

From the past to the future:

$$\alpha_s \text{ in } e^+e^-$$

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MIT

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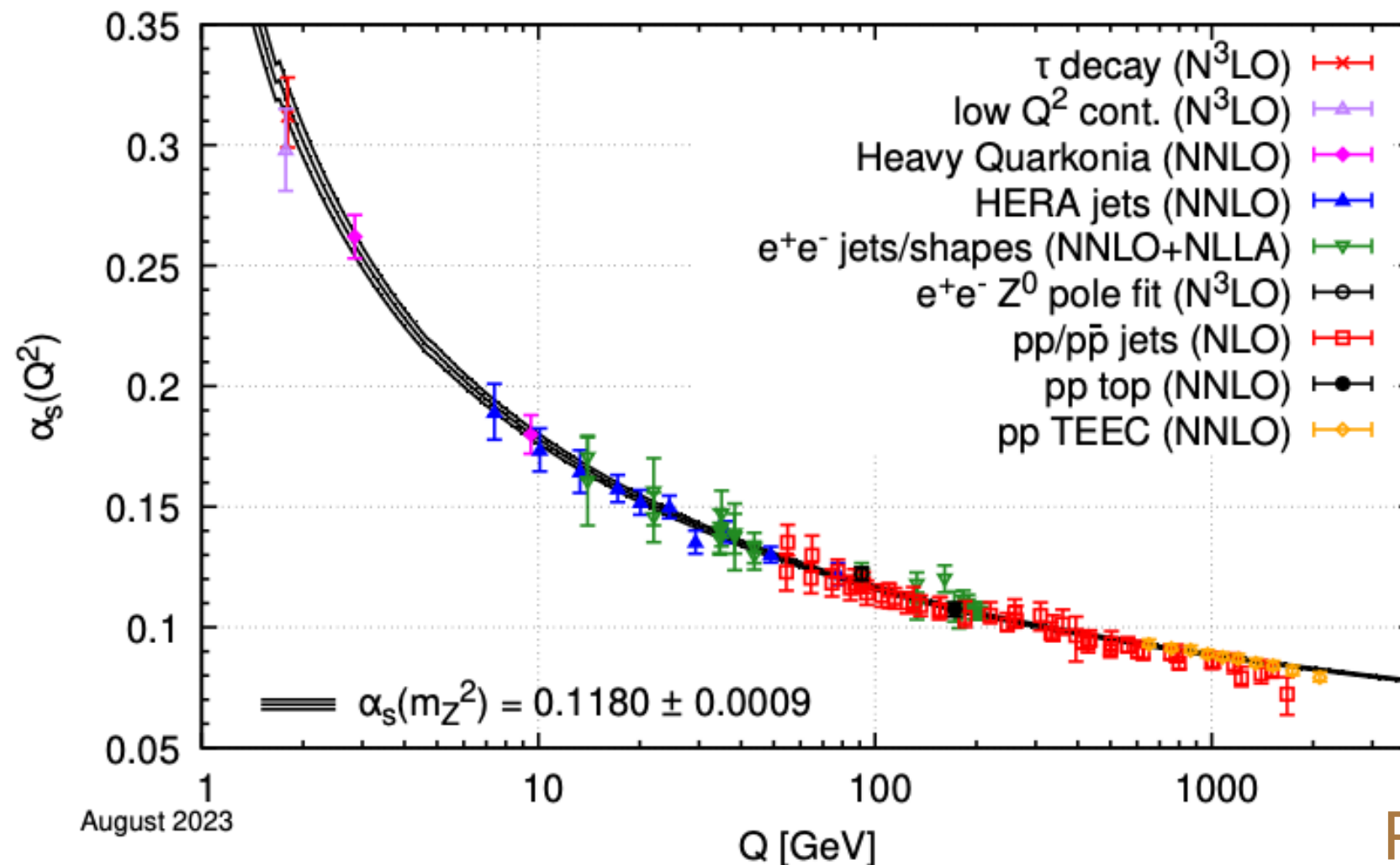
Massachusetts Institute of Technology

Outline

- α_s determinations: from LEP \rightarrow today
- Post LEP theory advances for $e^+e^- \rightarrow$ jets
- A puzzle in using precision theory to analyze LEP
- Future prospects (opportunities at FCC-ee)

Importance of α_s

- Key parameter for collider physics (MC, FO, resum., ...)
- Important for determining other SM parameters (e.weak, m_t , ...)
- Plays a role in search for new physics ($b \rightarrow s\gamma$, coupling unif., ...)



End of LEP era Inputs for $\alpha_s(m_Z)$

$$\alpha_s(m_Z) = 0.1170 \pm 0.0012$$

PDG 2005

e^+e^- collider
results

theory uncertainties
dominated

Process	Q [GeV]	$\alpha_s(Q)$	$\alpha_s(M_{Z^0})$	$\Delta\alpha_s(M_{Z^0})$		Theory
				exp.	theor.	
DIS [pol. SF]	0.7 - 8		$0.113^{+0.010}_{-0.008}$	± 0.004	$^{+0.009}_{-0.006}$	NLO
DIS [Bj-SR]	1.58	$0.375^{+0.062}_{-0.081}$	$0.121^{+0.005}_{-0.009}$	–	–	NNLO
DIS [GLS-SR]	1.73	$0.280^{+0.070}_{-0.068}$	$0.112^{+0.009}_{-0.012}$	$^{+0.008}_{-0.010}$	0.005	NNLO
τ -decays	1.78	0.345 ± 0.010	0.1215 ± 0.0012	0.0004	0.0011	NNLO
DIS [ν ; xF ₃]	2.8 - 11		$0.119^{+0.007}_{-0.006}$	0.005	$^{+0.005}_{-0.003}$	NNLO
DIS [e/ μ ; F ₂]	2 - 15		0.1166 ± 0.0022	0.0009	0.0020	NNLO
DIS [e-p \rightarrow jets]	6 - 100		0.1186 ± 0.0051	0.0011	0.0050	NLO
Υ decays	4.75	0.217 ± 0.021	0.118 ± 0.006	–	–	NNLO
$Q\bar{Q}$ states	7.5	0.1886 ± 0.0032	0.1170 ± 0.0012	0.0000	0.0012	LGT
e^+e^- [F ₂]	1.4 - 28		$0.1198^{+0.0044}_{-0.0054}$	0.0028	$^{+0.0034}_{-0.0046}$	NLO
e^+e^- [σ_{had}]	10.52	0.20 ± 0.06	$0.130^{+0.021}_{-0.029}$	$^{+0.021}_{-0.029}$	0.002	NNLO
e^+e^- [jets & shps]	14.0	$0.170^{+0.021}_{-0.017}$	$0.120^{+0.010}_{-0.008}$	0.002	$^{+0.009}_{-0.008}$	resum
e^+e^- [jets & shps]	22.0	$0.151^{+0.015}_{-0.013}$	$0.118^{+0.009}_{-0.008}$	0.003	$^{+0.009}_{-0.007}$	resum
e^+e^- [jets & shps]	35.0	$0.145^{+0.012}_{-0.007}$	$0.123^{+0.008}_{-0.006}$	0.002	$^{+0.008}_{-0.005}$	resum
e^+e^- [σ_{had}]	42.4	0.144 ± 0.029	0.126 ± 0.022	0.022	0.002	NNLO
e^+e^- [jets & shps]	44.0	$0.139^{+0.011}_{-0.008}$	$0.123^{+0.008}_{-0.006}$	0.003	$^{+0.007}_{-0.005}$	resum
e^+e^- [jets & shps]	58.0	0.132 ± 0.008	0.123 ± 0.007	0.003	0.007	resum
$p\bar{p} \rightarrow b\bar{b}X$	20.0	$0.145^{+0.018}_{-0.019}$	0.113 ± 0.011	$^{+0.007}_{-0.006}$	$^{+0.008}_{-0.009}$	NLO
$p\bar{p}, pp \rightarrow \gamma X$	24.3	$0.135^{+0.012}_{-0.008}$	$0.110^{+0.008}_{-0.005}$	0.004	$^{+0.007}_{-0.003}$	NLO
$\sigma(p\bar{p} \rightarrow \text{jets})$	40 - 250		0.118 ± 0.012	$^{+0.008}_{-0.010}$	$^{+0.009}_{-0.008}$	NLO
$e^+e^- \Gamma(Z \rightarrow \text{had})$	91.2	$0.1226^{+0.0058}_{-0.0038}$	$0.1226^{+0.0058}_{-0.0038}$	± 0.0038	$^{+0.0043}_{-0.0005}$	NNLO
e^+e^- 4-jet rate	91.2	0.1176 ± 0.0022	0.1176 ± 0.0022	0.0010	0.0020	NLO
e^+e^- [jets & shps]	91.2	0.121 ± 0.006	0.121 ± 0.006	0.001	0.006	resum
e^+e^- [jets & shps]	133	0.113 ± 0.008	0.120 ± 0.007	0.003	0.006	resum
e^+e^- [jets & shps]	161	0.109 ± 0.007	0.118 ± 0.008	0.005	0.006	resum
e^+e^- [jets & shps]	172	0.104 ± 0.007	0.114 ± 0.008	0.005	0.006	resum
e^+e^- [jets & shps]	183	0.109 ± 0.005	0.121 ± 0.006	0.002	0.005	resum
e^+e^- [jets & shps]	189	0.109 ± 0.004	0.121 ± 0.005	0.001	0.005	resum
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PDG 2005

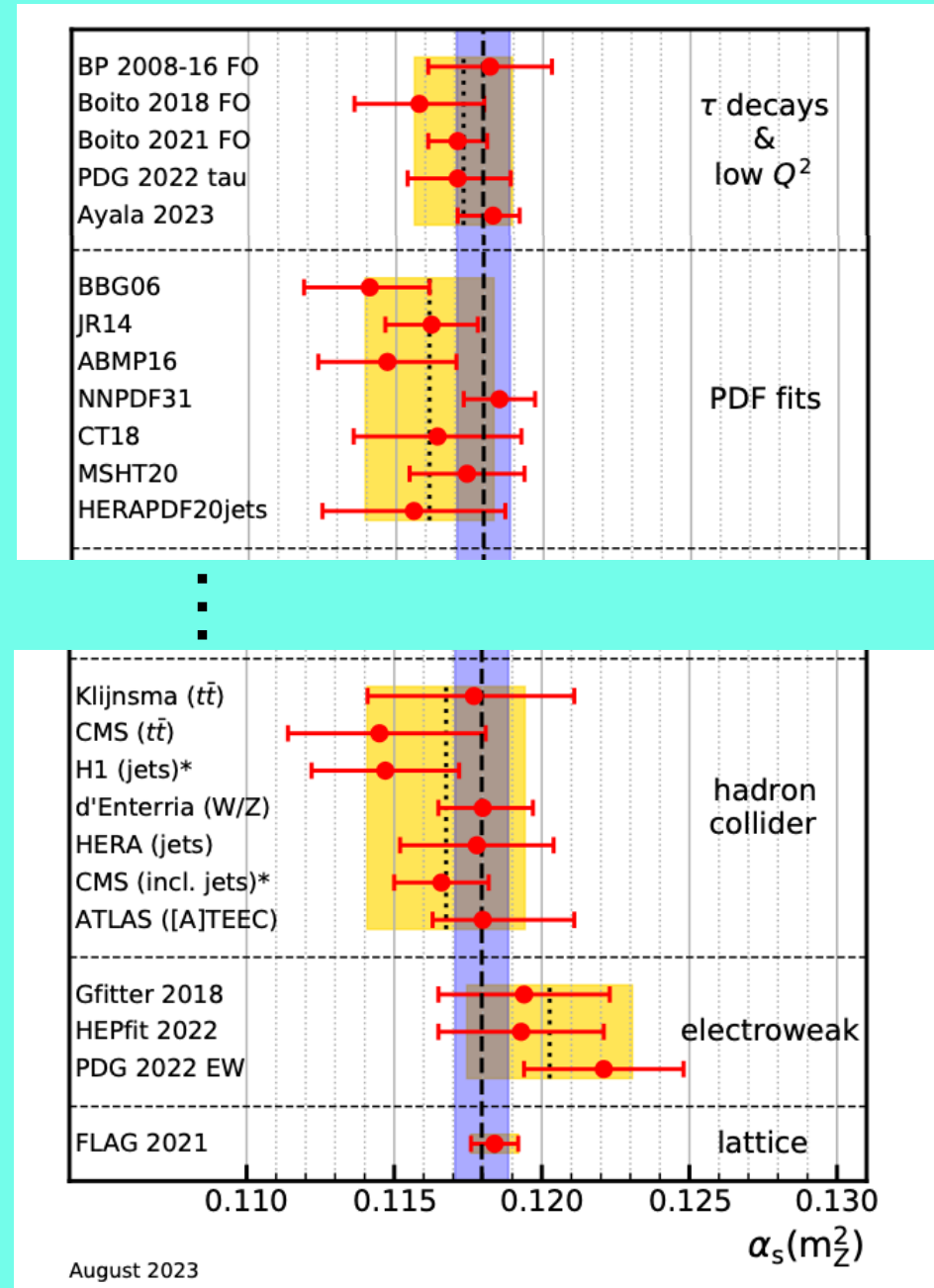
e^+e^- collider
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PDG 2023

$$\alpha_s(m_Z) = 0.1180 \pm 0.0009$$



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How has e^+e^-
theory progressed?

