



# nEDM from the clover-on-clover formulation

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# Introduction



# Introduction

## BSM Operators

Standard model CP violation in the weak sector.

Strong CP violation from dimension 4 operators anomalously small.

- Dimension 4:
  - CP violating mass  $m\bar{\psi}\gamma_5\psi$ .
  - Topological charge  $G_{\mu\nu}\tilde{G}^{\mu\nu}$ .
- Suppressed by  $v_{EW}/M_{BSM}^2$ :
  - Quark Electric Dipole Moment  $\bar{\psi}\Sigma_{\mu\nu}\tilde{F}^{\mu\nu}\psi$ .
  - Quark Chromo-electric Dipole Moment  $\bar{\psi}\Sigma_{\mu\nu}\tilde{G}^{\mu\nu}\psi$ .
- Suppressed by  $1/M_{BSM}^2$ :
  - Gluon Chromo-electric Dipole moment:  $G_{\mu\nu}G_{\lambda\nu}\tilde{G}_{\mu\lambda}$ .
  - Various four-fermi operators.



# Introduction

## States

Asymptotic states are 'free' particles: always have a  $P$  symmetry.

If interaction does not have these symmetries, the symmetry generator will be different on different asymptotic states.

Change the nucleon interpolating operator

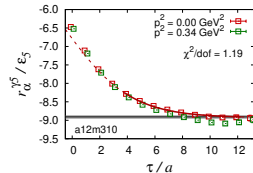
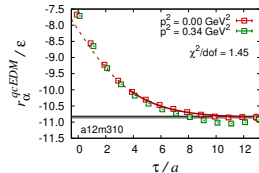
$$\hat{N} \rightarrow e^{-i\alpha_N \gamma_5} \hat{N}$$

to get standard parity for nucleon.

$\alpha_N$  can be chosen real if interactions have  $PT$  symmetry.

$\alpha_N$  can be calculated in chiral perturbation theory and as

$$\lim_{\tau \rightarrow \infty} [r_\alpha(\tau) \equiv \frac{\Im \text{Tr} \gamma_5 (1 + \gamma_4) \langle N(\tau) \bar{N}(0) \rangle}{\Re \text{Tr} (1 + \gamma_4) \langle N(\tau) \bar{N}(0) \rangle}]$$



arXiv:2304.09929 [hep-lat]

# Introduction

## Form Factors

Vector form-factors in standard representation for  $C$ ,  $P$ , and  $T$ :

Dirac  $F_1$ , Pauli  $F_2$ , Electric dipole  $F_3$ , and Anapole  $F_A$

Sachs electric  $G_E \equiv F_1 - (q^2/4M^2)F_2$  and magnetic  $G_M \equiv F_1 + F_2$

$$\begin{aligned} \langle N | V_\mu(q) | N \rangle = & \bar{u}_N \left[ \gamma_\mu F_1(q^2) + i \frac{[\gamma_\mu, \gamma_\nu]}{2} q_\nu \frac{F_2(q^2)}{2m_N} \right. \\ & + (2i m_N \gamma_5 q_\mu - \gamma_\mu \gamma_5 q^2) \frac{F_A(q^2)}{m_N^2} \\ & \left. + \frac{[\gamma_\mu, \gamma_\nu]}{2} q_\nu \gamma_5 \frac{F_3(q^2)}{2m_N} \right] u_N \end{aligned}$$

- The charge  $G_E(0) = F_1(0) = 0$ .
- $G_M(0)/2M_N = F_2(0)/2M_N$  is the (anomalous) magnetic dipole moment.
- $F_3(0)/2m_N$  is the electric dipole moment.
- $F_A$  violates PT;  $F_3$  violates CP.



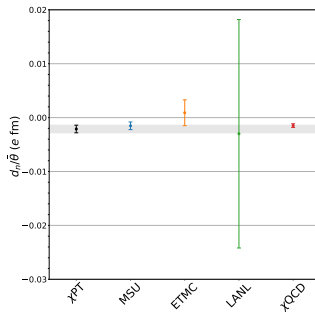
# Lattice Calculation



# Lattice Calculation

## Theta term

Calculation is essentially a reweighting of the 3-pt function by the topological charge.  
Difficult to obtain a signal at physical quark masses.  
Systematics are being studied.



