Coupling a Cosmic String to a TQFT

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(Based on works in progress with T. Daniel Brennan and Liantao Wang)

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Symmetry plays the central role in theoretical physics!

- **Symmetry** plays the central role in theoretical physics!
- Recently, notion/concept of Symmetry has gone through explosive generalizations!
 - Ordinary Symmetry → Higher-form Symmetry
 0-form 0-form 1-form, 2-form, ...
 particle or local operator extended object: line, surface, ...
 - → Higher-Group, Non-Invertible Symmetries
 - Conserved in entire → Subsystem Symmetries
 Spacetime (Fracton)
 - More...

Ο

Group

(Q1) Are there generalized symmetries in (3+1)d QFTs that relevant for particle physics?

(Q2) Can there be observable signals (even in principle) associated with (due to) the presence of those generalized symmetries?

(Q3) Can generalized symmetry provide novel or meaningful solutions to problems in particle physics?

Today's Talk

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(Q3) Can generalized symmetry provide novel or meaningful solutions to problems in particle physics?

<u>Outline</u>

- I. Axion-Maxwell Theory
- **II. TQFT-Coupling 1:** Axion-Portal to a Z_n TQFT IR-Universal Observables from TQFT-Coupling Anomaly-Inflow, Fermion Zero Modes
- **III. TQFT-Coupling 2:** Gauging Discrete Subgroup of Axion Shift Modification of Cosmic String Spectrum
- **IV. Conclusion**

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IV. Conclusion

$$S = \int \frac{1}{2} \left(\partial_{\mu} a \right)^{2} + \int \frac{1}{2g^{2}} F_{\mu\nu} F^{\mu\nu} - \int \frac{iK}{16\pi^{2}} \frac{a}{f} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

cf)
$$F_{\mu\nu}\tilde{F}^{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}F_{\rho\sigma}$$

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- This very familiar theory enjoys a large set of GGS:
 0-form axion shift
 - $\,\circ\,$ 2-form axion winding
 - \circ 1-form electric
 - 1-form magnetic
 - ✤ 3-group
 - * Non-invertible symmetries (Cordova, Ohmori '22)

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$$\theta = \frac{a}{f} \rightarrow \theta + c$$

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 - \circ 2-form axion winding

•
$$Q(\Sigma_1) = \oint_{\Sigma_1} \frac{d\theta}{2\pi} = m$$
 : Winding number

Charged object: cosmic string (Mst₂)



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$$+\int \frac{in}{2\pi} B_2 \wedge dB_1 - \int \frac{iK_{AB}}{4\pi^2} \frac{a}{f} F_A \wedge F_B - \int \frac{iK_B}{8\pi^2} \frac{a}{f} F_B \wedge F_B$$



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- (ii) Can be obtained from Abelian Higgs Model (AHM)

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$$\mathcal{L} \sim \Lambda^2(nB_1) \wedge * (d\varphi) \sim \frac{in}{2\pi} B_1 \wedge dB_2 \qquad (dB_2 \sim * d\varphi)$$

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- (i) Describes discrete (\mathbb{Z}_n) gauge theory
- (ii) Can be obtained from Abelian Higgs Model (AHM)
- (iii) Describes non-trivial vacuum structure and selection rules in terms of topological degrees of freedom B_1 , B_2

$$dB_1 = 0, \quad dB_2 = 0$$

$$\langle W_1(\Sigma_1, m) W_2(\Sigma_2, \ell) \rangle \sim e^{i \frac{2\pi}{n} m \ell \operatorname{Link}(\Sigma_1, \Sigma_2)}$$

$$W_1(\Sigma_1, m) = e^{im \oint_{\Sigma_1} B_1}, \quad W_2(\Sigma_2, \ell) = e^{i\ell \oint_{\Sigma_2} B_2}$$

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(Q1) Can there be any IR-Universal (local) observable effect?

- (Q2) Is this very exotic / pure academic setup? Or can this arise as IR-EFT of some standard UV QFT relevant for particle physics?
 - ✓ Illustrate importance of studying carefully the effects of remnant TQFT-couplings (GGS = essential tools)

***** Anomaly Inflow

Anomaly Inflow : W/O TQFT-Coupling [Callan and Harvey '85]

$$S = \int \frac{1}{2} da \wedge * da + \int \frac{1}{2g_A^2} F_A \wedge * F_A - \int \frac{iK_A}{8\pi^2} \frac{a}{f} F_A \wedge F_A$$



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$$\circ \ \delta S = i \int \delta^{(2)} (M_2^{st}) \wedge \left(\lambda \frac{K_A}{4\pi} F_A\right)$$



• Violation of $U(1)_A$ - inv localized on the string core

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- This 'gauge anomaly' must be cancelled by an anomalous 2d QFT on M_2^{st}
- Anomalous TQFT₄ in "bulk" + Anomalous QFT_2 on the "boundary" = 0

***** Anomaly Inflow : W/O TQFT-Coupling [Callan and Harvey '85]

$$S \supset \frac{iK_A}{8\pi^2} \int da \wedge A_1 \wedge F_A = i \int A_1 \wedge J_1$$

$$\circ * J_{1} = \frac{K_{A}}{4\pi} da \wedge F_{A}$$

$$\circ d * J_{1}(bulk) = \frac{K_{A}}{4\pi} F_{A} \wedge \delta^{(2)}(M_{2}^{st})$$

$$\circ \vec{J}_{1} \sim \nabla a \times \vec{E} \quad \text{(Hall-like current)}$$



