

# Nuclear Entanglement Entropy in terms of Short Range Correlations

2022 SRC Collaboration Meeting  
August 4

Ehoud Pazy

# Outline

## Brief Introduction:

- What is entanglement entropy?
  - Area law
- Orbital Entanglement
- What has already been calculated for nuclear structure?
  - Neutron-Proton Entanglement entropy
  - $^4\text{He}$  and  $^6\text{He}$
- The General Contact Formalism (GCF) in a nutshell.

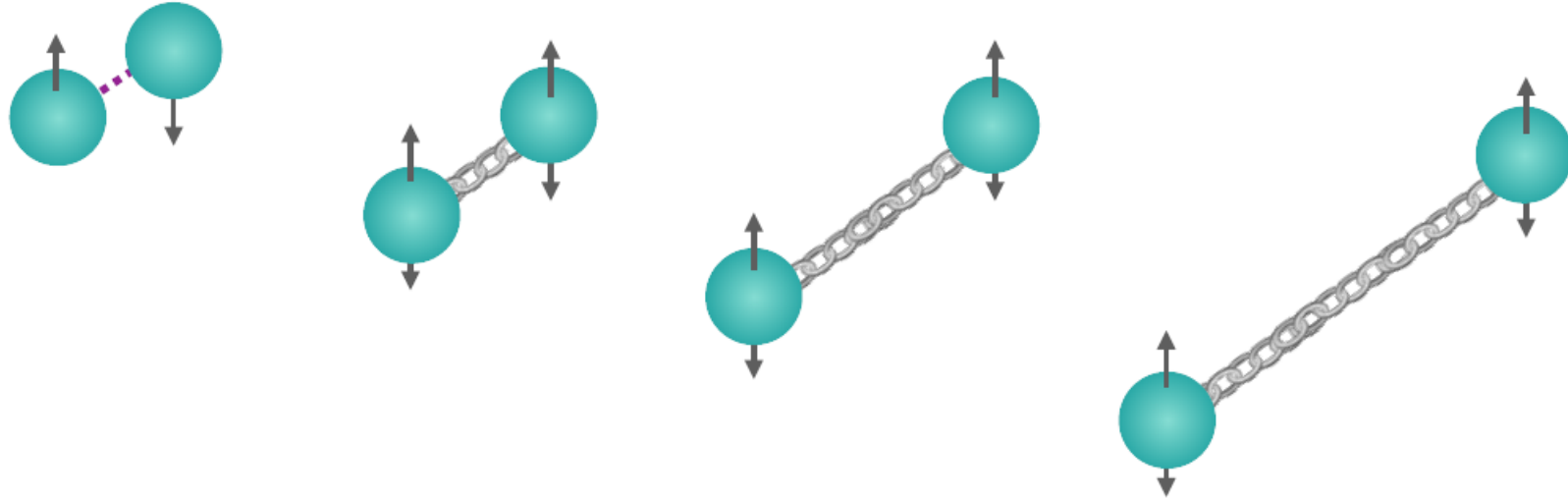
# Outline

## Calculating Entanglement Entropy for SRC:

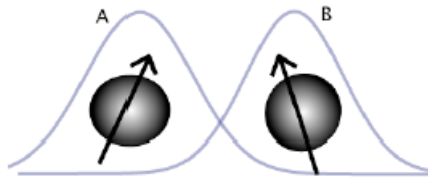
- Calculating the entanglement entropy of a single SRC pair.
- Summing up the Entanglement Entropy of SRC pairs.
- Comparing results with the  $^4\text{He}$  calculations.
- Entropy from the single particle occupation.
- The Fermi Gas part.

# Entanglement

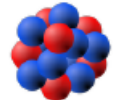
*Entanglement describes the non-local, purely quantum correlations of a system*



*Entanglement at short distance, when particles have overlapping wave functions:*



*ex: nucleons in the nucleus*



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# Entanglement Entropy



Is not a thermal entropy that originates  
from a lack of knowledge about the microstate of the system.