E. Mereghetti, personal communication

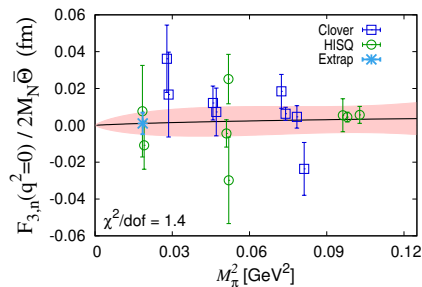
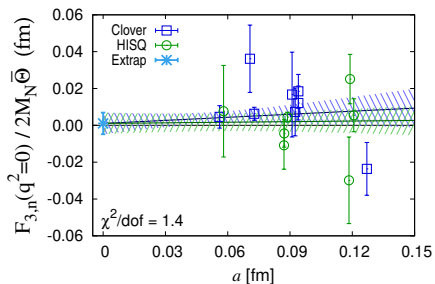


# Lattice Calculation

## Theta term: Clover-on-clover

New calculation: simultaneous fit to Clover-on-HISQ and Clover-on-Clover results  
PRELIMINARY (errors only statistical):

$$d_N = c_1 M_\pi^2 + c_2 M_\pi^2 \log\left(\frac{M_\pi^2}{M_N^2}\right) + c_3^{\text{HISQ}} a + c_3^{\text{Clover}} a = 0.0010(59)$$



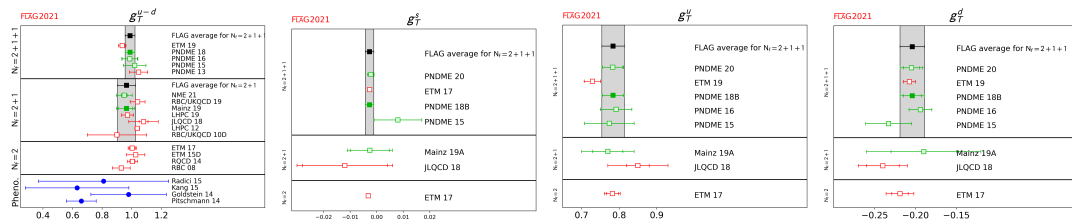


# Lattice Calculation

## Quark EDM

nEDM due to the quark EDM operator proportional to the tensor charge.

Lattice calculations dominate determination



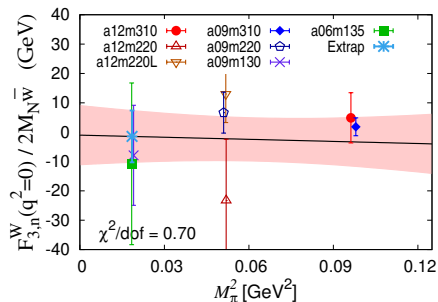
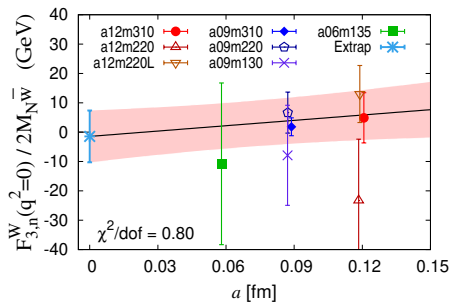
New clover-on-clover analysis in Phys. Rev. D105 (2022) 054505.



# Lattice Calculation

## Weinberg term

PRELIMINARY (errors only statistical):



Mixing with the Theta term to be removed.

Data collected with both clover-on-HISQ and clover-on-clover.

Analysis in progress.



