \text { data } \\
n_{2} f\left(Q^{2}\right), \text { for } 2 \mathrm{GeV} \text { data }
\end{array}\right.
$$

Using rational $(1,1)$
$f\left(Q^{2}\right)=\frac{1+p_{1} Q^{2}}{1+p_{2} Q^{2}}$

- $\mathrm{G}_{\mathrm{E}}^{\prime}$ as normalized electric Form factor: $\left\{\begin{array}{l}G_{E} / n_{1}, \text { for } 1 \mathrm{GeV} \text { data } \\ G_{E} / n_{2}, \text { for } 2 \mathrm{GeV} \text { data }\end{array}\right.$

$$
r_{p}=0.831+/-0.007 \text { (stat.) }+/-0.012 \text { (syst.) fm }
$$



Xiong et al., Nature 575, 147-150 (2019)


$$
n_{1}=1.0002+/-0.0002 \text { (stat.) +/- } 0.0020 \text { (syst.), }
$$

$$
n_{2}=0.9983+/-0.0002 \text { (stat.) }+/-0.0013 \text { (syst.) }
$$

## Proton radius at the time of PRad publication

- PRad result $\mathrm{r}_{\mathrm{p}}: 0.831+/-0.0127 \mathrm{fm}$, Xiong et al., Nature 575, 147-150 (2019)
- H Lamb Shift: 0.833 +/- 0.010 fm, Bezginov et al., Science 365, 1007-1012 (2019)
- CODATA 2018 value of $\mathrm{r}_{\mathrm{p}}: 0.8414+/-0.0019 \mathrm{fm}$, E. Tiesinga et al., RMP 93, 025010(2021)


CODATA has also shifted the value of the Rydberg constant.

More from ordinary hydrogen spectroscopy



Bezginov et al., Science 365, 1007 (2019) $\mathrm{r}_{\mathrm{p}}=0.833(10) \mathrm{fm}$


Grinin et al., Science 370, 1061 (2020)

$$
r_{p}=0.8482(38) \mathrm{fm}
$$

## Proton radius from ordinary and muonic H spectroscopy



| Experiment | Type | Transition(s) | $\sqrt{\left\langle r_{E p}^{2}\right\rangle}(\mathrm{fm})$ | $r_{\infty}\left(\mathrm{m}^{-1}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
| Pohl 2010 | $\mu \mathrm{H}$ | $2 S_{1 / 2}^{F=1}-2 P_{3 / 2}^{F=2}$ | $0.84184(67)$ |  |
| Antognini 2013 | $\mu \mathrm{H}$ | $2 S_{1 / 2}^{F=1}-2 P_{3 / 2}^{F=2}$ | $0.84087(39)$ |  |
| Beyer 2017 | H | $2 S_{1 / 2}^{F=0}-2 P_{3 / 2}^{F=1}$ | $2 S-4 P$ | $0.8335(95)$ |
| with $(1 S-2 S)$ |  | $10973731.568076(96)$ |  |  |
| Fleurbaey 2018 | H | $1 S-3 S$ | $0.877(13)$ | $10973731.56853(14)$ |
| Bezginov 2019 | H | with $(1 S-2 S)$ <br> $2 S_{1 / 2}-2 P_{1 / 2}$ <br> Grinin 2020 | H | $1 S-3 S$ |
|  |  | $0.833(10)$ |  |  |
| with $(1 S-2 S)$ | $0.8482(38)$ | $10973731.568226(38)$ |  |  |

## (Re)analyses of e-p scattering data



Gao and Vanderhaeghen, arXiv:2105.00571 (to appear in RMP)

## e-p scattering: magnetic spectrometer and calorimetric method



- Upgrade HyCal
- Adding $2^{\text {nd }}$ GEM

PRad-II Experimental Setup (Side View)