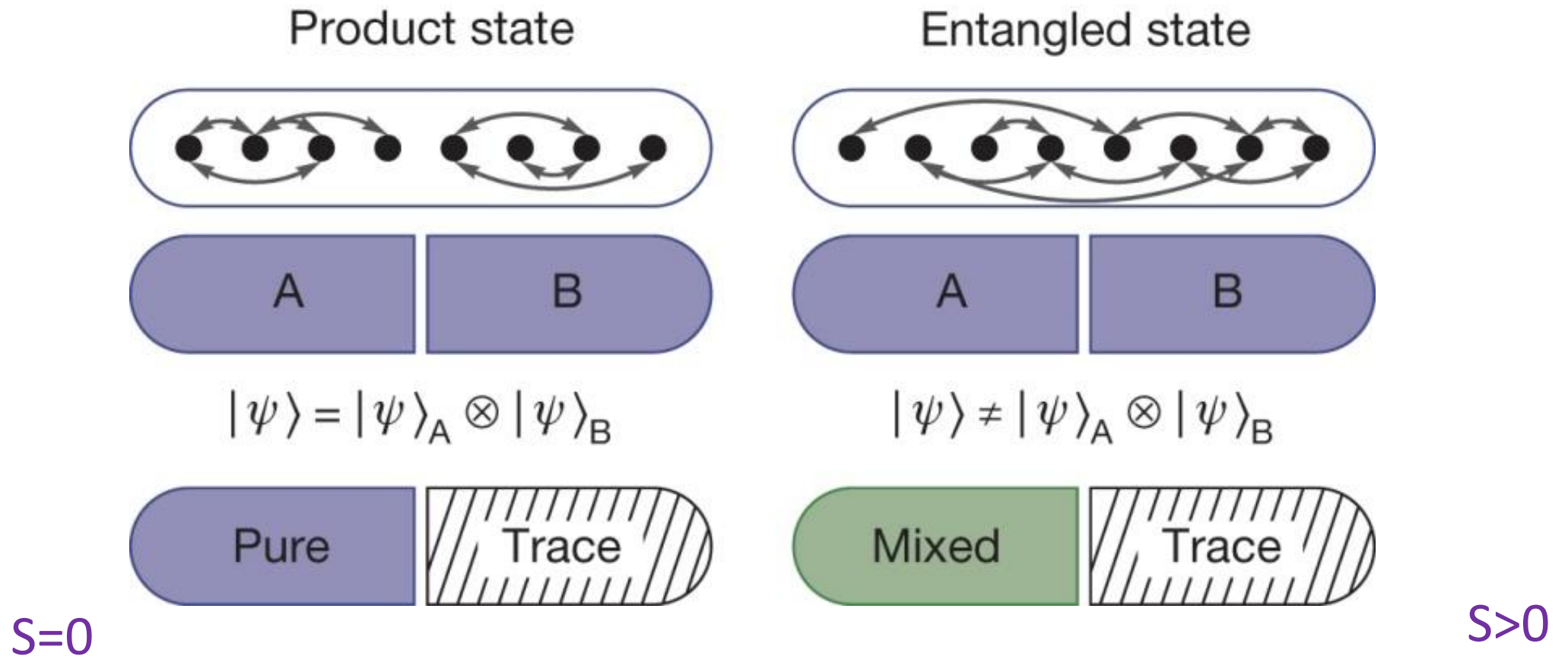
Results from tracing out part of the system

This entropy arises because of the quantum correlations-Entanglement

Even at zero temperature we will encounter a non-zero entropy!

.

# Entanglement Entropy: Product versus Entangled



# Definition of Entanglement Entropy

Divide a given quantum system into two parts **A** and **B**.  
Then the total Hilbert space becomes factorized

$$H_{tot} = H_A \otimes H_B$$

We define the reduced density matrix  $\rho_A$  for **A** by

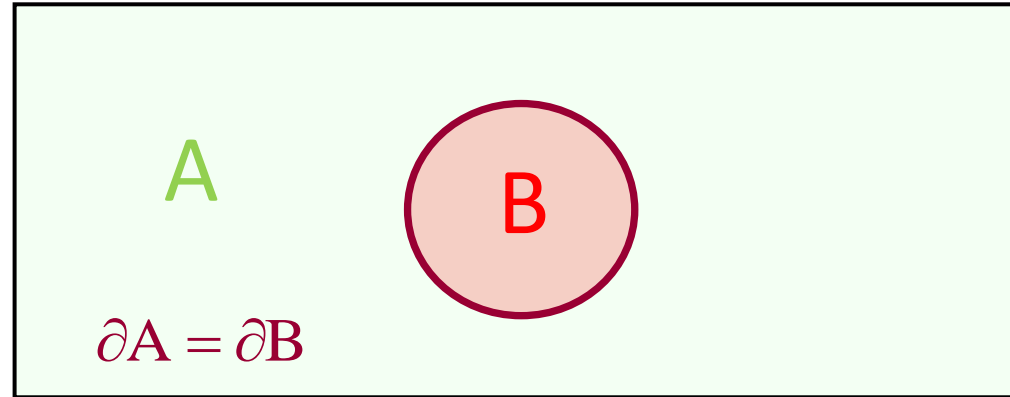
$$\rho_A = \text{Tr}_B \rho_{tot}$$

Tracing over the Hilbert space of **B** .

Now the entanglement entropy  $S_A$  is defined by the  
von Neumann entropy

$$S_A = -\text{Tr}_A \rho_A \log \rho_A$$

# Entangled Entropy Area Law



Expect that the entropy of a distinguished region **B**, be extensive  
Such a behavior is referred to as a volume scaling and is observed for thermal states.

One typically finds an area law, or an area law with a small (often logarithmic) correction:

The scaling of the entropy of a region is merely linear in the boundary area of the region.

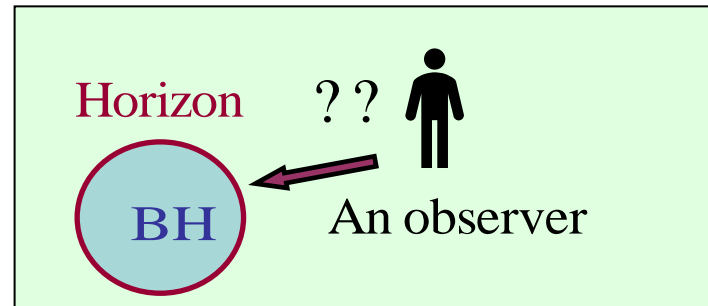
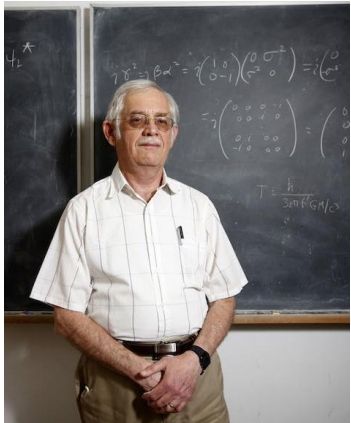


# The Holographic Principle and Black Hole Entropy

The Bekenstein-Hawking entropy of a black hole which is proportional to its boundary surface.