# Lattice Calculation

## Technique: quark chromo-EDM

The quark chromo-EDM operator can be handled by the Schwinger source method

$$P \equiv \left[ \not{D} + m - \frac{r}{2} D^2 + c_{sw} \Sigma^{\mu\nu} G_{\mu\nu} \right]^{-1} \longrightarrow \left[ \not{D} + m - \frac{r}{2} D^2 + \Sigma^{\mu\nu} (c_{sw} G_{\mu\nu} + i\epsilon \tilde{G}_{\mu\nu}) \right]^{-1}$$

The fermion determinant

$$\exp \text{Tr} \ln \left[ 1 + i\epsilon \Sigma^{\mu\nu} \tilde{G}_{\mu\nu} (\not{D} + m - \frac{r}{2} D^2 + c_{sw} \Sigma^{\mu\nu} G_{\mu\nu})^{-1} \right]$$

vanishes for the isovector case

Schwinger-source does not work with gradient-flow! Instead, we can use

$$\left[ \not{D} + m - \frac{r}{2} D^2 + \Sigma^{\mu\nu} (c_{sw} G_{\mu\nu} + i\epsilon \tilde{G}_{\mu\nu}) \right]^{-1} \approx P - i\epsilon P \Sigma^{\mu\nu} \tilde{G}_{\mu\nu} P$$



# Lattice Calculation

## Three-point function

$$e^{i\epsilon} \text{ (circle with red X) } \times$$

$$\left( \begin{array}{cc} \text{Diagram 1} & \text{Diagram 2} \\ \text{Diagram 3} & \text{Diagram 4} \end{array} + \begin{array}{cc} \text{Diagram 5} & \text{Diagram 6} \\ \text{Diagram 7} & \text{Diagram 8} \end{array} \right)$$

The diagrams show various quark propagator insertions on a lattice. Diagrams 1-4 show insertions on the top quark line (u) with momentum  $P$  or  $P_\epsilon$ . Diagrams 5-8 show insertions on the bottom quark line (d) with momentum  $P$  or  $P_\epsilon$ . Diagrams 5 and 7 have yellow loops labeled  $p^{seq}$  connecting the top and bottom lines. Diagrams 6 and 8 have yellow loops labeled  $p_\epsilon^{seq}$  connecting the top and bottom lines.

$\epsilon$  small to avoid multiple insertions.

# Lattice Calculation

## quark chromo-EDM status

- HISQ analysis of isovector chromo-EDM available. (arXiv:)
- Developed techniques for subtraction of power-divergence.
- Large lattice corrections because of  $O(a^2)/am$  corrections.
- Available at arXiv:2304.09929 [hep-lat]
- Isoscalar analysis in progress.



# Quark Chromoelectric Dipole Moment



# Quark Chromoelectric Dipole Moment

## Power divergence

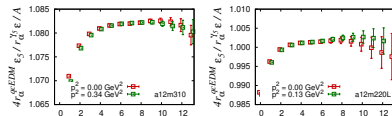
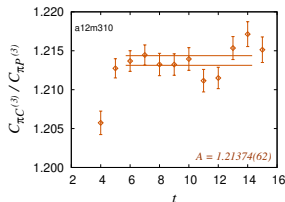
Quark chromo-EDM operator is power-divergent:

$$\tilde{C} = i\bar{\psi}\Sigma^{\mu\nu}\gamma_5 G_{\mu\nu} T^a \psi - i\frac{A}{a^2}\bar{\psi}\gamma_5 T^a \psi$$

Demanding  $\langle \Omega | \tilde{C} | \pi \rangle = 0$  fixes  $A$ :

Leading order  $\chi$ PT:

$$\alpha_N(\tilde{C}) = 0 \implies \frac{1}{A} \frac{\alpha_N(C)}{\alpha_N(\gamma_5)} = 1$$



Ensemble	$c_{SW}$	$a$ (fm)	$t$ -range	$A$
a12m310	1.05094	0.1207(11)	6–14	1.21374(62)
a12m220L	1.05091	0.1189(09)	7–14	1.21800(33)
a09m310	1.04243	0.0888(08)	8–22	0.99621(30)
a06m310	1.03493	0.0582(04)	14–30	0.77917(24)

# Quark Chromoelectric Dipole Moment

## Multiplicative renormalization

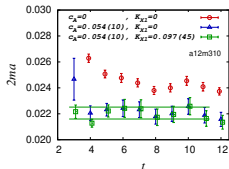
Isvector pseudoscalar can be rotated away up to  $O(a)$  effects!

