

Review

Leading Quark TMDPDFs  Nucleon Spin  Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \text{Unpolarized}$		$h_1^\perp = \text{Boer-Mulders}$
	L		$g_1 = \text{Helicity}$	$h_{1L}^\perp = \text{Worm-gear}$
	T	$f_{1T}^\perp = \text{Sivers}$	$g_{1T}^\perp = \text{Worm-gear}$	$h_1 = \text{Transversity}$ $h_{1T}^\perp = \text{Pretzelosity}$

sign flip

$$\begin{aligned} (f_{1T}^\perp)^{SIDIS} &= - (f_{1T}^\perp)^{DY}, & \text{rest equal} \\ (h_1^\perp)^{SIDIS} &= - (h_1^\perp)^{DY}, & \text{in SIDIS \& DY} \end{aligned}$$

Leading Quark TMDFFs  Hadron Spin  Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Unpolarized (or Spin 0) Hadrons		$D_1 = \text{Unpolarized}$		$H_1^\perp = \text{Collins}$

Before we turn to ingredients for phenomenology -15-
 there is one more interesting question

TMD Integral

Handbook § 2.9

$$\int d^2 k_T f_{i/p}(x, k_T, \mu, y) \stackrel{?}{=} f_i(x, \mu)$$

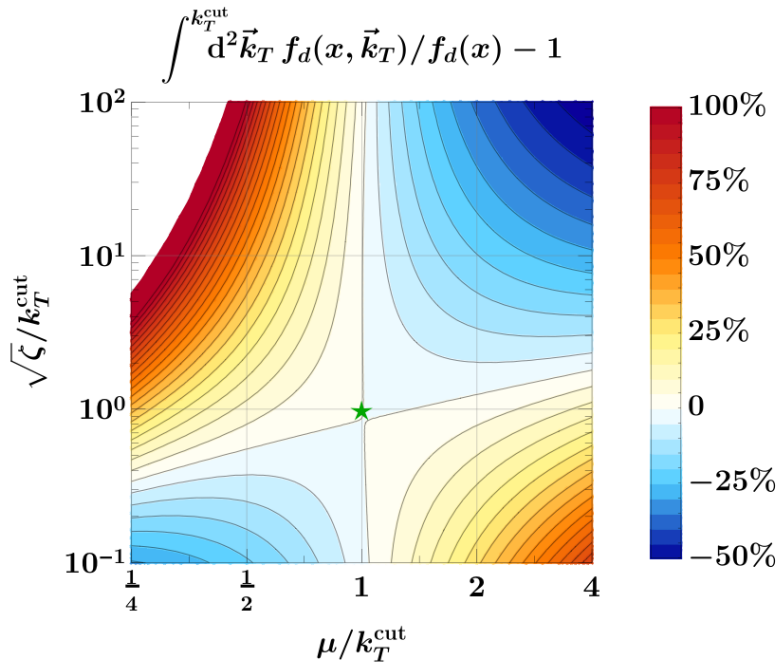
Naïve : Yes

UV \neq RGE : No , $\mu \neq y$ dependence differs
 ($k_T \rightarrow \infty$) (μ, y) on LHS \neq RHS

In fact

$$\int_0^{k_T^{cut}} d^2 k_T f_{i/p}(x, \vec{k}_T, \mu = k_T^{cut}, y = (k_T^{cut})^2) \simeq f_i(x, \mu = k_T^{cut})$$

at percent level :



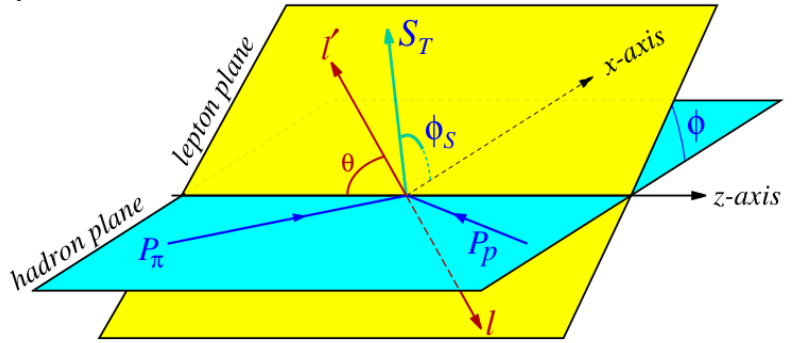
See Zhiguan Sun's talk tonight.