## Projections for PRad-II



## Ongoing and Future Experiments



| Experiment | Beam | Laboratory | $Q^{2}(\mathrm{GeV} / \mathrm{c})^{2}$ | $\delta r_{p}(\mathrm{fm})$ | Status |
| :---: | :---: | :---: | :---: | :---: | :---: |
| MUSE | $e^{ \pm}, \mu^{ \pm}$ | PSI | $0.0015-0.08$ | 0.01 | Ongoing |
| AMBER | $\mu^{ \pm}$ | CERN | $0.001-0.04$ | 0.01 | Future |
| PRad-II | $e^{-}$ | Jefferson Lab | $4 \times 10^{-5}-6 \times 10^{-2}$ | 0.0036 | Future |
| PRES | $e^{-}$ | Mainz | $0.001-0.04$ | $0.6 \%$ (rel.) | Future |
| A1@MAMI (jet target) | $e^{-}$ | Mainz | $0.004-0.085$ |  | Ongoing |
| MAGIX@MESA | $e^{-}$ | Mainz | $\geq 10^{-4}-0.085$ |  | Future |
| ULQ $^{2}$ | $e^{-}$ | Tohoku University | $3 \times 10^{-4}-8 \times 10^{-3}$ | $\sim 1 \%$ (rel.) | Future |



## Orbital motion - Nucleon Structure from 1D to 3D <br> $\rightarrow$ Nucleon Spin <br> $\rightarrow$ Quark Spin



Corresponding fragmentation functions

Generalized parton distribution (GPD)

Transverse momentum dependent parton distributions (TMD-PDFs)

Connected via 5-D Wigner distribution

Transverse momentum dependent Fragmentation functions (TMD-FFs)



Separation of Collins, Sivers and Pretzelosity through angular dependence

SIDIS SSAs depend on 4-D variables ( $x, Q^{2}, z, P_{T}$ ) and small asymmetries demand large acceptance + high luminosity allowing for measuring symmetries in 4-D binning with precision!

$$
A_{U T}\left(\phi_{h}, \phi_{S}\right)=\frac{1}{P_{t, p o l}} \frac{N^{\uparrow}-N^{\downarrow}}{N^{\uparrow}+N^{\downarrow}}
$$


$A_{U T}^{\text {Sivers }}$
$\propto\left\langle\sin \left(\phi_{h}-\phi_{S}\right)\right\rangle_{U T} \propto f_{1 T}^{\perp} \otimes D_{1}$ $\qquad$ Unpolarized fragmentation function

Haiyan Gao

## SoLID@12-GeV JLab: QCD at the intensity frontier

SoLID will maximize the science return of the $12-\mathrm{GeV}$ CEBAF upgrade by combining...
High Luminosity
$10^{37-39} / \mathrm{cm}^{2} / \mathrm{s}$
$[>100 \times$ CLAS12 $][>1000 \times$ EIC $]$

## Large Acceptance

 Full azimuthal $\phi$ coverageResearch at SoLID will have the unique capability to explore the QCD landscape while complementing the research of other key facilities

- Pushing the phase space in the search of new physics and of hadronic physics
- 3D momentum imaging of a relativistic strongly interacting confined system (nucleon spin)
- Superior sensitivity to the differential electro- and photo-production cross section of $\mathrm{J} / \psi$ near threshold (proton mass)

Synergizing with the pillars of EIC science (proton spin and mass) through high-luminosity valence quark tomography and precision $\mathrm{J} / \psi$ production near threshold


Haiyan Gao

## SIDIS with polarized "neutron" and proton @ SoLID

SoLID (SIDIS He3)


SoLID (SIDIS NH3)


E12-10-006: $\quad$ Single Spin Asymmetries on Transversely Polarized ${ }^{3} \mathrm{He} @ 90$ days Rating A Spokespersons: J.P. Chen, H. Gao (contact), J.C. Peng, X. Qian

E12-11-007: Single and Double Spin Asymmetries on Longitudinally Polarized ${ }^{3} \mathrm{He} @ 35$ days Rating A

E12-11-108: Rating A

Spokespersons: J.P. Chen (contact), J. Huang, W.B. Yan

Single Spin Asymmetries on Transversely Polarized Proton @ 120 days Spokespersons: J.P. Chen, H. Gao (contact), X.M. Li, Z.-E. Meziani

Run group experiments approved for TMDs, GPDs, and spin

## QCD intensity frontier with SoLID: large-acceptance \& high luminosity


X. Qian et al., PRL107, 072003(2011)

Quantum leap: 4-D binning for the first time!
SoLID-SIDIS program: Large acceptance, Full azimuthal coverage + High luminosity