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 $\circ * J_1 = \frac{\kappa_A}{4\pi} \, da \wedge F_A$ $\circ d * J_1(bulk) = \frac{K_A}{4\pi} F_A \wedge \delta^{(2)}(M_2^{st})$ $\circ \vec{j}_1 \sim \nabla a \times \vec{E}$ (Hall-like current)

 \circ 2d chiral fermions { $\alpha_i(z,t)$ }

$$d * J_1(2d) = -\frac{K_A}{4\pi} F_A$$

 $\sum Q_i^2 = K_A$

$$S = \int \frac{1}{2} da \wedge * da + \int \frac{1}{2g_A^2} F_A \wedge * F_A - \int \frac{iK_A}{8\pi^2} \frac{a}{f} F_A \wedge F_A$$
$$+ \int \frac{in}{2\pi} B_2 \wedge dB_1 - \int \frac{iK_{AB}}{4\pi^2} \frac{a}{f} F_A \wedge F_B - \int \frac{iK_B}{8\pi^2} \frac{a}{f} F_B \wedge F_B$$

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$$A_1 \rightarrow A_1 + d\lambda_A$$

 $\delta_A S = i \int \delta^{(2)} (M_2^{st}) \wedge \lambda_A \left(\frac{K_A}{4\pi} F_A + \frac{K_{AB}}{2\pi} F_B \right)$

* Anomaly Inflow : With TQFT-Coupling [Brennan, Hong, Wang '22]

$$S = \int \frac{1}{2} da \wedge * da + \int \frac{1}{2g_A^2} F_A \wedge * F_A - \int \frac{iK_A}{8\pi^2} \frac{a}{f} F_A \wedge F_A$$
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2. $B_1 \rightarrow B_1 + d\lambda_B$, $\lambda_B = \frac{2\pi}{n}\kappa$, $\kappa = 0, 1, \cdots, n-1$ $\delta_B S = i \int \delta^{(2)}(M_2^{st}) \wedge \lambda_B \left(\frac{K_{AB}}{2\pi}F_A + \frac{K_B}{4\pi}F_B\right)$

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$$J^{\mu} \rightarrow S \supset -i \int A_{\mu} J^{\mu}$$



$$J^{\mu} \rightarrow S \supset -i \int A_{\mu} J^{\mu} \rightarrow S \supset -i \int A \wedge * J + \frac{1}{2g^2} \int F \wedge * F$$



$$S \supset -i \int A_1 \wedge J_1 = \frac{1}{2} \int (da - f \mathcal{A}_1) \wedge (da - f \mathcal{A}_1)$$



$$S \supset -i \int A_1 \wedge * J_1 = \frac{1}{2} \int (da - f \mathcal{A}_1) \wedge * (da - f \mathcal{A}_1)$$

$$\Downarrow$$

$$S \supset \frac{1}{2} \int (da - f \mathcal{A}_1) \wedge * (da - f \mathcal{A}_1) + \frac{iM}{2\pi} \int \mathcal{B}_2 \wedge d\mathcal{A}_1$$

Gauged SymmetryGauge TheoryJ(a)
$$iM \over 2\pi \int \mathcal{B}_2 \wedge d\mathcal{A}_1$$

- (ii) What happen to axion strings?
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Note: axion-string = 2d defects charged under $U(1)^{(2)}$

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$$\left\langle U(\Sigma_{1}, \alpha) \ V(m, M_{2}^{st}) \right\rangle = e^{i\alpha m \operatorname{Link}(\Sigma_{1}, M_{2}^{st})} V(m, M_{2}^{st})$$



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After we gauge $Z_M^{(0)} \subset Z_K^{(0)}$:

$$U(\Sigma_1, \alpha) = e^{i\alpha \int_{\Sigma_1} \frac{da}{2\pi f}}$$
 is inv under gauge $\frac{a}{f} \rightarrow \frac{a}{f} + \frac{2\pi}{M}$ only for $\alpha = 2\pi M$

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$$\left\langle U(\Sigma_{1}, \alpha = 2\pi M) \ V(m, M_{2}^{st}) \right\rangle = e^{i2\pi Mm} \operatorname{Link}(\Sigma_{1}, M_{2}^{st}) \ V(m, M_{2}^{st})$$

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Gauging : Global Cosmic String → Local Cosmic String



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Conclusion and Outlook

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- ✓ Non-trivial TQFT-Coupling or Topological Modification of a QFT can lead to interesting observable consequences.
- ✓ Generalized Global Symmetry and their anomalies play the central role!
- ✓ This is just the beginning! (GGS in Particle Physics?)
 - > Applications: e.g. **DM** charged under topological Z_n force
 - Wealth of cosmic string physics
 - > Axion-string + BF-string + Composite-string
 - > Non-Kibble mechanism for composite string production?
 - > Gravitational Waves
 - Axion-YM (QCD) coupled to a TQFT (on-going) : very rich!
 - Non-invertible symmetries

$\begin{array}{c} \underline{A}_{n} \text{ TOFT} \\ \textbf{A}_{n} \text{ Thank you!} \\ B_{1}, B_{2} \end{array}$

(iii)
$$S \supset \frac{1}{2} \int (da - f \mathcal{A}_1) \wedge * (da - f \mathcal{A}_1) + \frac{iM}{2\pi} \int \mathcal{B}_2 \wedge d\mathcal{A}_1$$

- $\circ Z_M^{(0)} \subset Z_K^{(0)}$ gauging
- This is gauging of 0-form (acts on local a(x)) and non-linearly realized discrete symmetry
 - $(d\Phi A\Phi)^2$ vs $(d\phi A)^2$
- Still good notion of axion-local-fluctuations (vs AHM)