Polarized Drell-Yan $\pi(P_\pi) + p(P_p, S) \rightarrow \gamma^*/Z \rightarrow \ell^+ \ell^- X$

Angles θ, ϕ, ϕ_s defined in Collins-Soper frame

- $(\ell^+ \ell^-)$ at rest
- $P_{\pi T} = P_{pT} = \frac{q_T}{2}$

Study Leading terms for $q_T \ll Q$



Structure Functions

$F = F(x_\pi, x_p, q_T, Q^2)$

$$\frac{d\sigma}{d^4q d\Omega} = \frac{\alpha_{em}^2}{\mathcal{F} Q^2} \left\{ \left[(1 + \cos^2 \theta) F_{UU}^1 + \sin^2 \theta \cos(2\phi) F_{UU}^{\cos 2\phi} \right] \right.$$

$$+ S_L \sin^2 \theta \sin(2\phi) F_{UL}^{\sin 2\phi}$$

$$+ S_T (1 - \cos^2 \theta) \sin \phi_s F_{UT}^{\sin \phi_s}$$

$$\left. + S_T \sin^2 \theta \left[\sin(2\phi + \phi_s) F_{UT}^{\sin(2\phi + \phi_s)} + \sin(2\phi - \phi_s) F_{UT}^{\sin(2\phi - \phi_s)} \right] \right\}$$

$\uparrow \uparrow$ polarization of proton
 π unpol.

Factorization

- $F_{UU}^1 = \mathcal{B}[f_{1,\pi}^{(0)} f_{1,p}^{(0)}], \leftarrow$ Unpol.
- $F_{UU}^{\cos 2\phi} = M_\pi M_p \mathcal{B}[\tilde{h}_{1,\pi}^{\perp(1)} \tilde{h}_{1,p}^{\perp(1)}], \leftarrow$ Boer-Mulders²
- $F_{UL}^{\sin 2\phi} = -M_\pi M_p \mathcal{B}[\tilde{h}_{1,\pi}^{\perp(1)} \tilde{h}_{1L,p}^{\perp(1)}], \leftarrow$ " " & Worm-Gear
- $F_{UT}^{\sin \phi_s} = M_p \mathcal{B}[f_{1,\pi}^{(0)} \tilde{f}_{1T,p}^{\perp(1)}], \leftarrow$ Unpol & Sivers
- $F_{UT}^{\sin(2\phi - \phi_s)} = -M_\pi \mathcal{B}[\tilde{h}_{1,\pi}^{\perp(1)} \tilde{h}_{1,p}^{(0)}], \leftarrow$ Boer-Mulders & Transversity
- $F_{UT}^{\sin(2\phi + \phi_s)} = -\frac{M_\pi M_p^2}{4} \mathcal{B}[\tilde{h}_{1,\pi}^{\perp(1)} \tilde{h}_{1T,p}^{\perp(2)}]. \leftarrow$ Boer Mulders & Pretzelosity

$$\mathcal{B}[f_\pi^{(m)} f_p^{(n)}] \equiv \sum_i H_{i\bar{i}}(Q, \mu) \int_0^\infty \frac{db_T}{2\pi} b_T b_T^{m+n} J_{m+n}(q_T b_T) \tilde{f}_{i/p}^{(m)}(x_a, b_T, \mu, \zeta_a) \tilde{f}_{\bar{i}/\pi}^{(n)}(x_b, b_T, \mu, \zeta_b)$$

Polarized SIDIS

$$l(\lambda, \lambda) + p(\epsilon, \epsilon) \rightarrow l(\lambda') + h(P_h) + X$$

↑ massless
↑ helicity
↑ unpol.