τ

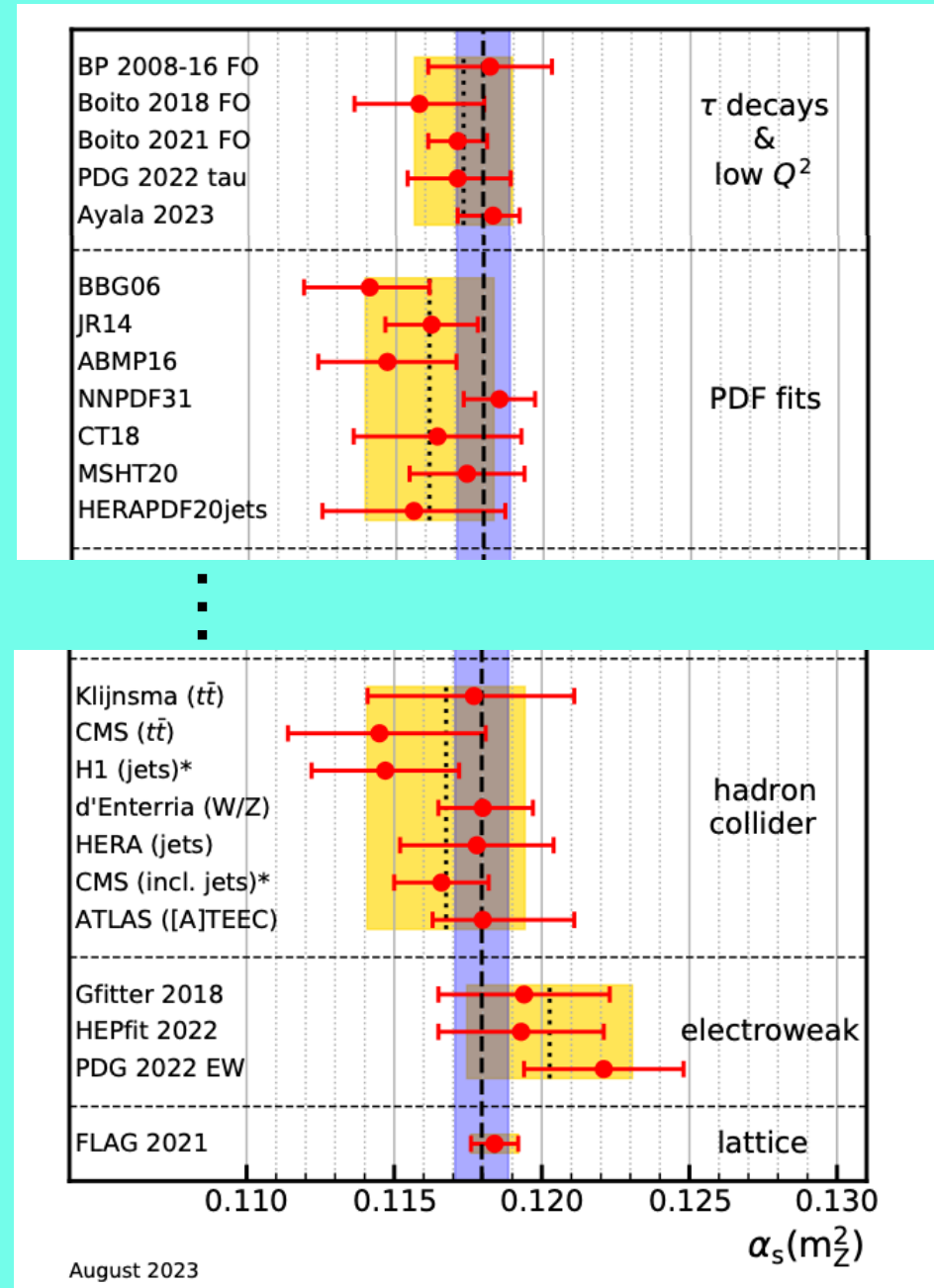
electroweak

jets*

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PDG 2023

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Jets are directly sensitive to α_s $e^+e^- \xrightarrow{Q} \text{jets}$ $(e^+e^- \rightarrow q\bar{q}g, \dots)$

Probed by event shapes (thrust, heavy-jet mass, C, ...)

$$\tau = \sum_i \tau_i$$

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In the period 2005-2015 the theoretical description of event shapes $d\sigma/d\tau$, got boosted

- Fixed-order predictions @ α_s^3 '07-'09 (Gehrmann-De Ridder et.al.; Weinzierl)
- Resummation at N³LL (SCET) '08-15 (Becher, Schwartz, Chien, Hoang, Mateu, IS, ...)
- Dijet power corrections with operators '07 (Lee, Sterman)
- Renormalon free power correction scheme '08 (Hoang, IS)

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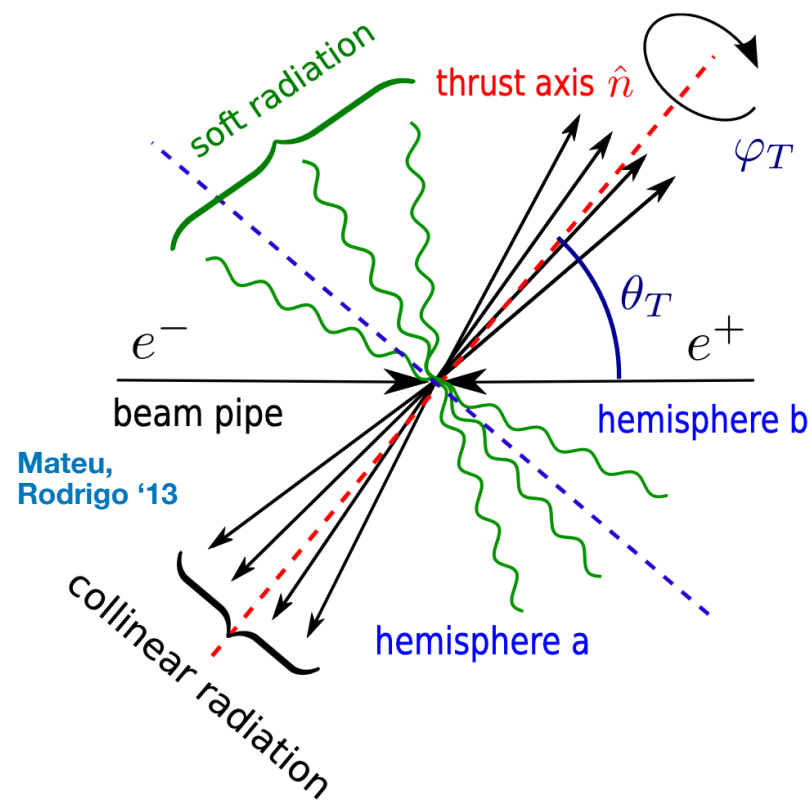
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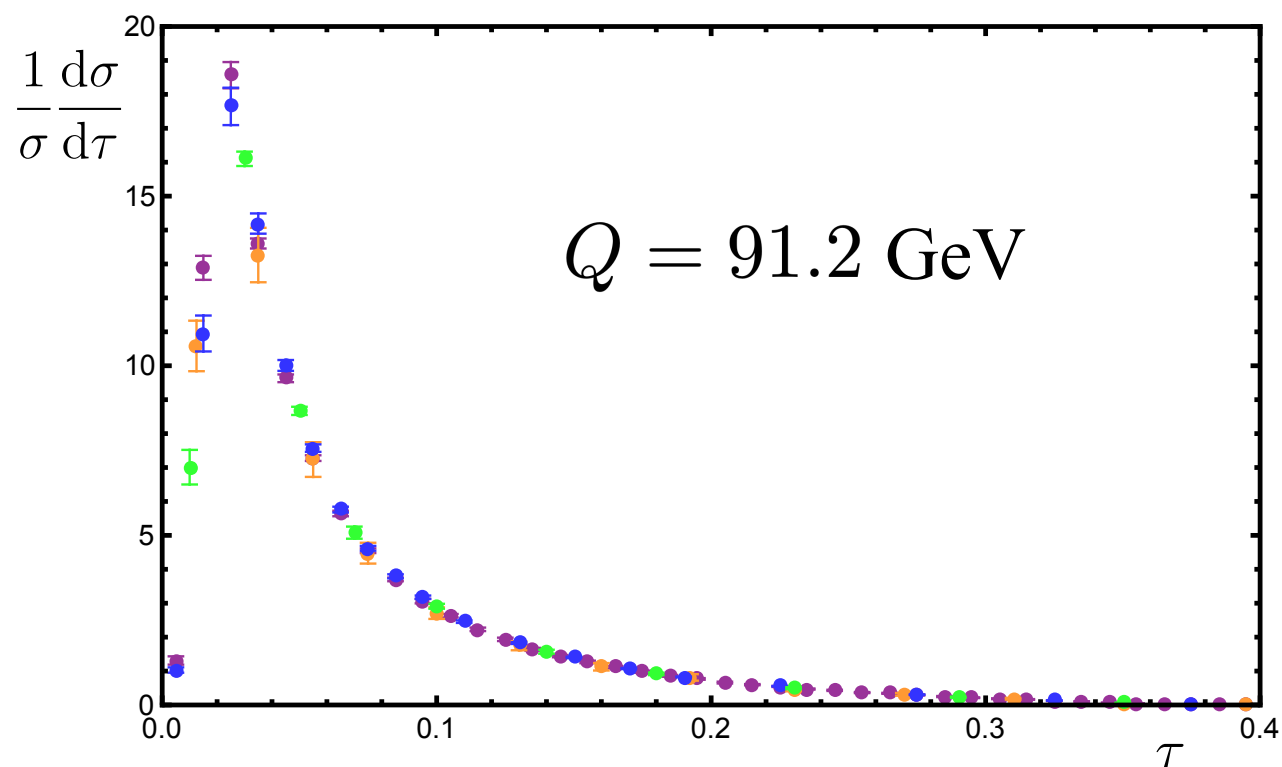
This motivated carrying out Global Fits to e^+e^- data at various c.m.(Q) to obtain $\alpha_s(m_Z)$

- Thrust (Abbate, Fickinger, Hoang, Mateu, IS 2011; Gehrmann, Luisoni, Monni 2013)
- Thrust moments (Abbate et.al. 2012)
- C-parameter (Hoang, Kolodrubetz, Mateu, IS 2015)
- Moments to many event-shapes (Gehrmann, Jaquier, Luisoni 2010)

Experimental data (eg. thrust fit)