- 4-D mapping of asymmetries with precision $\Delta z=0.05, \Delta P_{T}=0.2 \mathrm{GeV}, \Delta \mathrm{Q}^{2}=1 \mathrm{GeV}^{2}$, x bin sizes vary with median bin size 0.02 (statistical uncertainty for each bin: $\delta A \leq 0.02$ )
- Constrain models and forms of TMDs, Tensor charge, ...
- Lattice QCD, QCD dynamics, models

- More than 1400 bins in $x, Q^{2}, P_{T}$ and $z$ for $11 / 8.8 \mathrm{GeV}$ beam.


## Transversity and Tensor Charge

## Transversity distribution

$h_{1}$
 (Collinear \& TMD)

- Chiral-odd, unique for the quarks
- No mixing with gluons, simpler evolution effect
- A transverse counterpart to longitudinal spin $\mathrm{g}_{1}$, difference shows the relativistic effect
- Zeroth moment gives tensor charge:

$$
\begin{gathered}
\langle\mathrm{P}, \mathrm{~S}| \bar{\psi}_{q} i \sigma^{\mu \nu} \psi_{q}|\mathrm{P}, \mathrm{~S}\rangle=g_{T}^{q} \bar{u}(\mathrm{P}, \mathrm{~S}) i \sigma^{\mu \nu} u(\mathrm{P}, \mathrm{~S}) \\
\quad g_{T}^{q}=\int_{0}^{1}\left[h_{1}^{q}(x)-h_{1}^{\bar{q}}(x)\right] d x
\end{gathered}
$$

- A fundamental QCD quantity, valence quarks dominate
- Precisely calculated on the lattice
- Difference from nucleon axial charge is due to relativity
- High luminosity-large acceptance allows for high-precision test of LQCD predictions

FLAG review 2019: 1902.08191
Relative uncertainty 4\% (u), 7\% (d)



JAM20: arxiv:2002.08384

## Constraint on Quark EDMs <br> $d_{n}=g_{T}^{d} d_{u}+g_{T}^{u} d_{d}+g_{T}^{s} d_{s}$

Constraint on quark EDMs with combined proton and neutron EDMs

|  | $\mathbf{d}_{\mathbf{u}}$ upper limit | $\mathbf{d}_{\mathbf{d}}$ upper limit |
| :---: | :---: | :--- |
| Current $\mathrm{g}_{\boldsymbol{T}}+$ current EDMs | $1.27 \times 10^{-24} e \mathrm{~cm}$ | $1.17 \times 10^{-24} e \mathrm{~cm}$ |
| SoLID $\mathrm{g}_{\mathrm{T}}+$ current EDMs | $6.72 \times 10^{-25} e \mathrm{~cm}$ | $1.07 \times 10^{-24} e \mathrm{~cm}$ |
| SoLID $\mathrm{g}_{\mathrm{T}}+$ future EDMs | $1.20 \times 10^{-27} e \mathrm{~cm}$ | $7.18 \times 10^{-28} e \mathrm{~cm}$ |

Include 10\% isospin symmetry breaking uncertainty
Sensitivity to new physics

$$
d_{q} \sim e m_{q} /\left(4 \pi \Lambda^{2}\right)
$$

Three orders of magnitude improvement on quark EDM limit

Probe to $30 \sim 40$ times higher scale

Current quark EDM limit: $10^{-24} \mathrm{ecm}$
Future quark EDM limit: $10^{-27} e \mathrm{~cm}$ Image credit: D. Pitonyak

$$
\sim 1 \mathrm{TeV}
$$

$$
30 \sim 40 \mathrm{TeV}
$$

## TENSOR CHARGE AT THE EIC AND JLAB

L. Gamberg, Z. Kang, D. Pitonyak, A. Prokudin, N. Sato Phys.Lett.B 816 (2021)


JAM20: Cammarota, Gamberg, Kang, Miller,
Pitonyak, Prokudin, Rogers, Sato, Phys.Rev.D 102 (2020)


- JLab 12 data will allow to have complementary information on tensor charge to test the consistency of the extraction and expand the kinematical region
A. Prokudin@PANIC 2021


## Thank you for your attention!

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