The holographic principle—the conjecture that the information contained in a volume of space can be represented by a theory which lives in the boundary of that region—could be related to the area law behavior of the entanglement entropy in microscopic theories

## An analogy with Black hole entropy



The boundary region  $\partial A^\sim$  the event horizon

# Why is Entanglement Entropy Important?

## Knowledge:

What role do genuine quantum correlations entanglement play in quantum many-body systems?

- It gives a measure of the correlations in the system. Specifically it is very useful to measure how entangled ground states are. An important measure is the entanglement entropy.
- Search for new order parameters

## Applications:


- Quantum information and quantum computing.

# Why is Entanglement Entropy Important?

## Nuclear physics:

Is there a simple picture in which we can understand nuclear properties?

Is there an efficient scheme in which to model nuclear structure for applications?

$$|\Psi\rangle = \sum_{\mu\nu} c_{\mu\nu} |p_\mu\rangle |n_\nu\rangle$$

$$|01101000\dots\rangle |10010100\dots\rangle$$

$$|\Psi\rangle = \sum_{\mu\nu} c_{\mu\nu} |p_\mu\rangle |n_\nu\rangle$$

Can we truncate for just a few components?

# Computational Impossibility

Contributed talk FM 8 : Johnson



SAN DIEGO STATE  
UNIVERSITY

Despite advances, it is easy to get to model spaces  
beyond our reach:

shells between 50 and 82 ( $0g_{7/2}$   $2s_{1/2}$   $0h_{11/2}$ )

$^{128}\text{Te}$ : dim 13 million (laptop)

$^{127}\text{I}$ : dim 1.3 billion (small supercomputer)

$^{128}\text{Xe}$ : dim 9.3 billion (supercomputer)

$^{129}\text{Cs}$ : dim 50 billion (haven't tried!)

# Computationally Important

$$|\Psi\rangle = \sum_{i_p, j_p} \Psi_{i_p, j_n} |i_p\rangle |j_n\rangle \longrightarrow |\Psi\rangle = \sum_{a, b} \tilde{\Psi}_{ab} |\pi_a\rangle |\nu_b\rangle$$

Initial Basis

Where the uncoupled basis states are Slater determinants

$$|\pi_a\rangle = \sum_{i_p} U_{ai_p}^{\pi} |i_p\rangle$$

Pure proton base

**Goal**  $|\Psi\rangle = \sum_{a, b=1}^N \tilde{\Psi}_{ab} |\pi_a\rangle |\nu_b\rangle, \quad N \ll \min(d_p, d_n) \quad |\Psi\rangle = \sum_{\alpha} c_{\alpha} |\Phi_{\alpha}\rangle \quad 1 = \sum_{\alpha} |c_{\alpha}|^2$

$d_p$  and  $d_n$  are the number of proton and neutron basis

The weights tell us how much we can truncate.

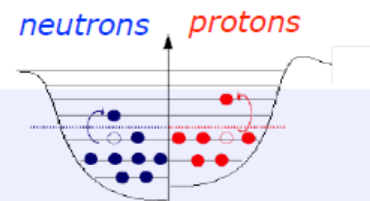
# Entanglement in the Nucleus



= Z protons + N neutrons

$$|\Psi\rangle = \sum_{\pi\nu} C_{\pi\nu} |\phi_\pi\rangle \otimes |\phi_\nu\rangle$$

$$= \sum_{n_{\pi_1} n_{\pi_2} \dots n_{\nu_1} n_{\nu_2} \dots} C_{n_{\pi_1} n_{\pi_2} \dots n_{\nu_1} n_{\nu_2} \dots} |n_{\pi_1} n_{\pi_2} \dots n_{\nu_1} n_{\nu_2} \dots\rangle$$



occupation numbers  $n_i = 0$  or  $1$

$$n_{\pi_1} + n_{\pi_2} \dots + n_{\nu_1} + n_{\nu_2} \dots = Z + N$$

Several types of entanglement are present in the nucleus:

- \* Entanglement between proton and neutron subsystems (distinguishable)

see e.g.: Papenbrock & Dean PRC 67, 051303(R) (2003), in the framework of DMRG;  
Gorton & Johnson (Gorton Master thesis 2018), in the traditional Shell Model

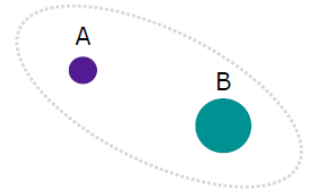
- \* Entanglement between modes (single-particle orbitals)

see e.g.: Legeza et al. PRC 92, 051303(R) (2015) in the framework of DMRG using Shell Model interactions;  
Kruppa et al. J. Phys. G: Nucl. Part. Phys. 48 025107 (2021) two-nucleon systems in the Shell Model