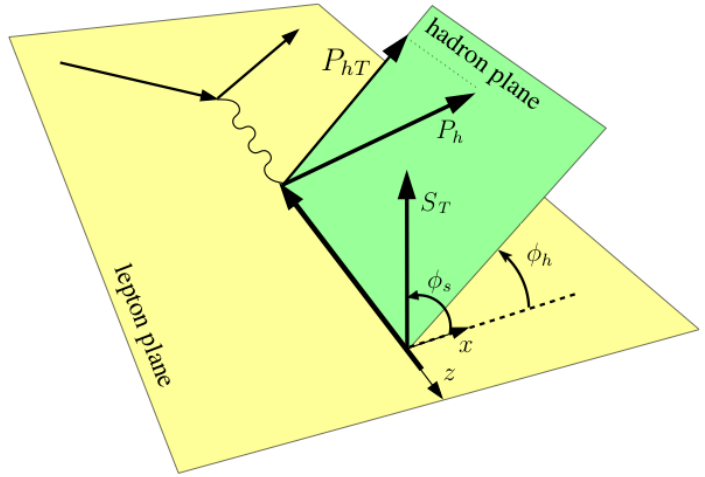
Frame for angles ϕ_h, ϕ_s

= Trento convention

- $\vec{q} \parallel \hat{z}$
- Leptons: x-z plane

Consider $Q \ll M_W, z$

Leading Order in $Q_T \ll Q$



Structure Functions

$$F = F(x, z_h, P_{hT}, Q^2)$$

$$\frac{d^6\sigma}{dx dy dz_h d\phi_s d\phi_h dP_{hT}^2} = \frac{\alpha_{em}^2}{x y Q^2} \left(1 - y + \frac{1}{2} y^2 \right) \left[F_{UU,T} + \cos(2\phi_h) p_1 F_{UU}^{\cos(2\phi_h)} + S_L \sin(2\phi_h) p_1 F_{UL}^{\sin(2\phi_h)} + S_L \lambda p_2 F_{LL} + S_T \sin(\phi_h - \phi_s) F_{UT,T}^{\sin(\phi_h - \phi_s)} + S_T \sin(\phi_h + \phi_s) p_1 F_{UT}^{\sin(\phi_h + \phi_s)} + \lambda S_T \cos(\phi_h - \phi_s) p_2 F_{LT}^{\cos(\phi_h - \phi_s)} + S_T \sin(3\phi_h - \phi_s) p_1 F_{UT}^{\sin(3\phi_h - \phi_s)} \right]$$

$$p_i = p_i(y)$$

beam pol. $l(\lambda)$ target pol. $p(\epsilon)$

Factorization

- $F_{UU}(x, z_h, P_{hT}, Q^2) = \mathcal{B} \left[\tilde{f}_1^{(0)} \tilde{D}_1^{(0)} \right],$ ← unpol.
- $F_{UU}^{\cos 2\phi_h}(x, z_h, P_{hT}, Q^2) = M_N M_h \mathcal{B} \left[\tilde{h}_1^{\perp(1)} \tilde{H}_1^{\perp(1)} \right],$ ← Boer-Mulders * Collins
- $F_{UL}^{\sin 2\phi_h}(x, z_h, P_{hT}, Q^2) = M_N M_h \mathcal{B} \left[\tilde{h}_{1L}^{\perp(1)} \tilde{H}_1^{\perp(1)} \right],$ ← Worm-Gear * Collins
- $F_{LL}(x, z_h, P_{hT}, Q^2) = \mathcal{B} \left[\tilde{g}_1^{(0)} \tilde{D}_1^{(0)} \right],$ ← Helicity * Unpol.
- $F_{LT}^{\cos(\phi_h - \phi_s)}(x, z_h, P_{hT}, Q^2) = M_N \mathcal{B} \left[\tilde{g}_{1T}^{\perp(1)} \tilde{D}_1^{(0)} \right],$ ← Worm-Gear * Unpol
- $F_{UT}^{\sin(\phi_h + \phi_s)}(x, z_h, P_{hT}, Q^2) = M_h \mathcal{B} \left[\tilde{h}_1^{(0)} \tilde{H}_1^{\perp(1)} \right],$ ← Transversity * Collins
- $F_{UT}^{\sin(\phi_h - \phi_s)}(x, z_h, P_{hT}, Q^2) = -M_N \mathcal{B} \left[\tilde{f}_{1T}^{\perp(1)} \tilde{D}_1^{(0)} \right],$ ← Sivers * Unpol
- $F_{UT}^{\sin(3\phi_h - \phi_s)}(x, z_h, P_{hT}, Q^2) = \frac{M_N^2 M_h}{4} \mathcal{B} \left[\tilde{h}_{1T}^{\perp(2)} \tilde{H}_1^{\perp(1)} \right],$ ← Pretzelosity * Collins