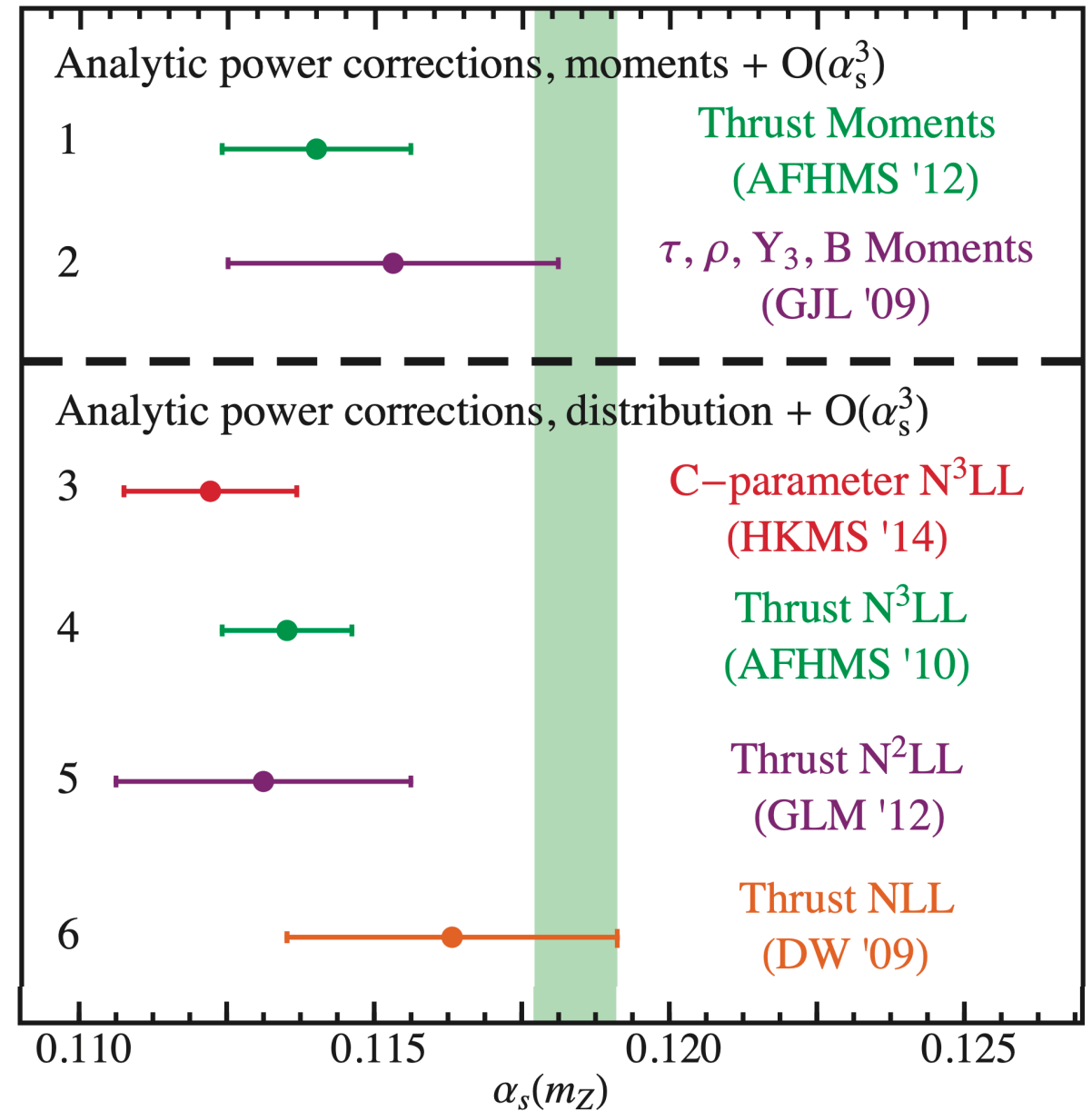
Experiment	Values of the center of mass energy Q
LEP	ALEPH {91.2, 133.0, 161.0, 172.0, 183.0, 189.0, 200.0, 206.0}
	DELPHI {45.0, 66.0, 76.0, 89.5, 91.2, 93.0, 133.0, 161.0, 172.0, 183.0, 189.0, 192.0, 196.0, 200.0, 202.0, 205.0, 207.0}
	OPAL {91.0, 133.0, 177.0, 197.0}
	L3 {41.4, 55.3, 65.4, 75.7, 82.3, 85.1, 91.2, 130.1, 136.1, 161.3, 172.3, 182.8, 188.6, 194.4, 200.0, 206.2}
SLAC	SLD {91.2}
DESY	TASSO {14.0, 22.0, 35.0, 44.0}
	JADE {35.0, 44.0}
KEK	AMY {55.2}



Standard dataset: $35 \text{ GeV} \leq Q$
 $\frac{6 \text{ GeV}}{Q} \leq \tau \leq 0.33$
 487 bins

Correlations treated in minimal overlap model

Precision e^+e^- Jet Puzzle



Although α_s was obtained with high accuracy, the central values are **uncomfortably small** (compared to world average)

Low values confirmed by 2023 analysis (Bell, Lee, Makris, Talbert, Yan)

Factorization theorems, resum logs, jet hadronization, combining regions, ...

$$\frac{d\sigma}{d\tau} = \int dk \left(\frac{d\hat{\sigma}_s}{d\tau} + \frac{d\hat{\sigma}_{ns}}{d\tau} + \frac{d\hat{\sigma}_b}{d\tau} \right) \left(\tau - \frac{k}{Q} \right) S_\tau^{\text{mod}}(k - 2\bar{\Delta}) + O\left(\sigma_0 \frac{\alpha_s \Lambda_{\text{QCD}}}{Q} \right)$$

$$\frac{\Delta\alpha_s}{\alpha_s} \sim 0.5\%$$

$$\frac{d\hat{\sigma}_s}{d\tau} = \sum_n \alpha_s^n \delta(\tau) + \sum_{n,l} \alpha_s^n \left[\frac{\ln^l \tau}{\tau} \right]_+$$

$$= H(\mu_H) \times J(\mu_J) \otimes S(\mu_S)$$

Singular partonic for massless quarks
QCD+QED final states

b mass corrections

nonperturbative soft function, $\left(\frac{\Omega_1}{Q\tau} \right)$

$$\frac{d\hat{\sigma}_{ns}}{d\tau} = \sum_{n,l} \alpha_s^n \ln^l \tau + \sum_n \alpha_s^n f_n(\tau) \quad \text{Nonsingular partonic}$$

Resummation for singular partonic

Single hadronization parameter Ω_1 when $\Lambda_{\text{QCD}} \ll Q\tau$

$$S(\tau) = S_{\text{pert}}(\tau) - S'_{\text{pert}}(\tau) \frac{2\Omega_1}{Q} \approx S_{\text{pert}}\left(\tau - \frac{2\Omega_1}{Q}\right) \quad \Omega_1 \sim \Lambda_{\text{QCD}}$$

$$\bar{\Omega}_1 \equiv \frac{1}{2N_c} \langle 0 | \text{tr} \bar{Y}_n(0) Y_n(0) i \partial_\tau Y_n^\dagger(0) \bar{Y}_n^\dagger(0) | 0 \rangle \quad \overline{\text{MS}}$$

Jets in e^+e^- with QFT

Factorization theorems, res

$$\frac{d\sigma}{d\tau} = \int dk \left(\frac{d\hat{\sigma}_s}{d\tau} + \frac{d\hat{\sigma}_{ns}}{d\tau} + \dots \right)$$

$$\frac{d\hat{\sigma}_s}{d\tau} = \sum_n \alpha_s^n \delta(\tau) + \sum_{n,l} \alpha_s^n \left[\frac{\ln^l \tau}{\tau} \right]_+$$

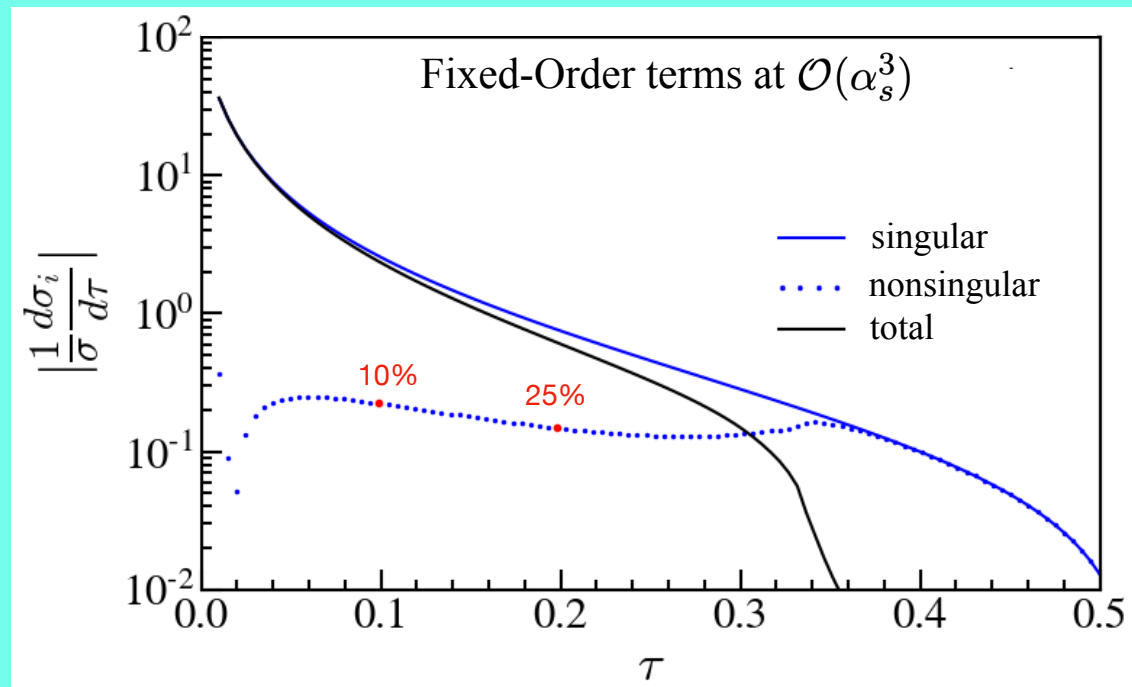
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Singular terms dominate



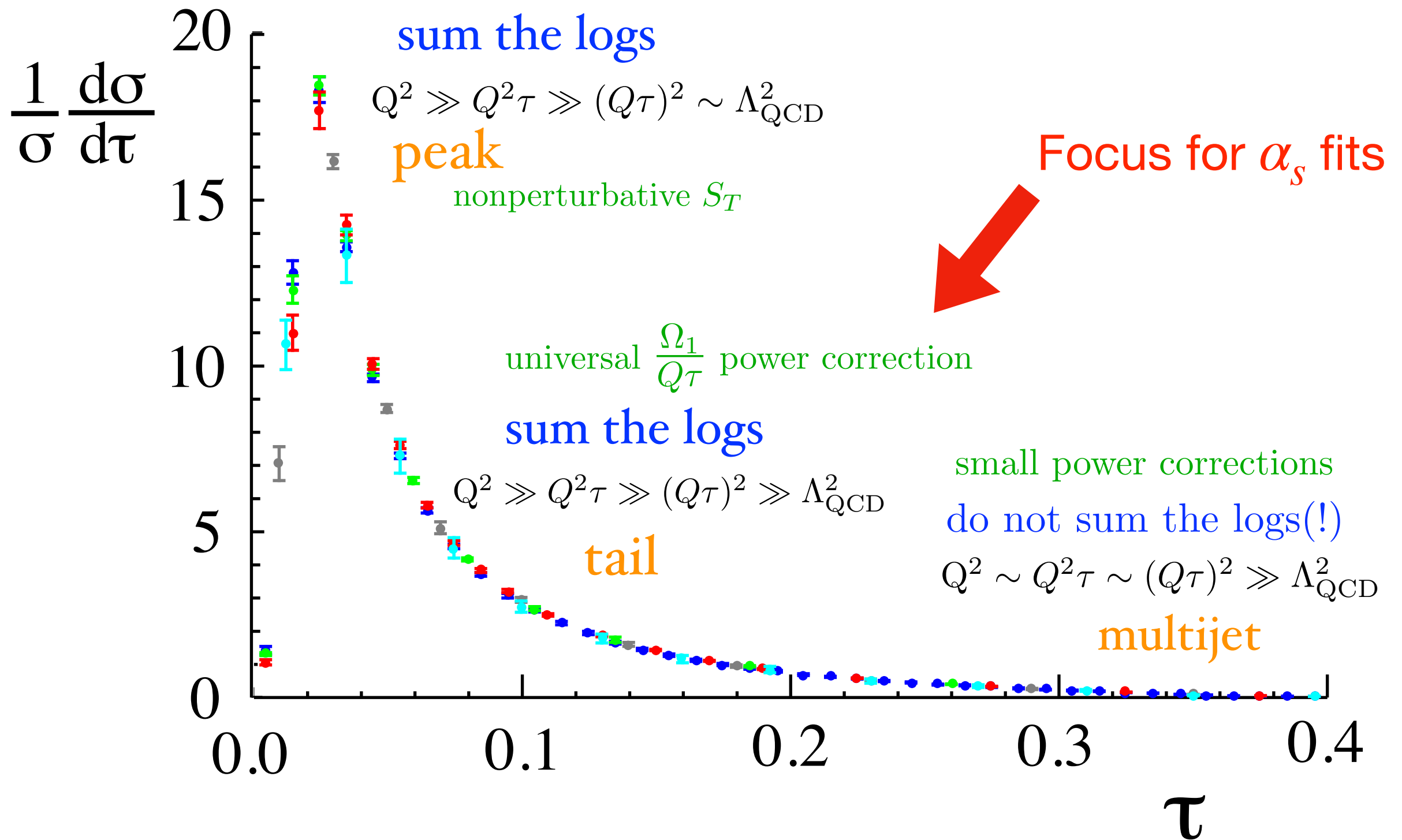
ons, ...