Caroline Robin

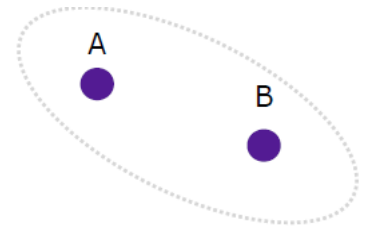
# A Complication-Orbital Entanglement

Entanglement in systems with distinguishable particles:  
Well understood – Hilbert Space has a tensor like structure



$$H_{tot} = H_A \otimes H_B \quad .$$

Entanglement in systems with indistinguishable particles:  
Not well understood-under debate



$$\mathcal{H} = \mathcal{S}(\mathcal{H}_A \otimes \mathcal{H}_B) \quad (\text{bosons}) \quad \text{or} \quad \mathcal{H} = \mathcal{A}(\mathcal{H}_A \otimes \mathcal{H}_B) \quad (\text{fermions})$$

Define entanglement between modes Rather than particles  
(second quantization)

# Calculating Orbital Entanglement

$$|\Psi\rangle = \sum_{\eta} \mathcal{A}_{\eta} |\phi_{\eta}\rangle \longrightarrow \rho^{(i)} = \begin{pmatrix} 1 - \gamma_{ii} & 0 \\ 0 & \gamma_{ii} \end{pmatrix} \longleftarrow \text{One Orbital Density Matrix}$$

$$\gamma_{ii} = \langle \Psi | a_i^{\dagger} a_i | \Psi \rangle$$

Slater determinant

$$|\phi_{\eta}\rangle = \prod_{i \in \eta}^A a_i^{\dagger} |0\rangle$$

$$S_i^{(1)} = -\text{Tr}[\rho^{(i)} \ln \rho^{(i)}] = -\sum_{k=1}^2 \omega_k^{(i)} \ln \omega_k^{(i)}$$

The  $\omega_k$  are eigenvalues of  $\rho_i$

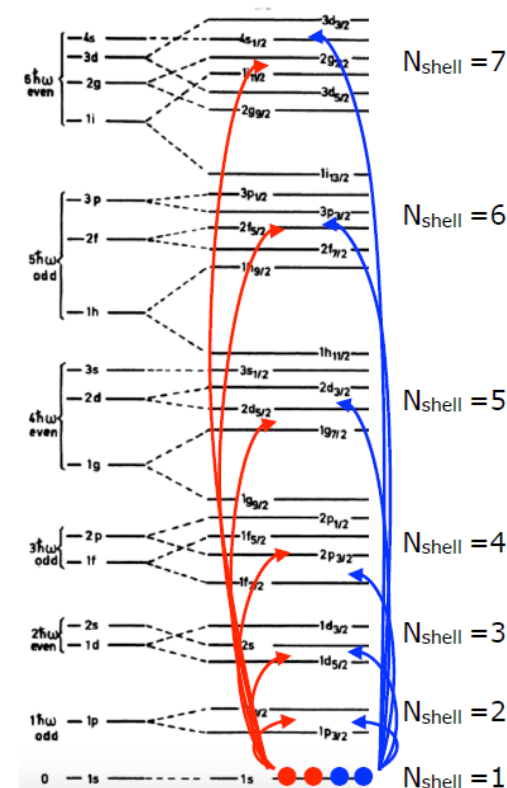
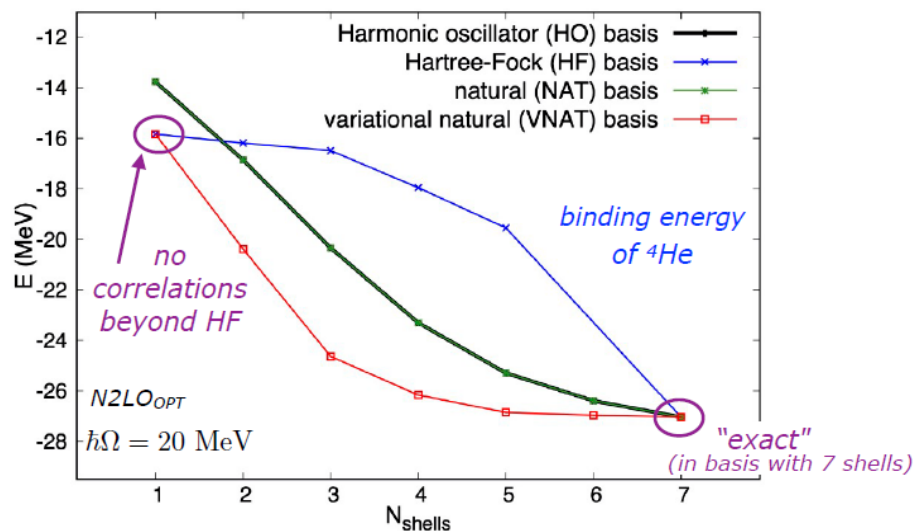


# What has been Calculated for Nuclear structure

★ application to  $^4\text{He}$  with a bare chiral interaction (2-body force, provided by P. Navrátil)