We can determine the  $O(a)$  effects non-perturbatively:

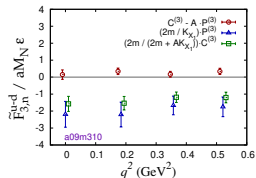
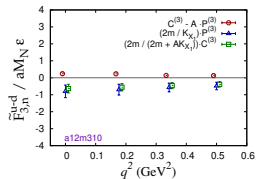
$$\frac{\langle \pi [ a \partial_\mu A^\mu - \bar{c}_A a^2 \partial^P + \bar{K} (a^2 C - AP) ] \rangle}{\langle \pi P \rangle} = 2\bar{m}a(1 + O(a^2))$$

So, on-shell zero-momentum

$$\text{M.E. of } P = \text{M.E. of } \frac{x \equiv a^2 \bar{K}}{y \equiv 2\bar{m}a + A\bar{K}} C.$$

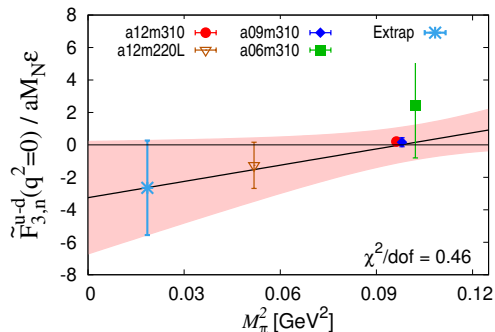
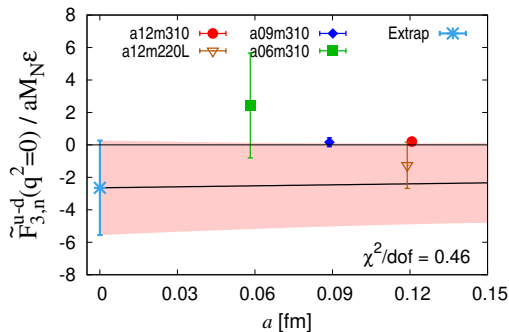


Ensemble	fit-range		$\chi^2/\text{d.o.f}$		$c_A$	$\bar{K}_{X1}$	$2\bar{m}a$	$\frac{2ma}{K_{X1}}$	$\frac{2ma}{2ma + AK_{X1}}$
	$c_A$	$\bar{K}_{X1}$	$c_A$	$\bar{K}_{X1}$					
a12m310	4-11	3-11	0.66	0.88	0.054(10)	0.097(45)	0.02205(46)	0.23(10)	0.158(58)
a12m220L	4-11	3-11	2.08	3.09	0.0342(77)	0.183(35)	0.01152(21)	0.063(12)	0.0491(86)
a09m310	5-15	4-15	0.99	1.09	0.0277(40)	0.047(15)	0.01684(15)	0.35(11)	0.263(61)
a06m310	6-20	5-20	0.29	1.53	0.0093(17)	0.0272(60)	0.010460(37)	0.385(87)	0.331(50)





# Quark Chromoelectric Dipole Moment Results



Preliminary data for power-divergent-subtracted tree-level matched operator  $\tilde{C}$  in  $\overline{\text{MS}}$  scheme 1-loop run to 2 GeV.



# Quark Chromoelectric Dipole Moment Request

Move away from mixed actions.

- $a = 0.127(2)$  fm,  $M_\pi = 285(3)$  MeV,  $32^3 \times 96$ , 2.3 M Sky-core-hrs/500 confs
- $a = 0.094(1)$  fm,  $M_\pi = 269(3)$  MeV,  $48^3 \times 128$ , 10.4 M Sky-core-hrs/500 confs
- $a = 0.0728(8)$  fm,  $M_\pi = 272(3)$  MeV,  $48^3 \times 128$ , 10.4 M Sky-core-hrs/500 confs

Using Chroma codebase.

Code works best on KNL; parts of the code not yet optimized on GPUs.