Single hadronization parameter Ω_1 when $\Lambda_{\text{QCD}} \ll Q\tau$

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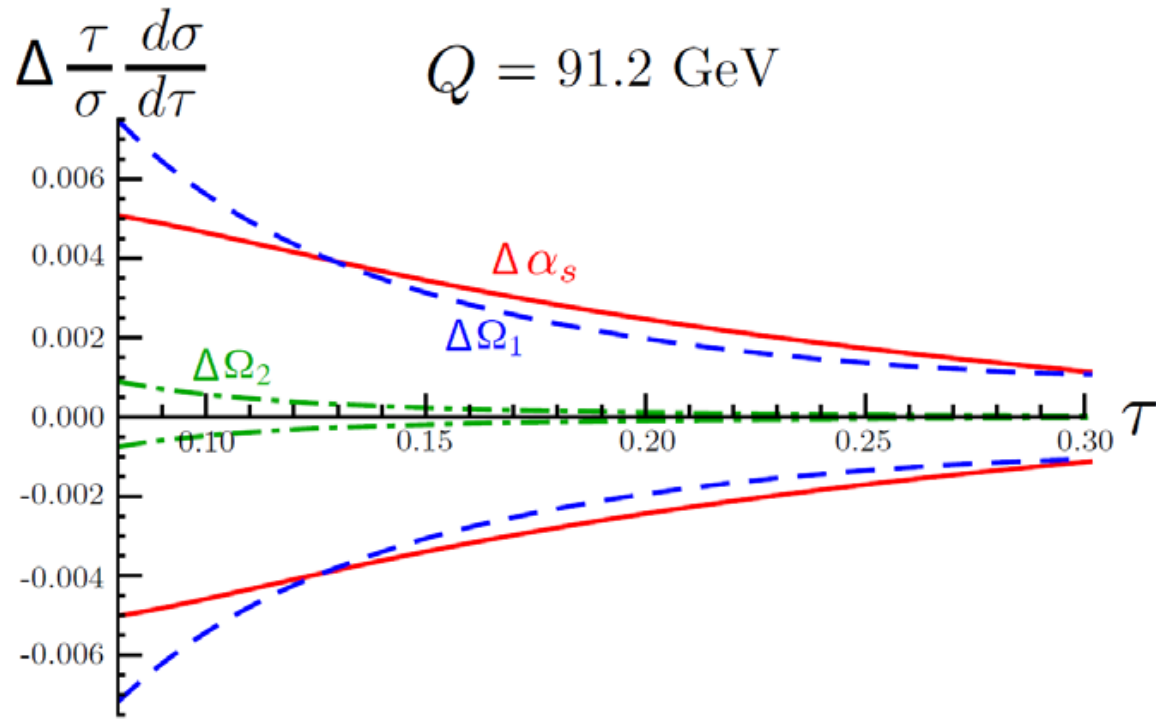
$$\bar{\Omega}_1 \equiv \frac{1}{2N_c} \langle 0 | \text{tr} \bar{Y}_n(0) Y_n(0) i \partial_\tau Y_n^+(0) \bar{Y}_n^+(0) | 0 \rangle \quad \overline{\text{MS}}$$

Three Regions:

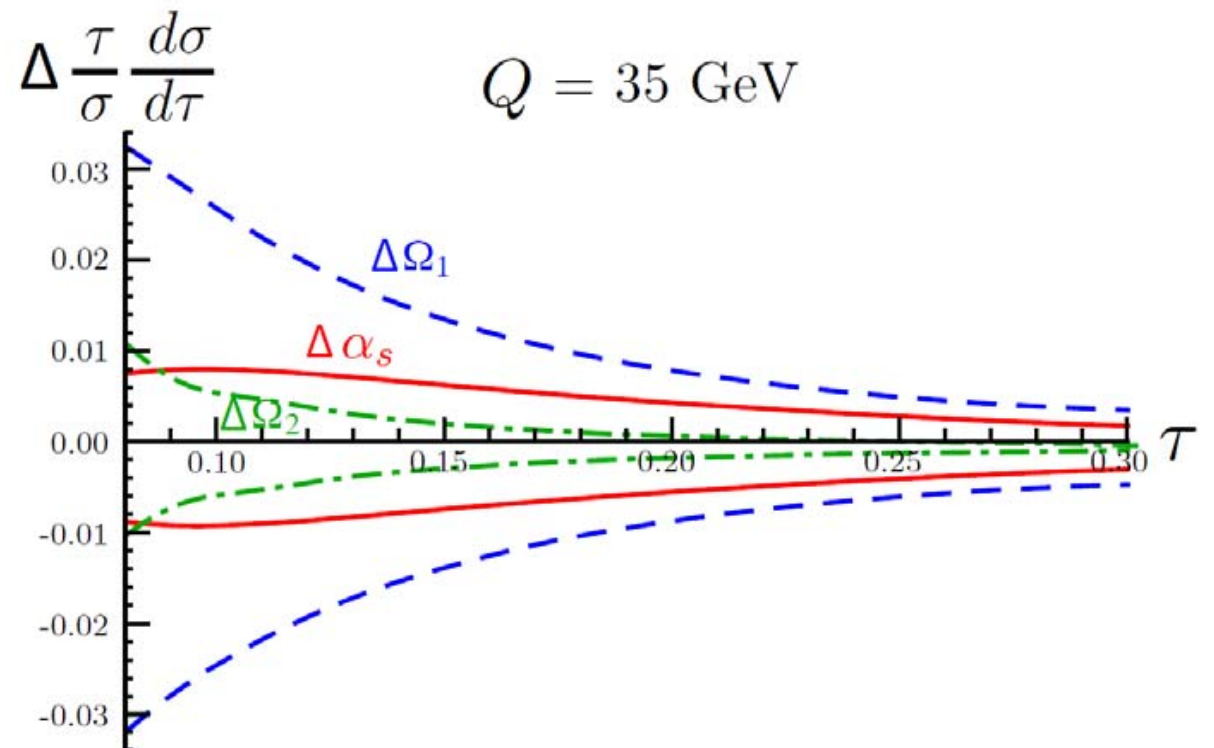
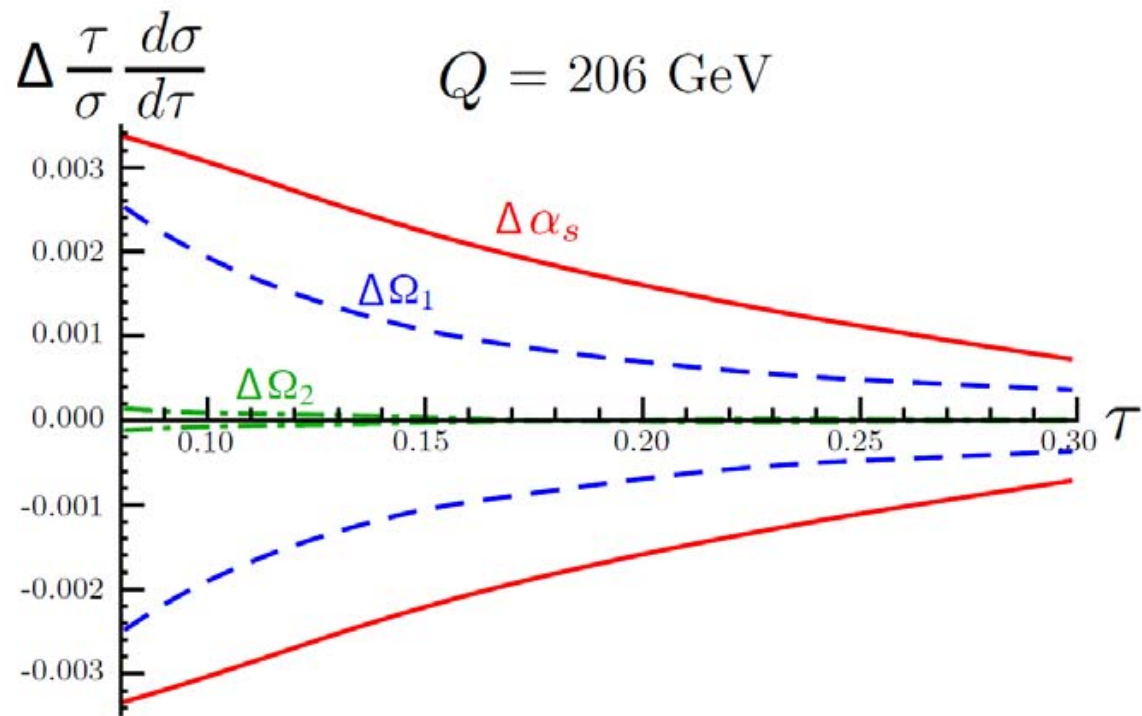