single-particle bases expanded on 7 HO major shells

full diagonalization in active space with  $N_{\text{shell}} \leq 7$



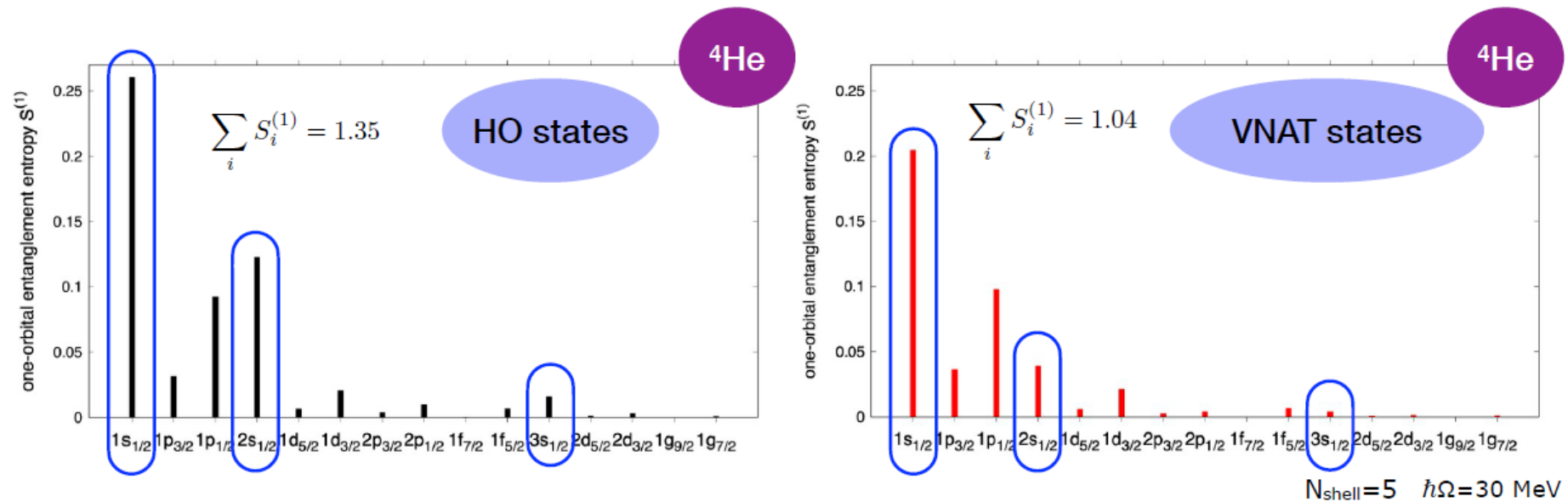
**Here: mode entanglement in (very) light nuclei with chiral EFT interaction**

C. Robin, M. J. Savage, N. Pillet, PRC 103, 034325 (2021), arXiv:2007.09157 [nucl-th] (2020)

# Single-Orbital Entanglement in $^4\text{He}$

► Single-orbital Von Neumann entropy:

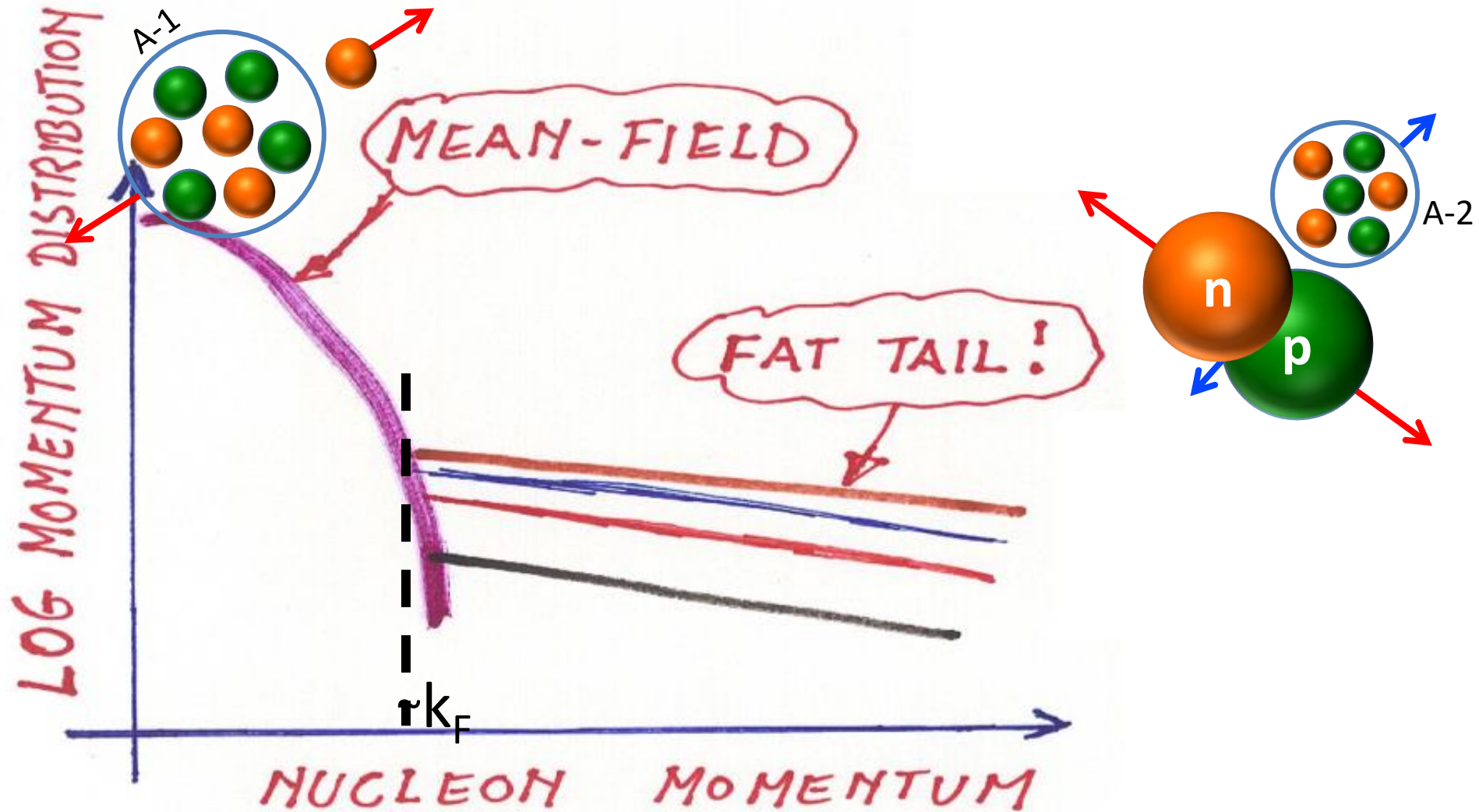
= measure of entanglement of one orbital with the rest of the system



