

Based on work in collaboration with W. Buchmüller, V. Domcke, and H. Murayama [1912.03695, 2009.10649, 2107.04578]

Kai Schmitz

Junior professor at the University of Münster, Germany Cambridge High Energy Workshop 2022, Harvard University Phase Transitions and Topological Defects in the Early Universe Center of Mathematical Sciences and Applications | August 3, 2022



Cosmic strings:

• Topological defects after *U*(1) breaking in the early Universe

[Ringeval: 1005.4842]



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- Network of long strings and closed loops in scaling regime

[Allen, Martins, Shellard: ctc.cam.ac.uk/outreach]



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[CERN]

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Gravitational waves (GWs):

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Gravitational waves (GWs):

• Loop oscillations + GW bursts from cusps and kinks on loops

Assumption: Energy loss via particle emission off closed loops is negligible [Matsunami, Pogosian, Saurabh, Vachaspati: 1903.05102] [Hindmarsh, Lizarraga, Urio, Urrestilla: 2103.16248]

Stable cosmic strings and NANOGrav

[Blasi, Brdar, KS: 2009.06607] [See also Ellis, Lewicki: 2009.06555]



 \odot Explain NANOGrav signal for $G\mu \sim 10^{-(10 \cdots 11)}$ and $lpha \sim 0.1$

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- \odot GUT scale $\Lambda \sim 10^{15 \cdots 16} \text{ GeV}$ points to $G\mu \sim 10^{-(7 \cdots 8)}$ (smaller α ?)
- Signal at higher frequencies too small for LIGO, Virgo, KAGRA

Cosmic strings and grand unification

[Dror, Hiramatsu, Kohri, Murayama, White: 1908.03227] [See also King, Pascoli, Turner, Zhou: 2005.13549, 2106.15634]



UV embedding of the seesaw mechanism in GUT models: Neutrino mass, leptogenesis, cosmic strings, GWs, proton decay

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Assumption: Inflation dilutes monopoles; otherwise string-monopole gas



Decay rate per string length:

[Vilenkin: Nucl. Phys. B 196 (1982) 240] [Preskill, Vilenkin: hep-ph/9209210] [Monin, Voloshin: 0808.1693]

$$\Gamma_d = \frac{d\#}{dtd\ell} = \frac{\mu}{2\pi} e^{-\pi\kappa}, \qquad \kappa = \frac{m^2}{\mu} \qquad (1)$$

String tension µ, monopole mass m



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Monopoles with and without unconfined magnetic flux:

- Unconfined flux: $M\bar{M}$ annihilation, emission of massless gauge bosons
- No unconfined flux: energy loss only via emission of gravitational waves

Possible scenarios

$$W_{B-L} = \lambda T \left(S\bar{S} - \frac{1}{2} v_{B-L}^2 \right) + \frac{h_i}{M_*} S^2 N_i^2$$
(2)

B-L phase transition after supersymmetric hybrid inflation:

• T: inflaton, S, \overline{S} : Higgs / waterfall fields, N_i : right-handed neutrinos

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 $\underset{\text{[Buchmüller: 2102.08923]}}{\text{Minimal alternative: }} SU(2) \times U(1) \stackrel{\text{triplet}}{\longrightarrow} U(1) \times U(1) \stackrel{\text{doublets}}{\longrightarrow} U(1)$

Strategy

End of scaling when long string segments begin to enter the horizon: [Leblond, Shlaer, Siemens: 0903.4686]

$$\Gamma_d \,\ell \, t_s \sim \Gamma_d H^{-1} \, t_s \sim \Gamma_d t_s^2 \sim 1 \quad \Rightarrow \quad t_s \sim \frac{1}{\sqrt{\Gamma_d}} \tag{3}$$



Scaling regime, $t < t_s$

- Loops: emit GWs, decay into segments negligible
- Long strings: decay into segments on superhorizon scales, chop off closed loops, GW emission negligible

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Decay regime, $t > t_s$

- Loops: emit GWs and decay into segments
- Segments from loops and long strings: emit GWs and decay into segments; no production of new loops

Kinetic equation for the number densities of loops and segments, $\stackrel{\circ}{n}$ and $\tilde{n}{:}$

$$\partial_t n(\ell, t) = S(\ell, t) - \partial_\ell \left[u(\ell, t) n(\ell, t) \right] - \left[3H(t) + \Gamma_d \ell \right] n(\ell, t) \quad (4)$$

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Source term *S*:

• Loops from long strings (loop production function): $S \propto t^{-4} \, \delta \left(\ell - lpha t
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Time derivative of the string length $u = \dot{\ell}$:

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Challenge: Solve set of partial integro-differential equations in both the scaling and decay regimes, match solutions at $t = t_s$. (Plus, RD / MD.)

Loop number density during the decay regime in the radiation era:

[Cf. Blanco-Pillado, Olum, Shlaer: 1309.6637] [Cf. Blanco-Pillado, Olum: 1709.02693]

$$\stackrel{\circ}{n_{>}}^{\rm rr}(\ell,t) = \frac{B \, e^{-\Gamma_{d} \left[\ell(t-t_{s})+\frac{1}{2}\Gamma G \mu(t-t_{s})^{2}\right]}}{t^{3/2} \left(\ell+\Gamma G \mu t\right)^{5/2}} \,\Theta\left(\alpha t_{s}-\bar{\ell}\left(t_{s}\right)\right) \,\Theta\left(t_{\rm eq}-t\right)$$
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• Exponential suppression at $\ell t > 1/\Gamma_d = t_s^2$ or $t^2 > 2/(\Gamma_d \Gamma G \mu) = t_e^2$ because of new exponential suppression factor:

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Similar results for $\overset{\circ}{n_{<}}^{rr}$, $\overset{\circ}{n_{<}}^{rm}$, $\overset{\circ}{n_{<}}^{rm}$, $\overset{\circ}{n_{>}}^{rm}$, $\overset{\circ}{n_{>}}^{rm}$, $\overset{\circ}{n_{>}}^{rm}$, $\overset{\circ}{n_{<}}^{rm}$, $\overset{\circ}{n_{<}}^{rr}$, $\overset{\circ}{n_{<}}^{(s) rr}$, $\overset{\circ}{n_{>}}^{(s) rr}$, $\overset{\circ}{n_{>}}^{(s)$

Compute GW spectrum following the standard procedure:

$$\Omega_{\rm gw}(f) = \frac{G\mu^2}{\rho_{\rm crit}} \sum_k P_k \frac{2k}{f} \int_{t_{\rm ini}}^{t_0} dt \left[\frac{a(t)}{a(t_0)}\right]^5 n\left(\frac{a(t)}{a(t_0)}\frac{2k}{f}, t\right)$$
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- Suppress spectrum in nHz range, explain NANOGrav for larger $G\mu$



Extrapolate spectrum to large *f* and compare with LIGO, Virgo, KAGRA:



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- Tilt at PTA frequencies correlated with amplitude at LVK frequencies
- LISA will probe the entire parameter space consistent with NANOGrav



Metastable cosmic strings:

• Prediction in many GUT models when combined with inflation to solve the monopole problem



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Thank you very much for your attention!