$$\mathcal{B}[\tilde{f}^{(m)} \tilde{D}^{(n)}] \equiv x \sum_i e_i^2 \mathcal{H}_{ii}(Q^2, \mu) \int_0^\infty \frac{db_T}{2\pi} b_T b_T^{m+n} J_{m+n}(q_T b_T) \tilde{f}_{i/N}^{(m)}(x, b_T, \mu, \zeta_1) \tilde{D}_{h/i}^{(n)}(z_h, b_T, \mu, \zeta_2)$$

Other Processes : $e^+ e^- \rightarrow h_1 + h_2 + X$ see § 2.11
 \rightarrow 2 TMD FFS

Implementation

Ⓐ Resummation $\alpha_s \ln\left(\frac{Q}{b_T}\right) \sim \alpha_s \ln(Q b_T) \sim 1$

$$\begin{aligned} \tilde{\sigma}^W(b_T) = f_q(x_1) f_{\bar{q}}(x_2) C[\alpha_s] \exp \left\{ \begin{aligned} & \frac{\alpha_s}{4\pi} \left(d_{12} L_b^2 + d_{11} L_b \right) \\ & + \left(\frac{\alpha_s}{4\pi} \right)^2 \left(d_{23} L_b^3 + d_{22} L_b^2 + d_{21} L_b \right) \\ & + \left(\frac{\alpha_s}{4\pi} \right)^3 \left(d_{34} L_b^4 + d_{33} L_b^3 + d_{32} L_b^2 + d_{31} L_b \right) \end{aligned} \right\} + \dots, \end{aligned}$$

LL

NLL

NNLL

N³LL

Derived by Solving Renormalization Group Equations (RGE)

- $\mu \frac{d}{d\mu} \ln H(Q^2, \mu) = \gamma_H^H(Q, \mu) = 4 \Gamma_{\text{cusp}}[\alpha_s] \ln \frac{Q}{\mu} + \gamma_H^H[\alpha_s]$
- $\mu \frac{d}{d\mu} \ln \tilde{f}_{i/p} = \gamma_{\mu^2}^i(\mu, y) = -\Gamma_{\text{cusp}} \ln y / \mu^2 + \gamma_{\mu^2}^i[\alpha_s]$

Note: always perturbative for $\mu \gg \Lambda_{\text{QCD}}$.

- $y \frac{d}{dy} \ln \tilde{f}_{i/p} = \frac{1}{2} \gamma_y^i(\mu, b_T) \left(= \frac{1}{2} \tilde{K}(b_T; \mu) \right)$

↑
rapidity anom. dim

↑
Collins-Soper Kernel

$$\gamma_y^g(\mu, b_T) = -2 \int_{\gamma_{b_T}}^{\mu} \frac{d\mu'}{\mu'} \Gamma_{\text{cusp}}[\alpha_s(\mu')] + \gamma_y^g[\alpha_s(\gamma_{b_T})]$$

Note: $\gamma_y^g(\mu, b_T)$ is non-perturbative for $b_T^{-1} \sim \Lambda_{\text{QCD}}$!

(b) Combine Pert. & Non-Pert

$$b_T^{-1} \gg \Lambda_{\text{QCD}} \quad b_T^{-1} \sim \Lambda_{\text{QCD}}$$

$$f_{i/h}(x, b_T, \mu, y) = f_{i/h}^{\text{pert}}(x, b^*(b_T), \mu, y) f^{\text{NP}^*}(x, b_T)$$

fit to TMD data

$$\sum_j \int \frac{dz}{z} C_{ij}\left(\frac{x}{z}, b_T, \mu, y\right) f_j(z, \mu)$$

pert.
Global fits

- $b^*(b_T)$ shields pert. from Landau Pole $\alpha_s(\Lambda_{\text{QCD}}) = \infty$
- f^{NP^*} 's meaning depends on choice of b^*

(c) Y term (eg. DY)

$$\frac{d\sigma}{dQ dY d^2q_T} = \frac{d\sigma^W}{dQ dY d^2q_T} + \frac{d\sigma^Y}{dQ dY d^2q_T}$$