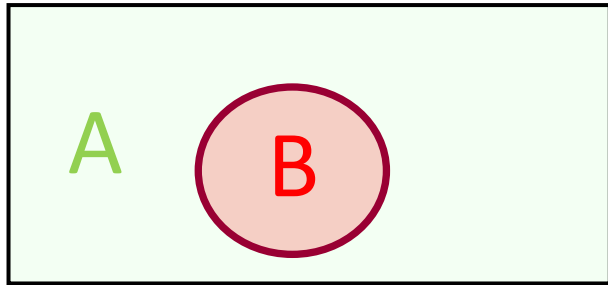
• Entanglement of the  $s$  states are decreased:  $(1s)_{\text{VNAT}}$  contains most important information

$$(1s)_{\text{VNAT}} = a_1(1s)_{\text{HO}} + a_2(2s)_{\text{HO}} + a_3(3s)_{\text{HO}} \dots$$

# General Contact Formalism (GCF)



# General Contact Formalism (GCF)



ansatz

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \varphi(\mathbf{r}_{ij}) A(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$

Two Body Universal  
wave function

in terms of:

Relative coordinate  $\mathbf{r}_{ij} = \mathbf{r}_i - \mathbf{r}_j$

A function of the A-2  
remaining nucleons

also in terms of:

Center of mass  
coordinate  $\mathbf{R}_{ij} = (\mathbf{r}_i + \mathbf{r}_j)/2$

The  
Contact

$$C_{pn} = N(A, Z) \langle A_{pn}^{\alpha_1} | A_{pn}^{\alpha_1} \rangle$$

$$\langle A | A \rangle = \int d^3 R_{ij} \prod_{k \neq i,j} d^3 r_k A^\dagger(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j}) A(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$

The contact traces over all of the degrees of freedom aside from the pair

# Calculating the SRC Entanglement Entropy

- The SRC Entanglement entropy is the sum of the Entanglement Entropy of SRC pairs.
- Calculating the entanglement entropy of a single pair.
- Comparing results with the  $^4\text{He}$  calculations.
- The Fermi Gas part.

# SRC Entanglement Entropy is a sum of the Entanglement Entropy of Single SRC

Assumption: the total SRC entanglement entropy is the sum of the number of SRC pairs-  
Meaning assuming SRC pairs are not entangled between themselves

$$S_A^{SRC} = \sum_i^N S_i^{SRC}$$

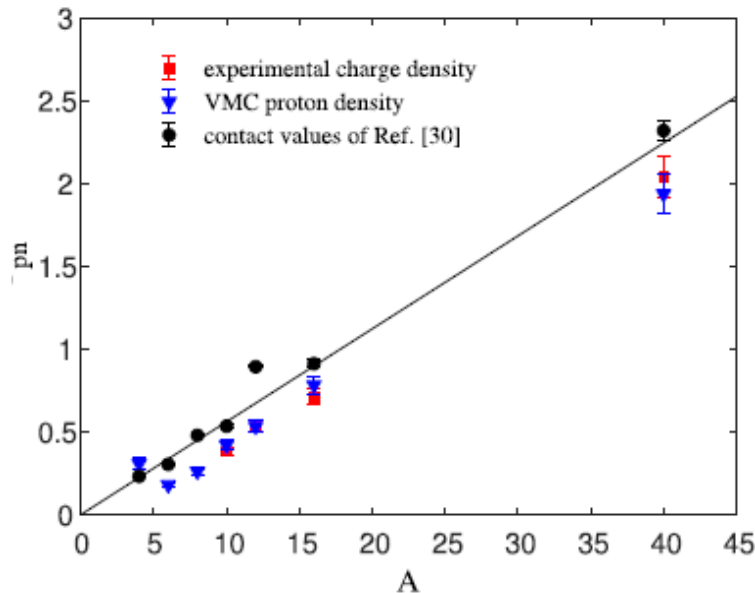
The probability for obtaining a single SRC pair is given by the normalized contact

$$c_{pn} \equiv \frac{C_{pn}}{(A/2)}$$

# The Normalized Contact

The probability for obtaining a single SRC pair

$$c_{pn} \equiv \frac{C_{pn}}{(A/2)}$$



This normalization of the contact  
Gives the fraction of the one body  
momentum density above  $K_F$

**Calculations for the single SRC entanglement entropy will be done with the normalized contact which is  $A$  independent**

# Calculating the single SRC Entanglement Entropy

$$\rho^{(i)} = \begin{pmatrix} 1 - \gamma_{ii} & 0 \\ 0 & \gamma_{ii} \end{pmatrix}$$



$$S_i^{(1)} = -\text{Tr}[\rho^{(i)} \ln \rho^{(i)}] = -\sum_{k=1}^2 \omega_k^{(i)} \ln \omega_k^{(i)}$$