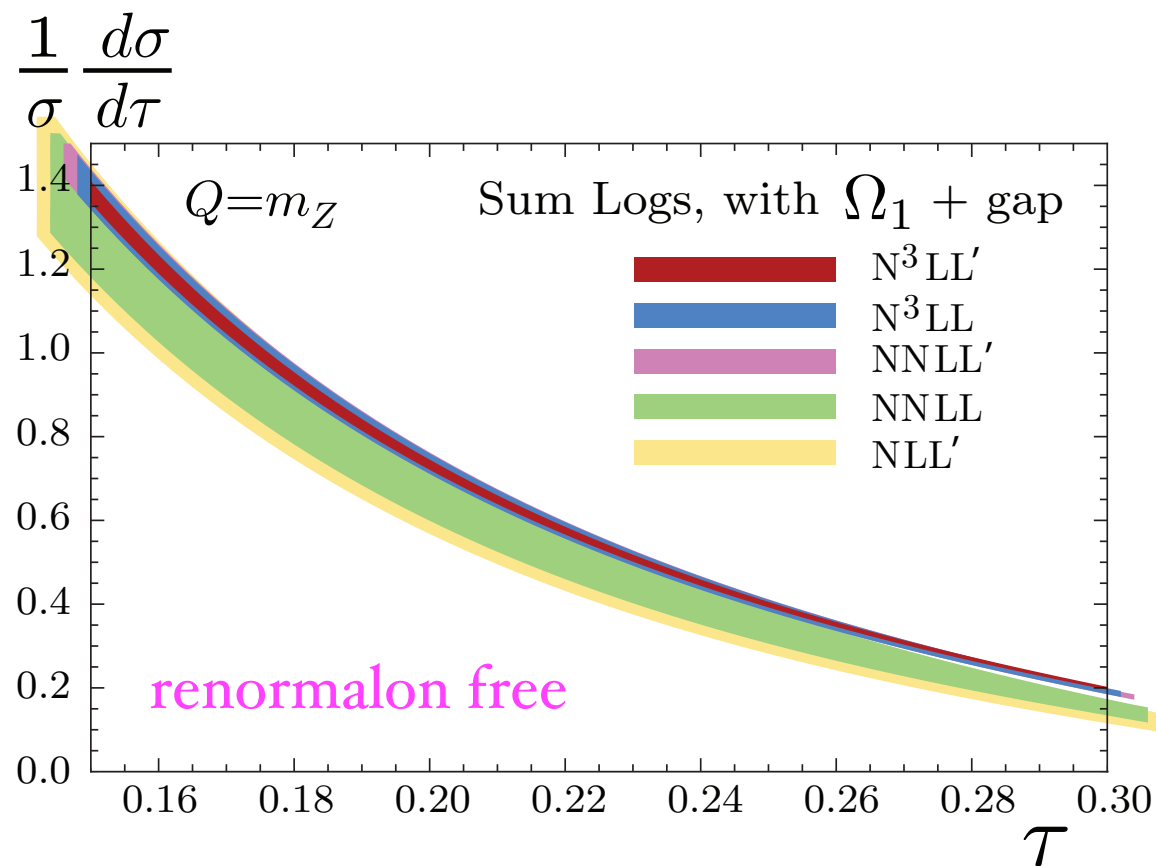
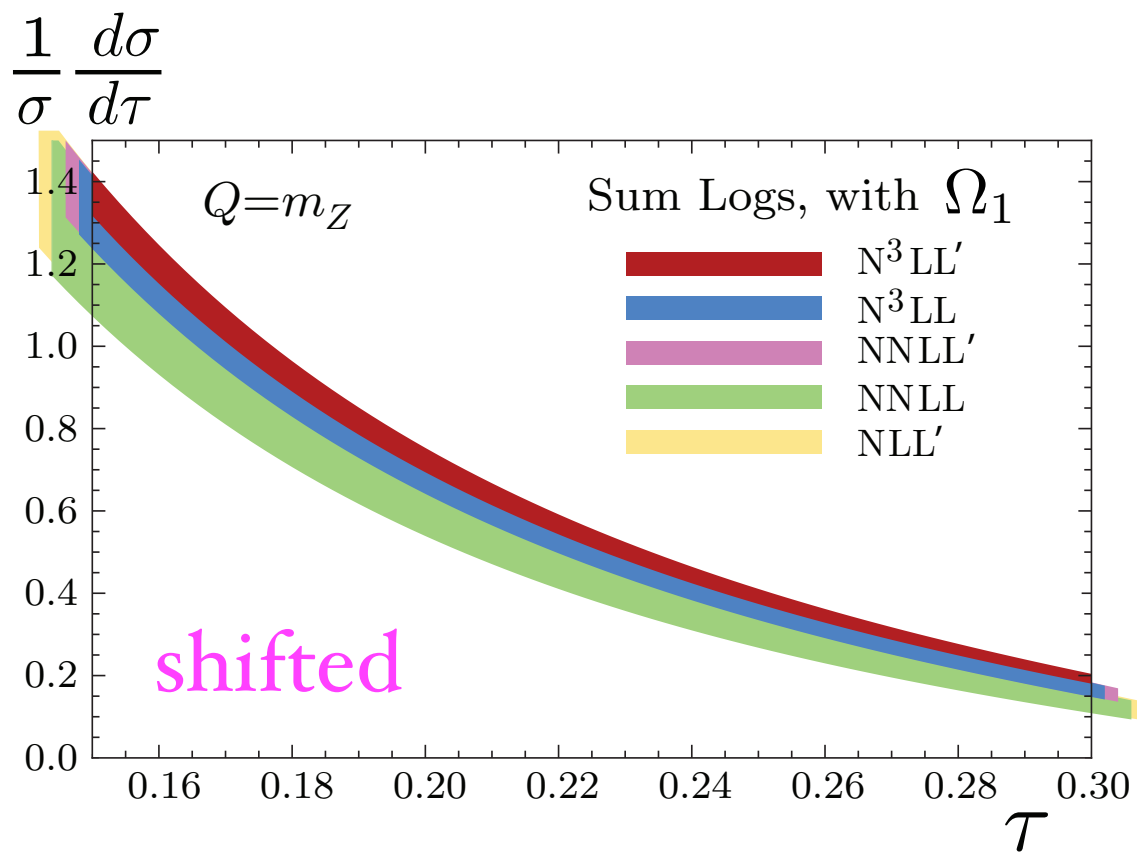
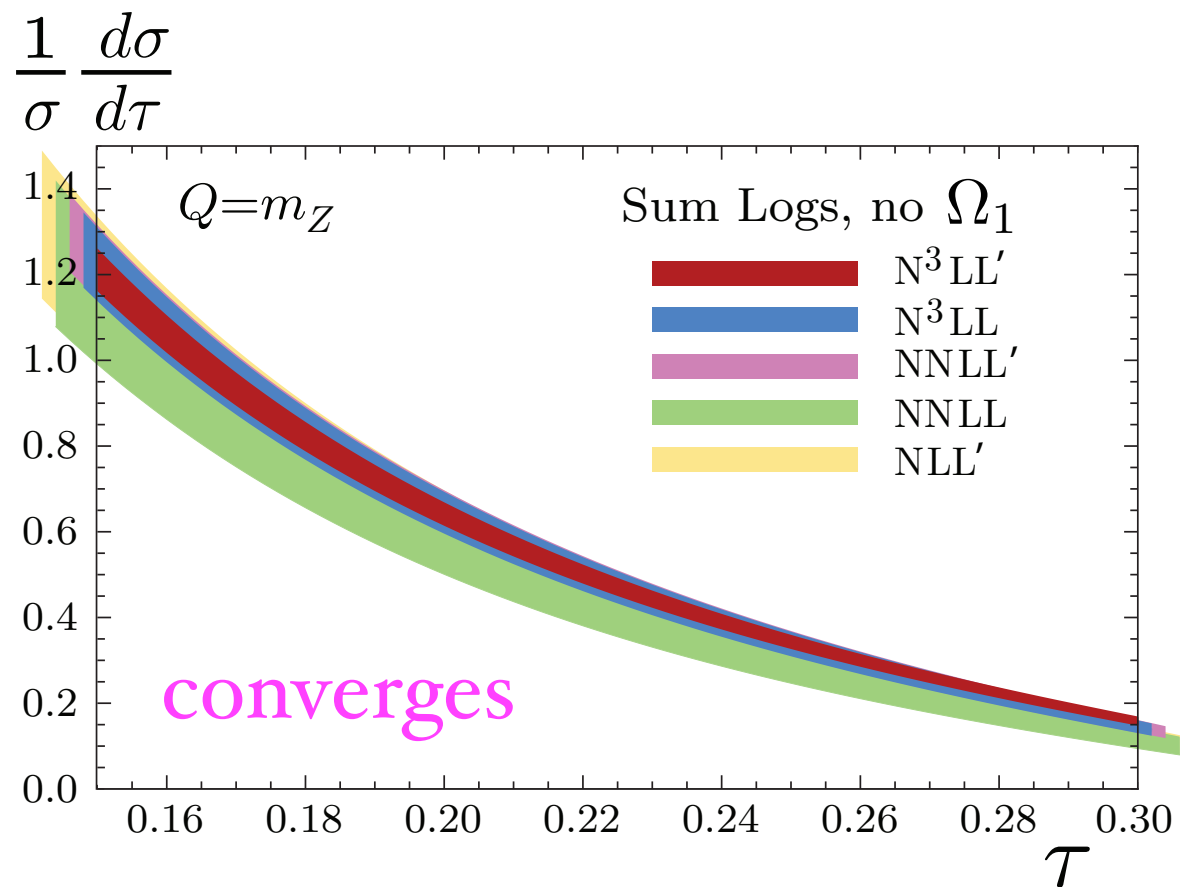
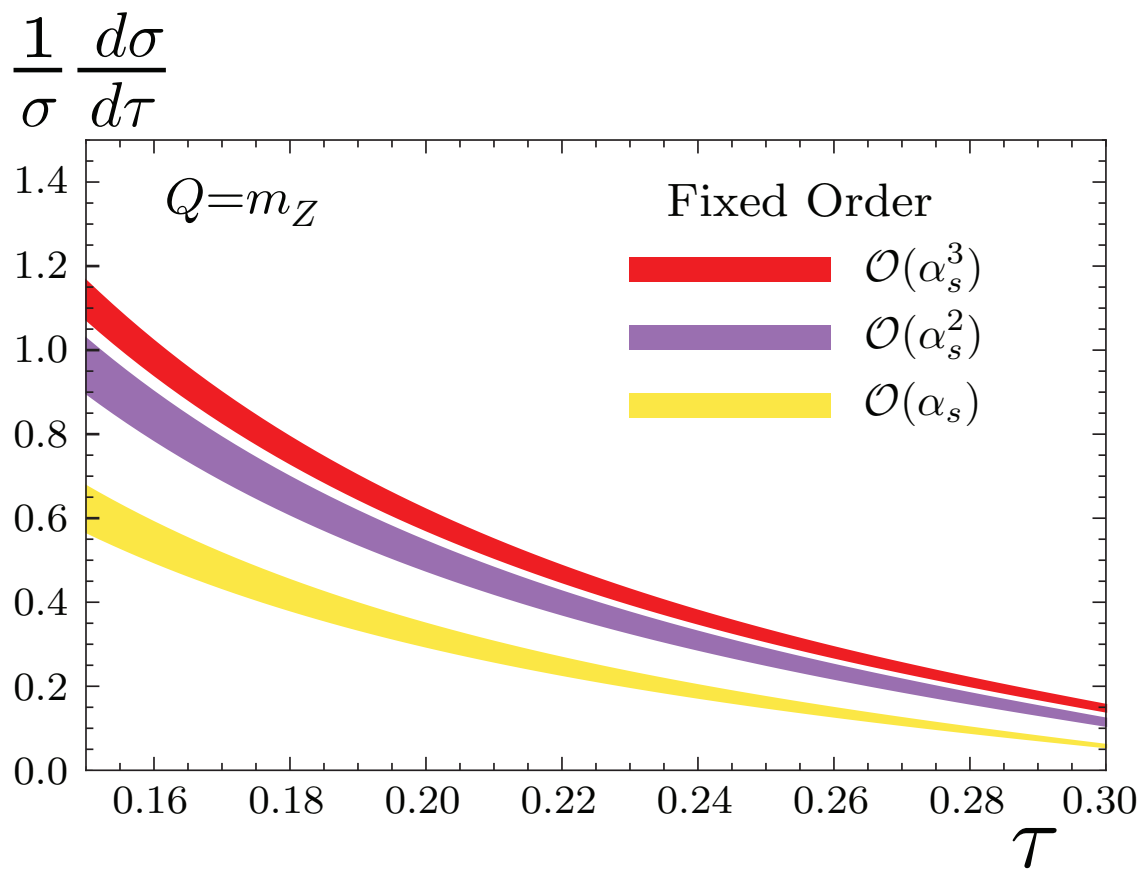
Two fit parameters $\{\alpha_s, \Omega_1\}$