\uparrow Factorized, dominant $q_T \ll Q$
 \uparrow $O(q_T^2/Q^2) + \dots$
important when $q_T \sim Q$

(d) Pert. Orders LO, NLO (α_s), NNLO (α_s^2), ...
in H_{ij} , C_{ij} & Y-term

SV19 = Scimemi, Vladimirov (1912.06532)

Pavia19 = Bachetta, Bertone, Bissolotti, Bozzi, Delcarro,
Piacenza, Radici (1912.07550)

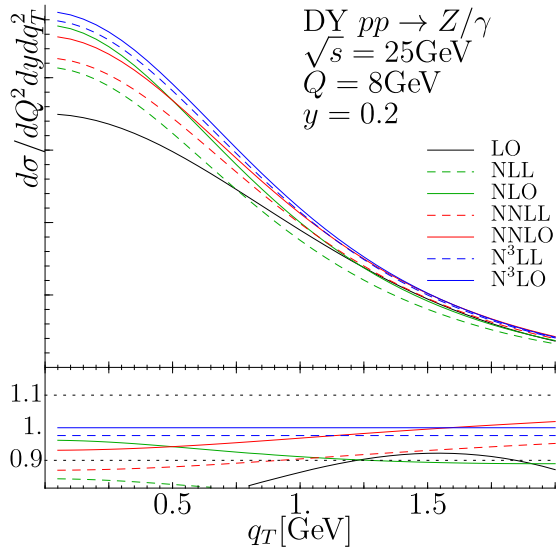
Common features:

- Unpolarized Data with constraint: $q_T/Q < 0.2 - 0.25$ (4-6% power corrections)
- Perturbative accuracy: N3LL resummation + NNLO matching to PDF
- Neglect small contributions from Boer-Mulders terms (higher twist for pert. b_T)
- Common b_T dependence for all flavors

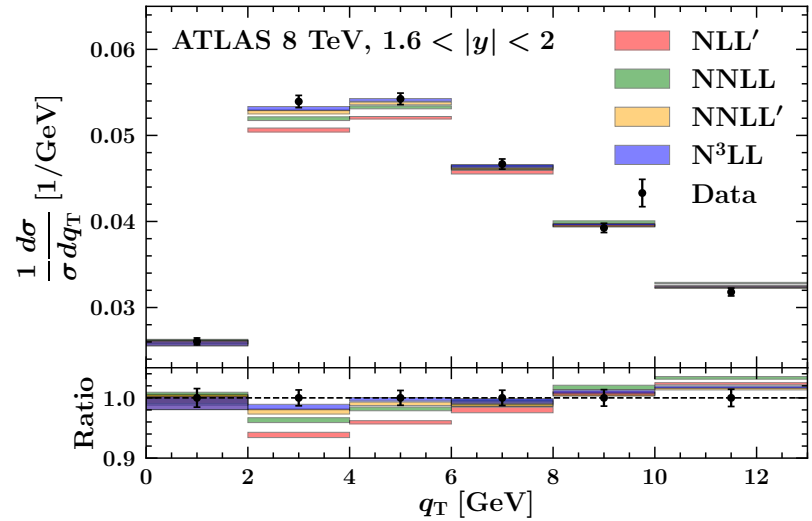
Common features:

Good Perturbative convergence:

SV19



Pavia19



Global Fits

Differences:

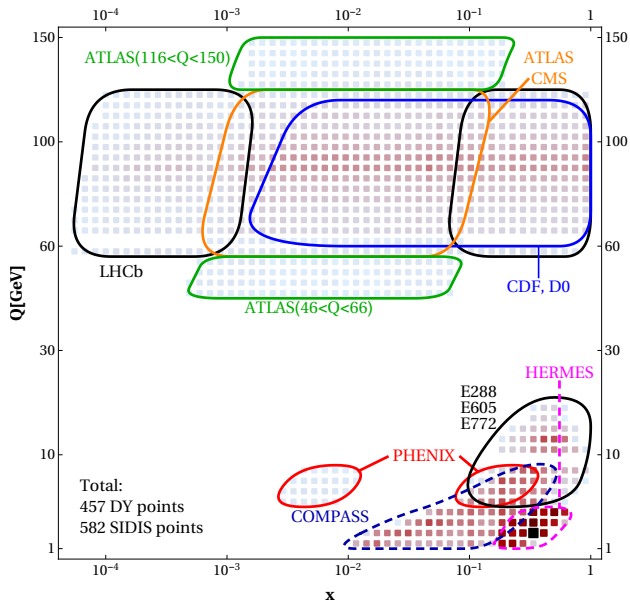
Some differences in solution of evolution equations