One Orbital Density Matrix

$$\gamma_{ii} = \langle \Psi | a_i^\dagger a_i | \Psi \rangle$$

SRC

The probability an SRC is occupied

$$\gamma_{SRC} = c_{pn}$$

$$\rho^{(\alpha_1)} = \begin{pmatrix} 1 - \gamma_{SRC} & 0 \\ 0 & \gamma_{SRC} \end{pmatrix}$$



$$S_{pn}^{SRC} = - \left[ c_{pn} \ln \left( \frac{c_{pn}}{1 - c_{pn}} \right) + \ln (1 - c_{pn}) \right]$$



# Calculating the SRC Entanglement Entropy

$$S_A^{SRC} = -\frac{A}{2} \left[ c_{pn} \ln \left( \frac{c_{pn}}{1 - c_{pn}} \right) + \ln (1 - c_{pn}) \right]$$



The SRC Entanglement Entropy is extensive  $\sim A$

# Comparing to Previous Results for ${}^4\text{He}$

$$S_{tot}^{(1)} = \sum_i S_i^{(1)}$$

$$S_A^{SRC} = -\frac{A}{2} \left[ c_{pn} \ln \left( \frac{c_{pn}}{1 - c_{pn}} \right) + \ln (1 - c_{pn}) \right]$$

$N_{tot}$	HO	HF	NAT	VNAT
2 shells	0.596	0.270	0.596	0.441
3 shells	1.143	0.487	0.929	0.746
4 shells	1.065	0.686	0.928	1.063
5 shells	1.348	2.327	1.036	1.042
6 shells	1.264	3.434	0.972	0.963
7 shells	1.217	1.069	1.006	1.006

$$S_A^{SRC} = 0.72$$

$$c_{pn} = 0.12$$

$$S_A^{SRC} = 0.84$$

$$c_{pn}^{exp} = 0.15$$

- Proton-Proton and Neutron –Neutron pairs were not considered.
- Two orbital entanglement was not considered.

**Quite a good agreement with previous results**

# What has been Calculated for Nuclear structure

Entropy, single-particle occupation probabilities, and short-range correlations



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(Dated: June 9, 2022)

Both theoretical and experimental studies have shown that the fermion momentum distribution has a generic behavior  $n(k) = C/k^4$  at momenta larger than the Fermi momentum, due to their short-range interactions, with approximately 20% of the particles having momenta larger than the Fermi momentum. It is shown here that short-range correlations, which induce high-momentum tails of the single-particle occupation probabilities, increase the entropy of fermionic systems, which in its turn will affect the dynamics of many reactions, such as heavy-ion collisions and nuclear fission.

$$S = -g \int \frac{d^3k}{(2\pi)^3} n(k) \ln n(k) - g \int \frac{d^3k}{(2\pi)^3} [1 - n(k)] \ln[1 - n(k)], \quad (5)$$

**The Entropy in Terms of the Canonical momentum occupation function**

# The Entropy in Terms of the Canonical momentum occupation function

$$S = -g \int \frac{d^3 k}{(2\pi)^3} n(k) \ln n(k) - g \int \frac{d^3 k}{(2\pi)^3} [1 - n(k)] \ln[1 - n(k)], \quad (5)$$

$$n(k) = \eta(k_0) \begin{cases} n_{\text{mf}}(k), & \text{if } k \leq k_0 \\ n_{\text{mf}}(k_0) k_0^4 \frac{1}{k^4}, & \text{if } k_0 < k < \Lambda \end{cases}, \quad (9)$$

$$n(k) = C/k^4$$

where

$$C(k_0) = \eta(k_0) n_{\text{mf}}(k_0) k_0^4 \quad (10)$$

# Reconciling the two views: canonical momentum occupation and GCF

$$S = -g \int \frac{d^3k}{(2\pi)^3} n(k) \ln n(k) - g \int \frac{d^3k}{(2\pi)^3} [1 - n(k)] \ln[1 - n(k)], \quad (5)$$

$$n(k) = C/k^4$$

$$S = -\frac{g}{2\pi^2} \int_{k_0}^{\Lambda} \left[ \frac{C}{k^2} \ln \left( \frac{\frac{C}{k^4}}{1 - \frac{C}{k^4}} \right) + k^2 \ln \left( 1 - \frac{C}{k^4} \right) \right]$$

Estimating the integral as lower limit time  $K_0$

$$S \approx -\frac{g}{2\pi^2} \left[ \frac{C}{k_0} \ln \left( \frac{\frac{C}{k_0^4}}{1 - \frac{C}{k_0^4}} \right) + k_0^3 \ln \left( 1 - \frac{C}{k_0^4} \right) \right]$$