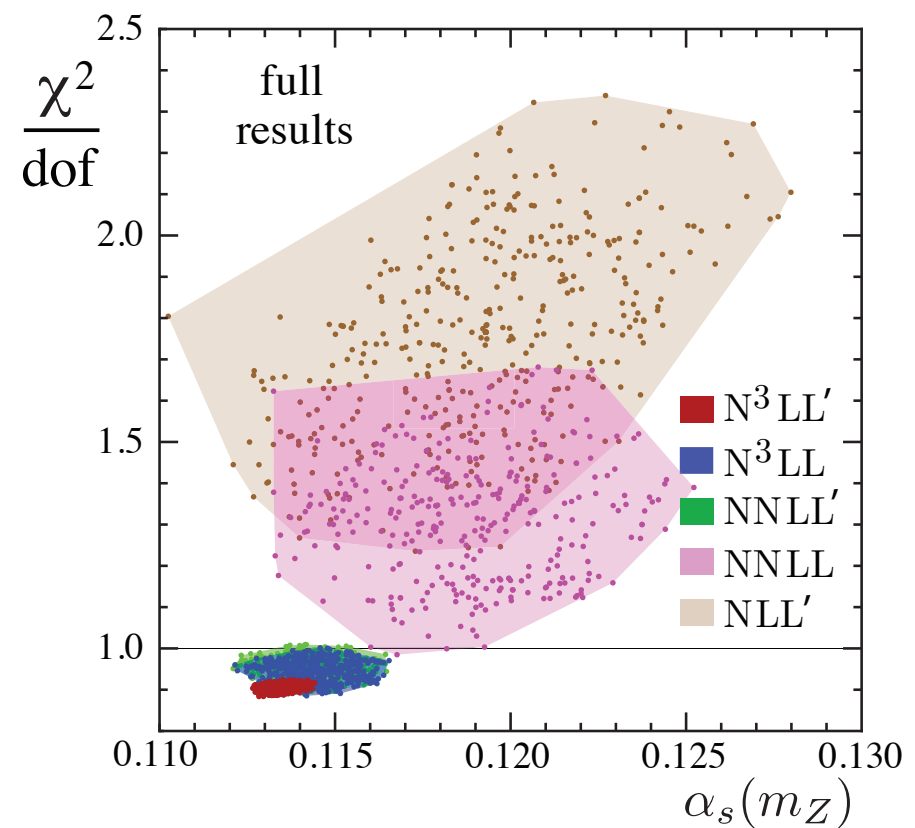
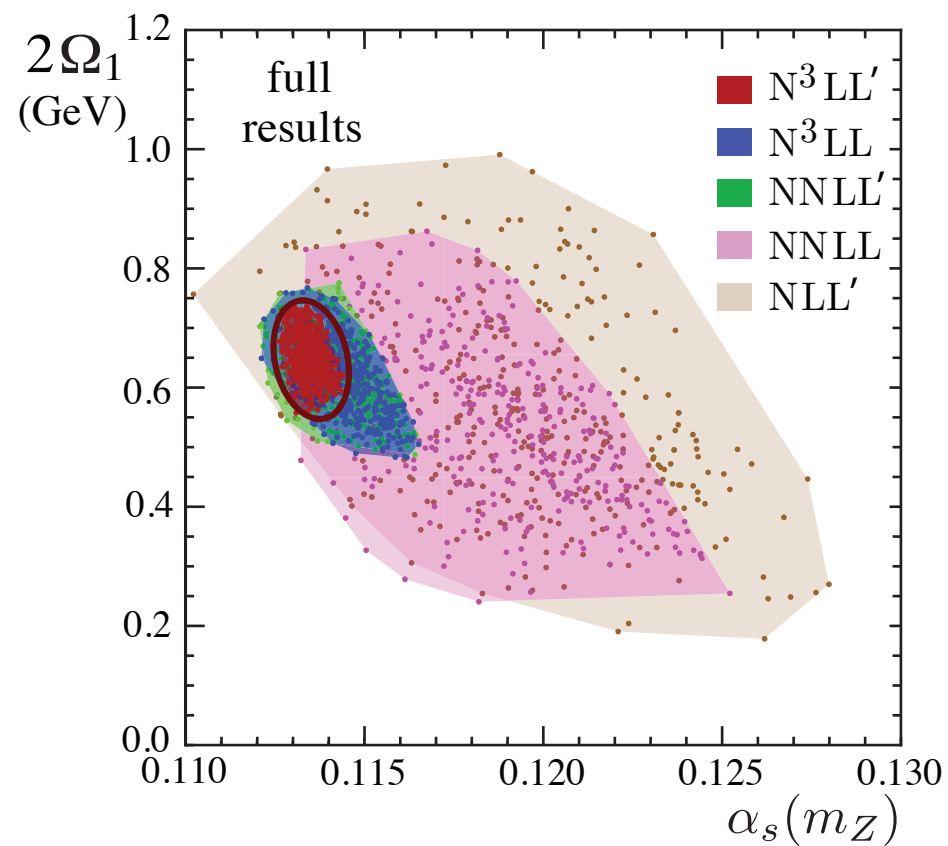
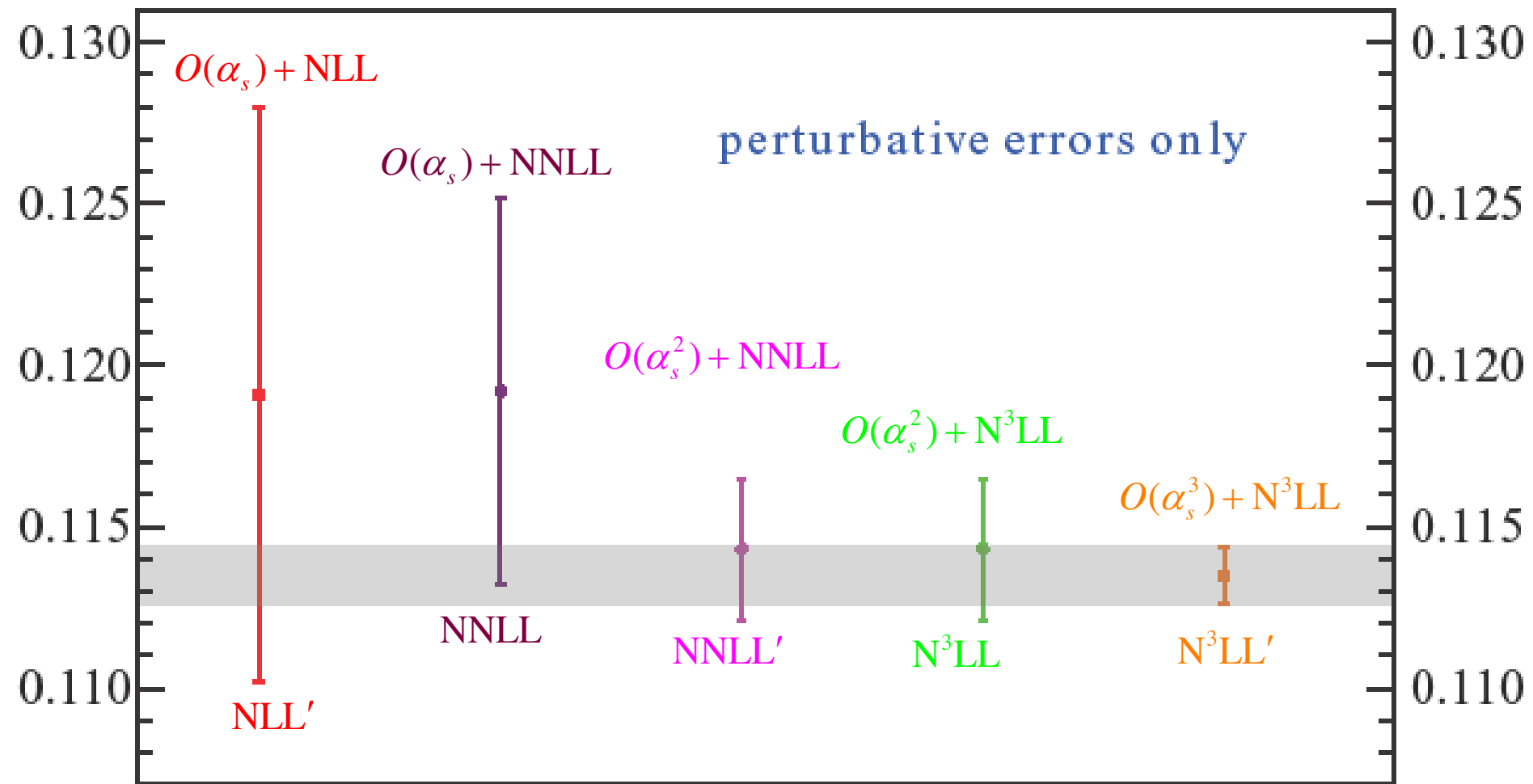
We fit for Ω_1 & $\alpha_s(m_Z)$ simultaneously. Strong degeneracy lifted by many Q's.



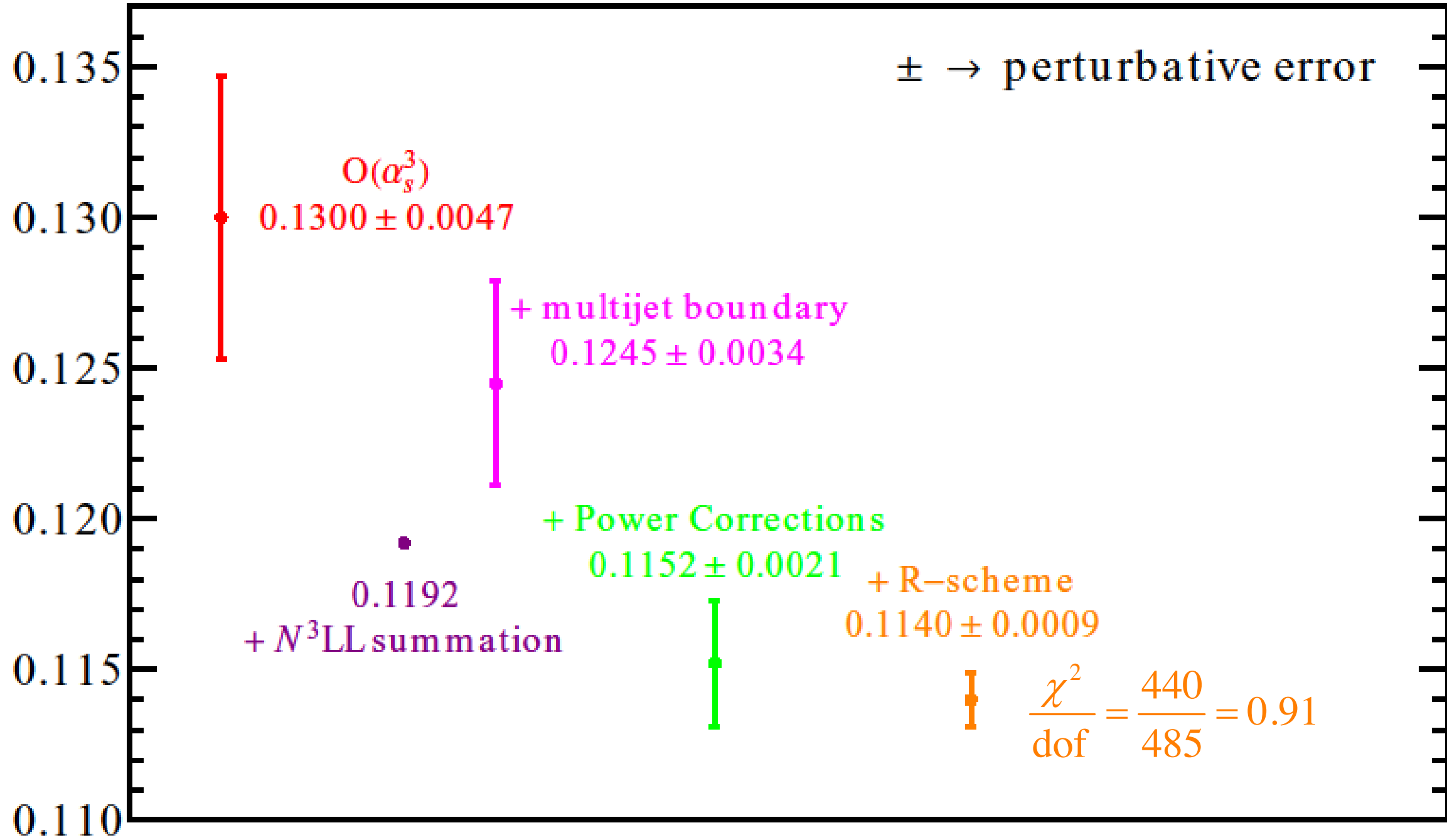
Power correction needed with 20% accuracy to get α_s at the 1% level