Datasets used

SV19

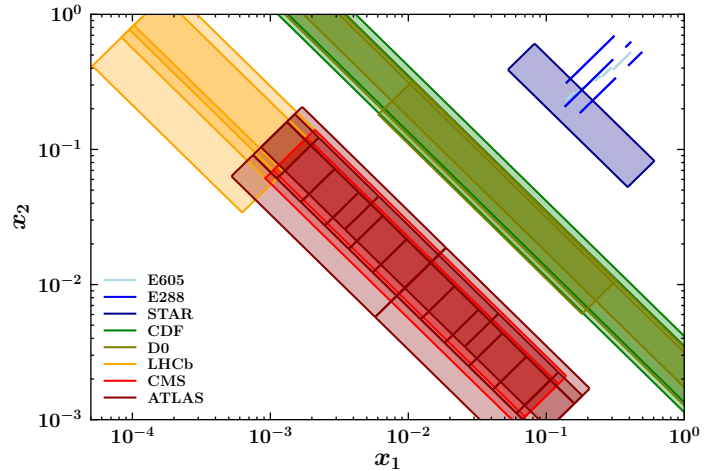
Drell-Yan (457 bins)

SIDIS (582 bins)



Pavia19

Drell-Yan (353 bins)



$$x_1 = Qe^y/\sqrt{s}, \quad x_2 = Qe^{-y}/\sqrt{s}$$

Global Fits

Differences:

Non-perturbative Models

SV19

TMDPDF: 5
TMDFF: 4
CS kernel: 2

$$f_{NP}(x, b) = \exp\left(-\frac{\lambda_1(1-x) + \lambda_2 x + x(1-x)\lambda_5}{\sqrt{1 + \lambda_3 x \lambda_4} b^2}\right)$$

$$D_{NP}(x, b) = \exp\left(-\frac{\eta_1 z + \eta_2(1-z) b^2}{\sqrt{1 + \eta_3 (b/z)^2} z^2}\right) \left(1 + \eta_4 \frac{b^2}{z^2}\right)$$

$$\gamma_\zeta^q(\mu, b) = \gamma_\zeta^{q \text{ pert}}(\mu, b^*) - \frac{1}{2} c_0 b b^*$$

$$b^*(b) = \frac{b}{\sqrt{1 + b^2/B_{NP}^2}}$$

Pavia19

TMDPDF: 7
CS kernel: 2

$$f_{NP}(x, b_T) = \left[\frac{1 - \lambda}{1 + g_1(x) \frac{b_T^2}{4}} + \lambda \exp\left(-g_{1B}(x) \frac{b_T^2}{4}\right) \right]$$

$$g_1(x) = \frac{N_1}{x\sigma} \exp\left[-\frac{1}{2\sigma^2} \ln^2\left(\frac{x}{\alpha}\right)\right]$$

$$g_{1B}(x) = \frac{N_{1B}}{x\sigma_B} \exp\left[-\frac{1}{2\sigma_B^2} \ln^2\left(\frac{x}{\alpha_B}\right)\right]$$

$$\gamma_\zeta^q(\mu, b) = \gamma_\zeta^{q \text{ pert}}(\mu, b_*) - \frac{1}{2} (g_2 b_T^2 + g_{2B} b_T^4)$$

$$b_*(b_T) = b_{\max} \left(\frac{1 - \exp\left(-\frac{b_T^4}{b_{\max}^4}\right)}{1 - \exp\left(-\frac{b_T^4}{b_{\min}^4}\right)} \right)^{\frac{1}{4}}$$

Note: model form for b^* used to split perturbative & non-perturbative parts

Fit Results:

SV19

$$\chi^2/N_{pt} = 1.06$$

NP-parameters	
RAD	$B_{NP} = 1.93 \pm 0.22$ $c_0 = (4.27 \pm 1.05) \times 10^{-2}$
TMDPDF	$\lambda_1 = 0.224 \pm 0.029$ $\lambda_2 = 9.24 \pm 0.46$ $\lambda_3 = 375. \pm 89.$ $\lambda_4 = 2.15 \pm 0.19$ $\lambda_5 = -4.97 \pm 1.37$
TMDFF	$\eta_1 = 0.233 \pm 0.018$ $\eta_2 = 0.479 \pm 0.025$ $\eta_3 = 0.472 \pm 0.041$ $\eta_4 = 0.511 \pm 0.040$

Low and High energy data are well described

RAD parameters are less sensitive to input PDF set

Universality of RAD satisfied by DY vs. SIDIS data

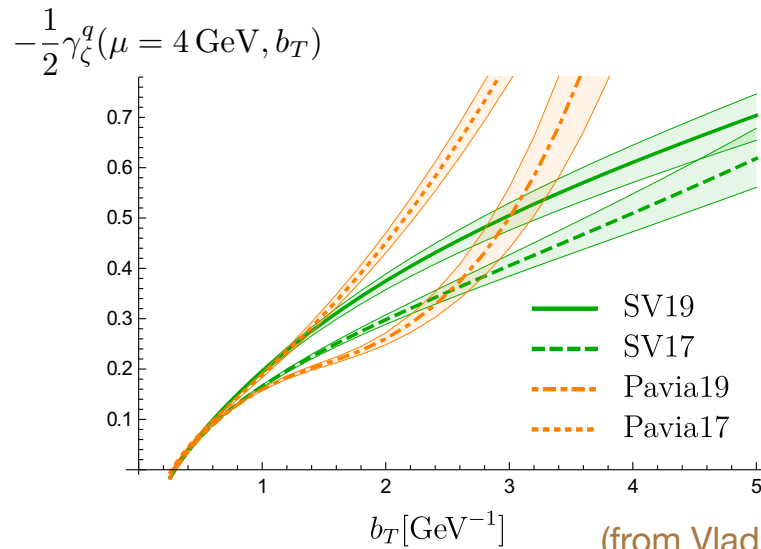
Pavia19

$$\chi^2/N_{pt} = 1.02$$

Parameter	Value
g_2	0.036 ± 0.009
N_1	0.625 ± 0.282
α	0.205 ± 0.010
σ	0.370 ± 0.063
λ	0.580 ± 0.092
N_{1B}	0.044 ± 0.012
α_B	0.069 ± 0.009
σ_B	0.356 ± 0.075
g_{2B}	0.012 ± 0.003

Fit Results:

Comparison of results for CS Kernel in non-perturbative regime:

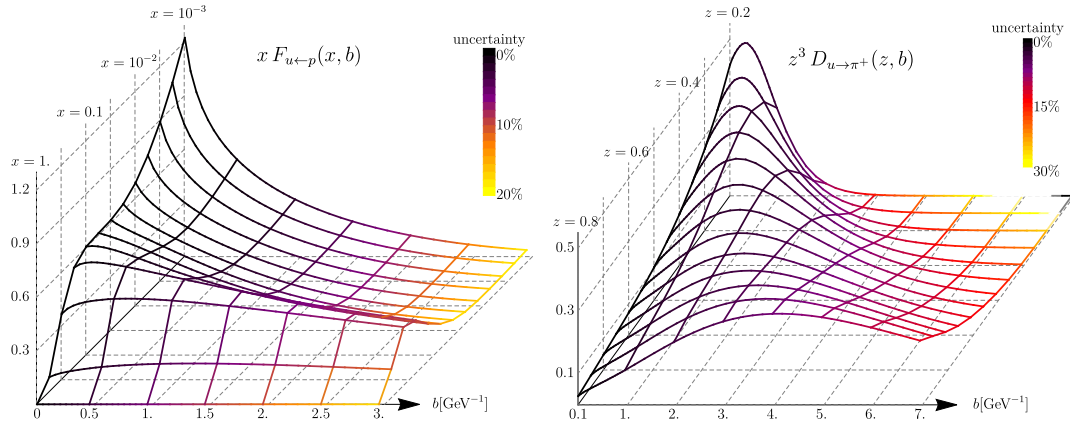


(from Vladimirov, 2003.02288)