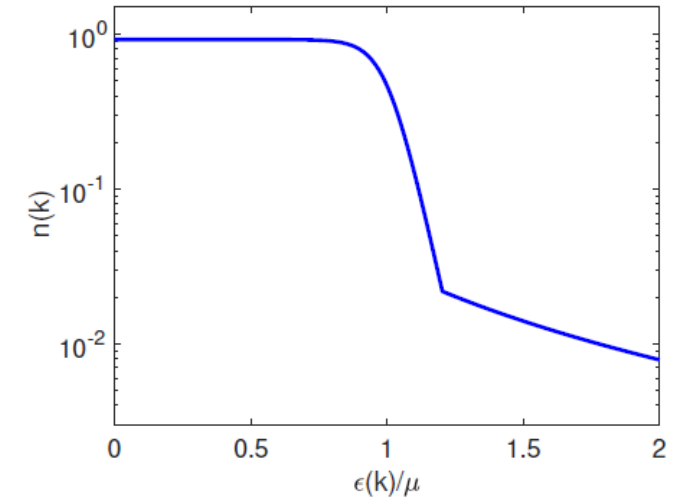
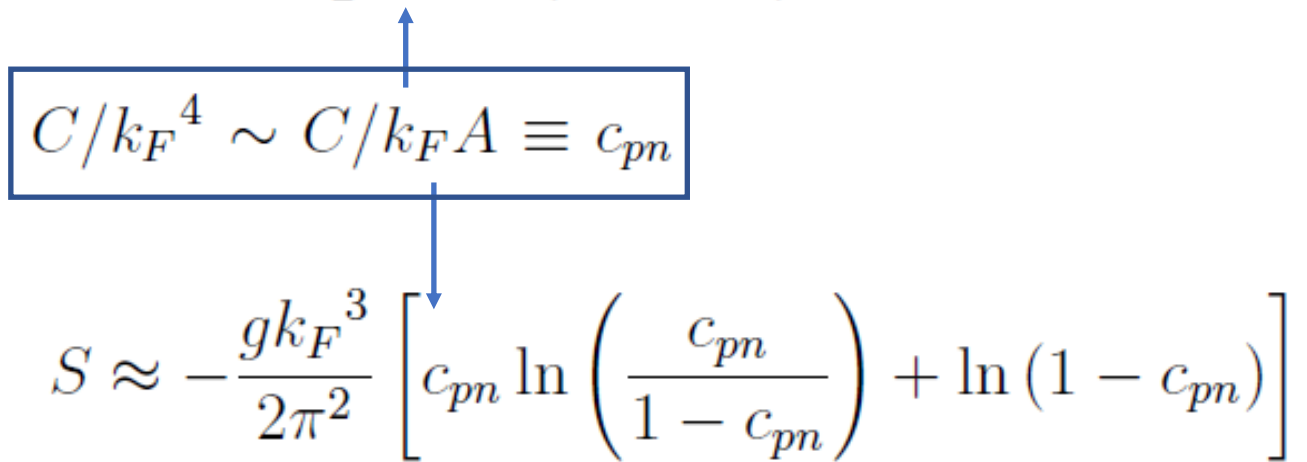


FIG. 5. The generic dependence of the canonical momentum occupation probability  $n(k)$  versus the kinetic energy  $\epsilon(k)$  after taking into the account the role of the short-range correlations, a result similar to the behavior established for finite systems in Ref. [2].

# Reconciling the two views

$$S \approx -\frac{g}{2\pi^2} \left[ \frac{C}{k_0} \ln \left( \frac{\frac{C}{k_0^4}}{1 - \frac{C}{k_0^4}} \right) + k_0^3 \ln \left( 1 - \frac{C}{k_0^4} \right) \right]$$

$C/k_F^4 \sim C/k_F A \equiv c_{pn}$

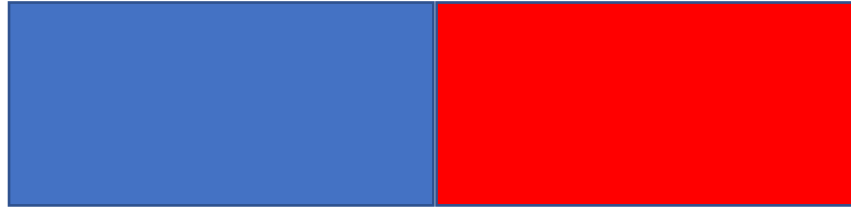


$$S \approx -\frac{gk_F^3}{2\pi^2} \left[ c_{pn} \ln \left( \frac{c_{pn}}{1 - c_{pn}} \right) + \ln (1 - c_{pn}) \right]$$

$$gk_F^3/6\pi^2 \approx A$$

# Comparing the SRC with the Fermi Sea Entanglement Entropy

$$S_{SRC|FS} = -\text{Tr}(\rho_{FS} \ln \rho_{FS}) = -\text{Tr}(\rho_{SRC} \ln \rho_{SRC})$$



# Calculating the Entanglement Entropy for the Fermi Sea

General Formula for the Entanglement Entropy of a Fermi Sea

$$S^{FS} = \frac{L^{d-1}}{(2\pi)^{d-1}} \frac{\log L}{12} \int \int |n_x \cdot n_k| dA_x dA_k$$

Nuclear structure Calculation

$$S_A^{FS} = (1 - c_{pn}) \frac{\log L}{12} \tilde{u}_s A^{\frac{2}{3}} = (1 - c_{pn}) \frac{\log A}{36} \tilde{u}_s A^{\frac{2}{3}}$$

The naïve Fermi sea calculation of the Entanglement Entropy satisfies an area law. Does not match SRC Entanglement Entropy!



# Summary

- A general expression was obtained for the SRC Entanglement Entropy.
- The SRC Entanglement Entropy was found to be extensive.
- The SRC entanglement entropy seems to fit previous  $^4\text{He}$  calculations.
- The naïve Fermi Gas Entanglement entropy approximation does not seem to work.

Thank you!