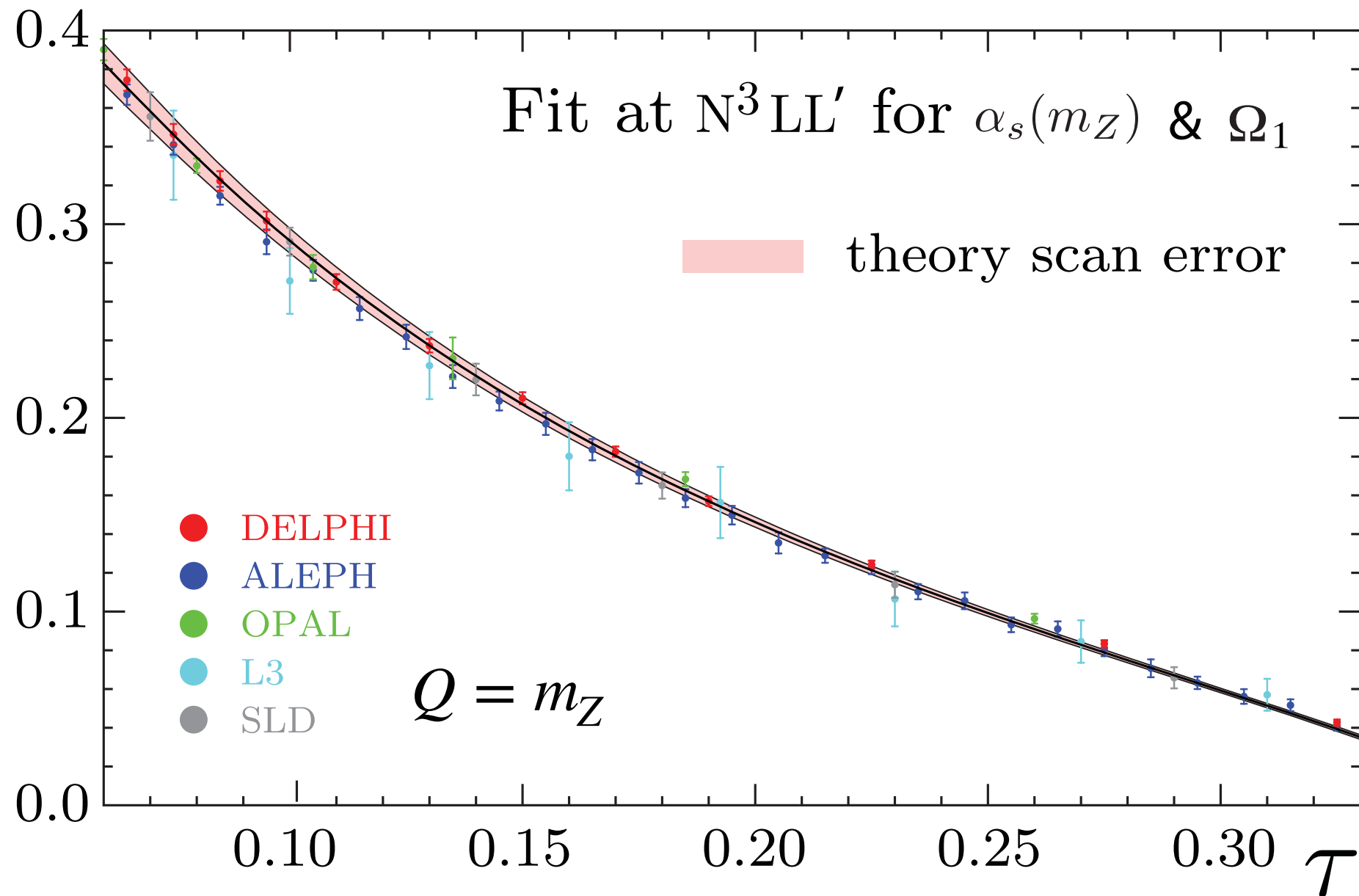


$\alpha_s(m_Z)$ from global thrust fits



Power correction has expected impact: $\frac{\delta\alpha_s}{\alpha_s} \approx -(9 \pm 3) \%$

$$\frac{\tau}{\sigma} \frac{d\sigma}{d\tau}$$



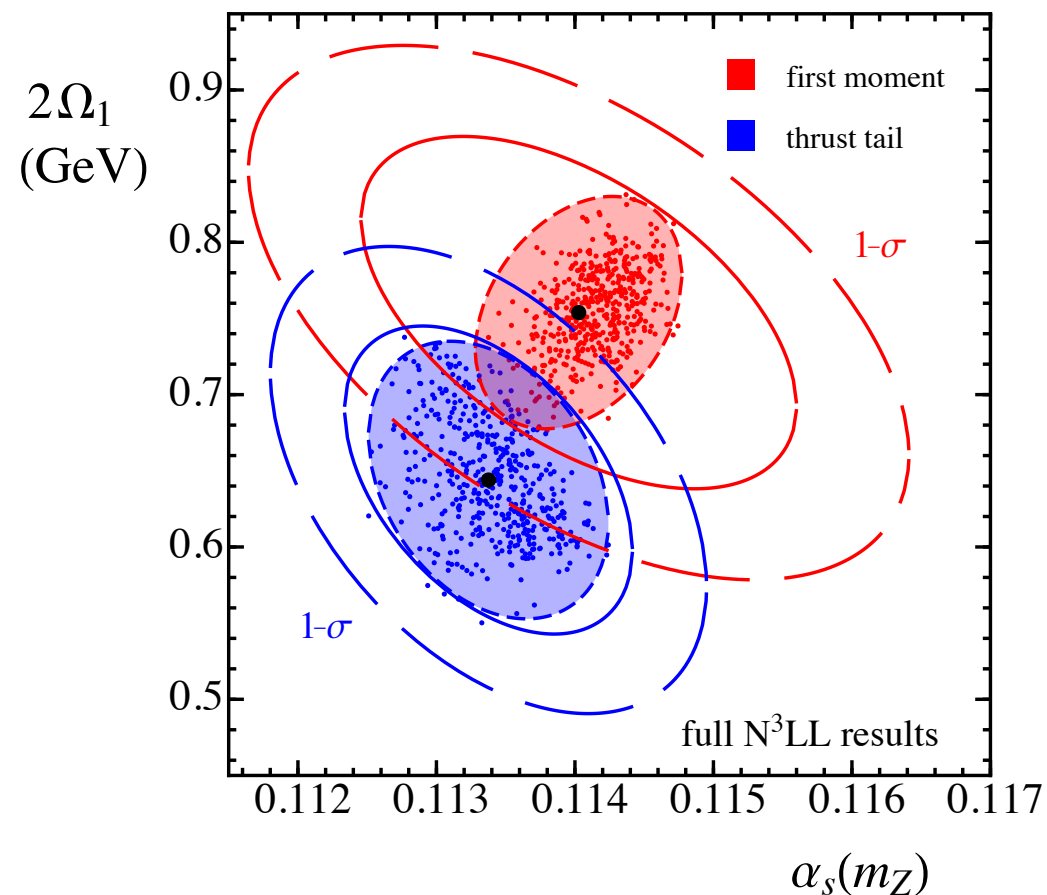
+ agreement for many other Q 's

Consistency checks

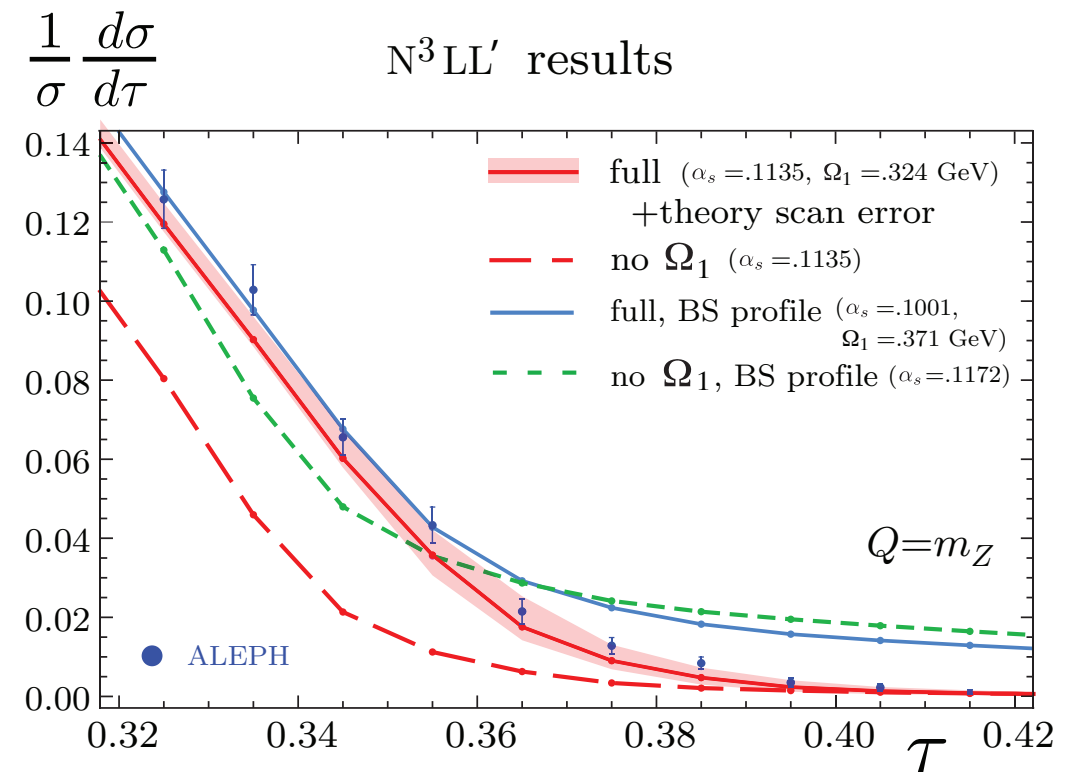
QED & b-mass effects small

$$\Delta\alpha_s(m_Z) = -0.0005$$

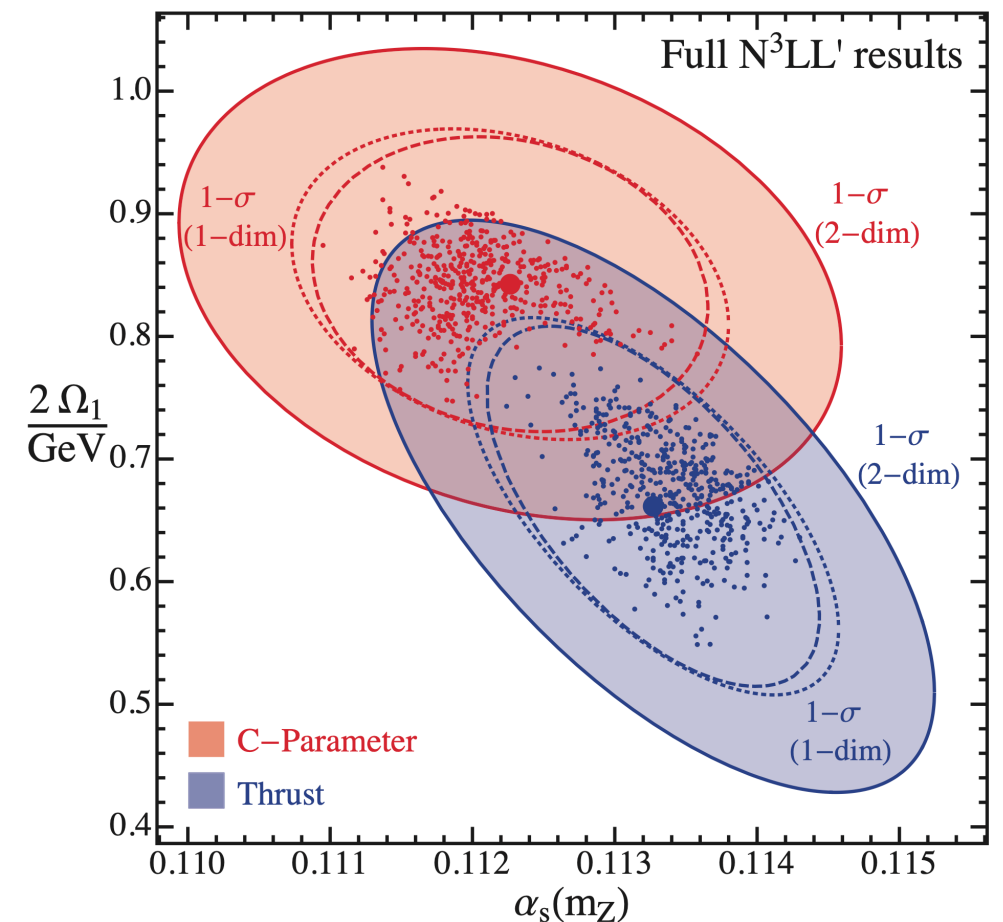
Thrust vs. thrust moments



Agreement beyond the fit region



Thrust vs. C-parameter

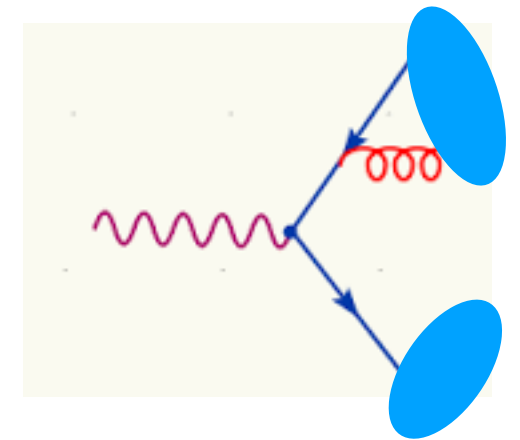
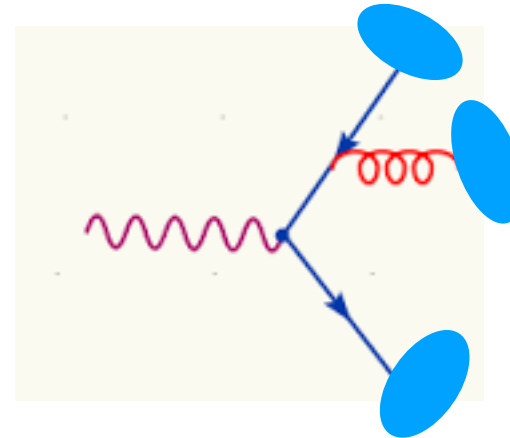


What could be going on?