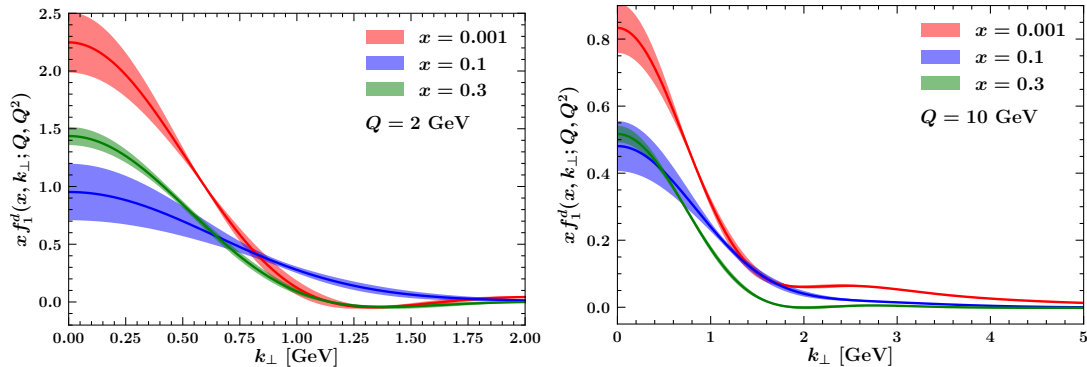
Fit Results:

Results for intrinsic TMDPDF (& TMDFF)

SV19



Pavia19



Quite precise determinations if we assume a given fit form.

Global Fits

Bury, Prokudin, Vladimirov (2012.05135)

Extraction of **Sivers function** from global fit to SIDIS, DY, and W/Z data
[76 bins: HERMES, COMPASS, Jlab (SIDIS); STAR(W/Z); COMPASS (DY)]

$$f_{1T}^{\perp \text{SIDIS}} = -f_{1T}^{\perp \text{DY}}$$

N3LL analysis following SV19

Flavor dependent parametrization (no matching)

$$f_{1T;q \leftarrow h}^{\perp}(x, b) = N_q \frac{(1-x)x^{\beta_q}(1+\epsilon_q x)}{n(\beta_q, \epsilon_q)} \exp\left(-\frac{r_0 + x r_1}{\sqrt{1+r_2 x^2} b^2} b^2\right)$$

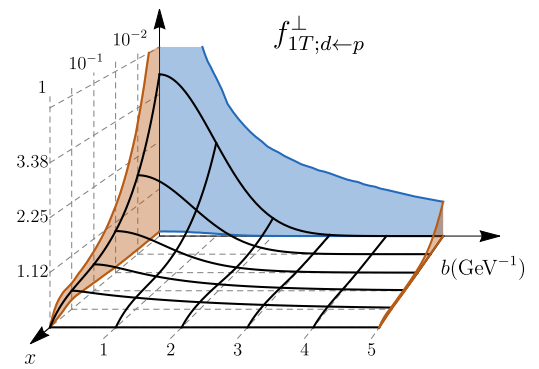
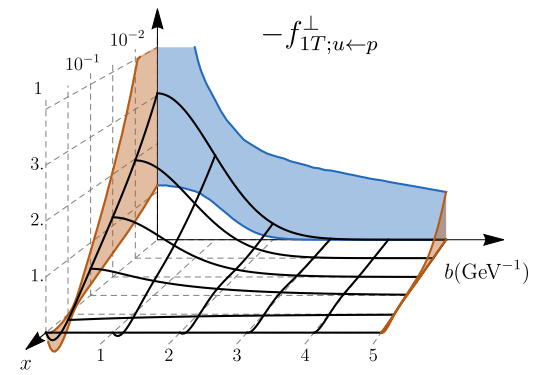
Results:

Good global fit: $\chi^2/N_{pt} = 0.88$

Opposite signs for up and down Sivers functions

Data not precise enough to confirm sign flip

$$f_{1T}^{\perp \text{SIDIS}} = +f_{1T}^{\perp \text{DY}} \quad \text{gives} \quad \chi^2/N_{pt} = 1.0$$



The End