Issue 1: Power corrections for 2-jets (Ω_1) versus 3-jets ($\neq \Omega_1$)

- perturbative dijet cross section dominates
- fit used dijet power correction for all τ

3-jet power corrections in fit region?



- Luisoni, Monni, Salam (2021)

Use dispersive model to “compute” power correction at $C = 3/4$

Find within that model that: $\Omega_1^{2\text{-jet}} \simeq 2\Omega_1^{3\text{-jet}}$ (model has one NP param.)

Note: In QCD 3-jet configuration is described by operators with 6 Wilson lines (generalizing the 4 lines for dijets), and hence ratio is non-perturbative too.

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- Caola et.al. (2022)

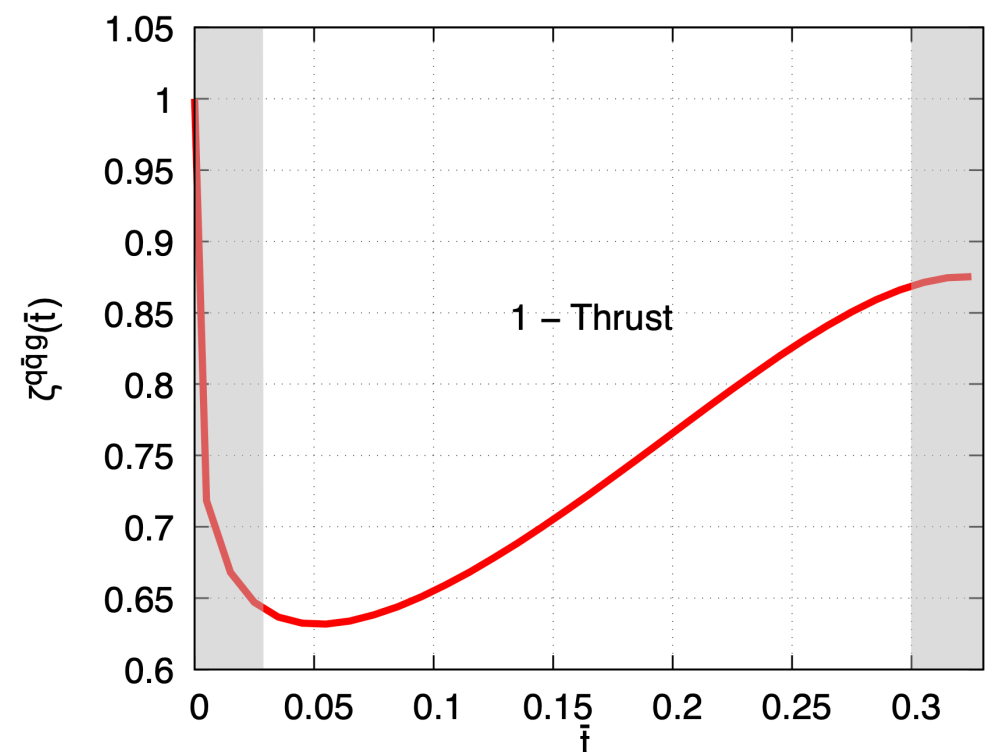
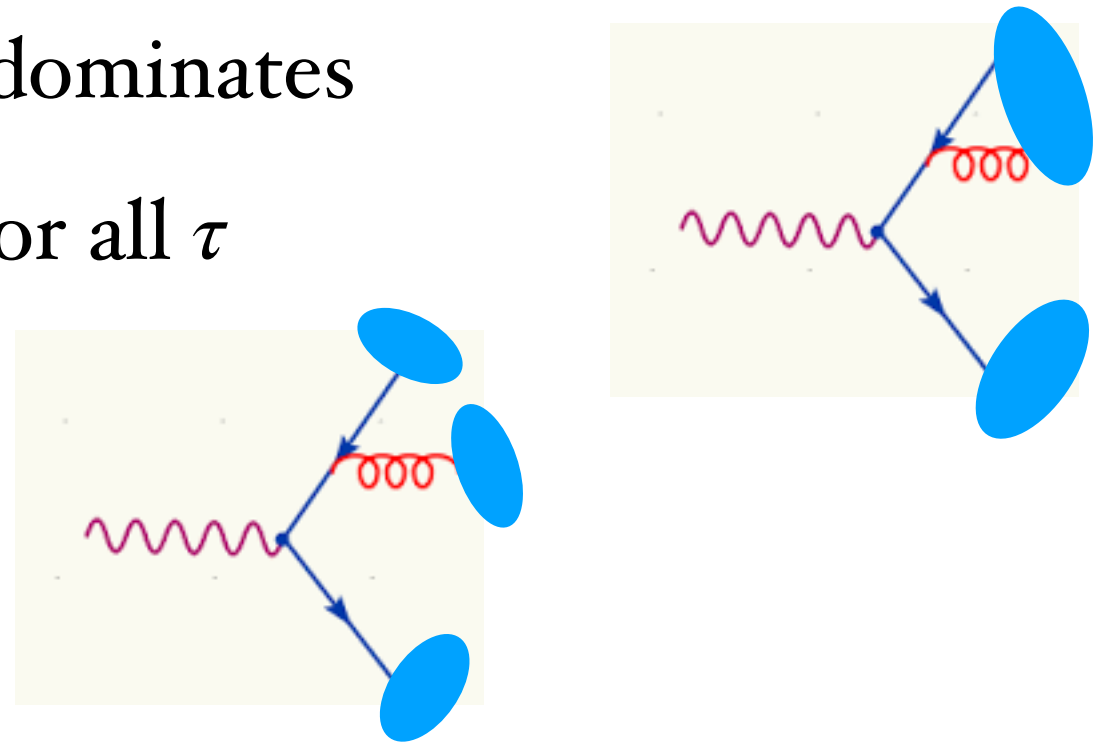
Large β_0 limit, probe power corrections with small gluon mass

$$\frac{d\sigma}{d\tau} = \frac{d\hat{\sigma}}{d\tau} \left(\tau - \frac{\zeta_\tau(\tau) 2\Omega_1}{Q} \right)$$

Rapid change at small τ .

Increases by 30% in fit region.

Note: method implicitly assumes 3-jets can always be resolved



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3-jet power corrections in fit region?

- Nason, Zanderighi (2023)

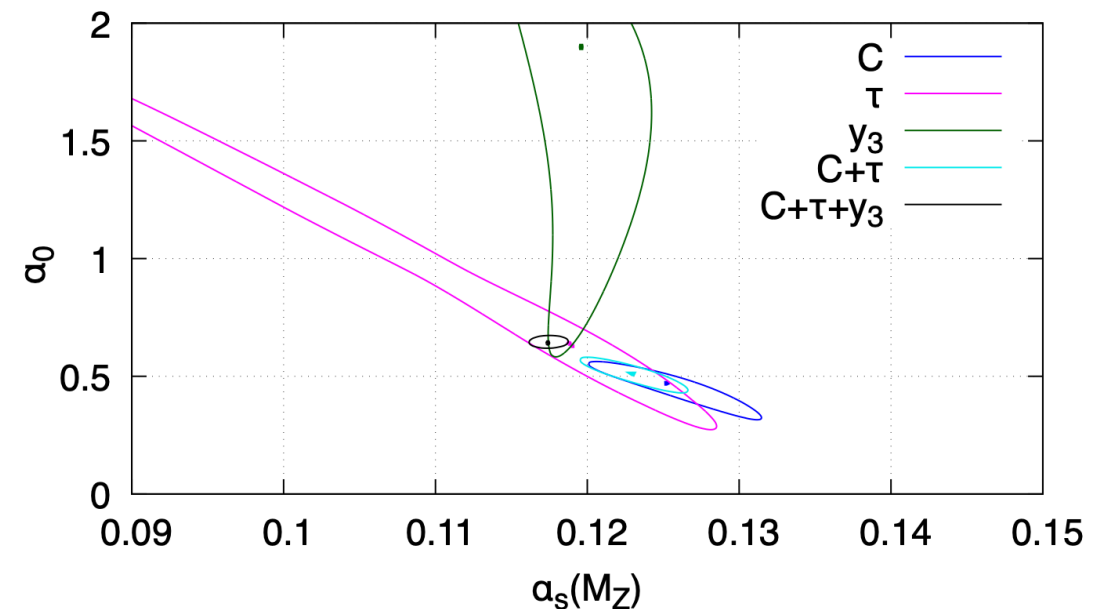
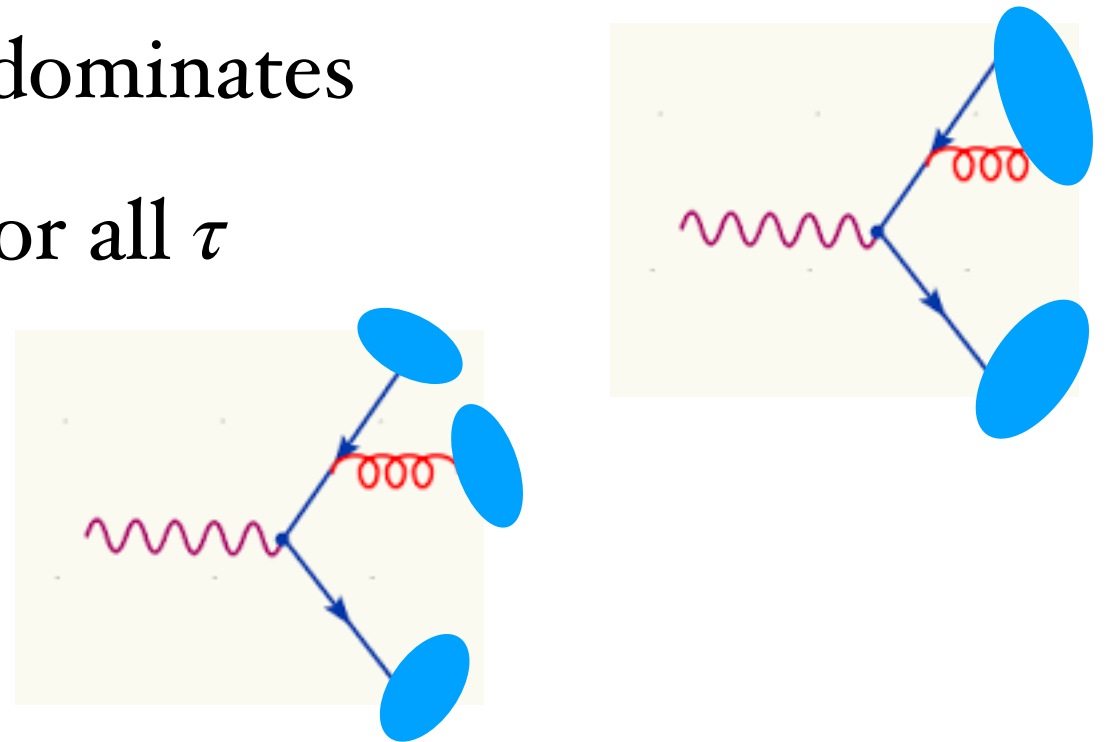
Analysis using ALEPH at $Q = m_Z$, various event shapes

Implements 3-jet $\zeta(\tau)$ model

Conclude that 3-jet power correction can lead to larger α_s

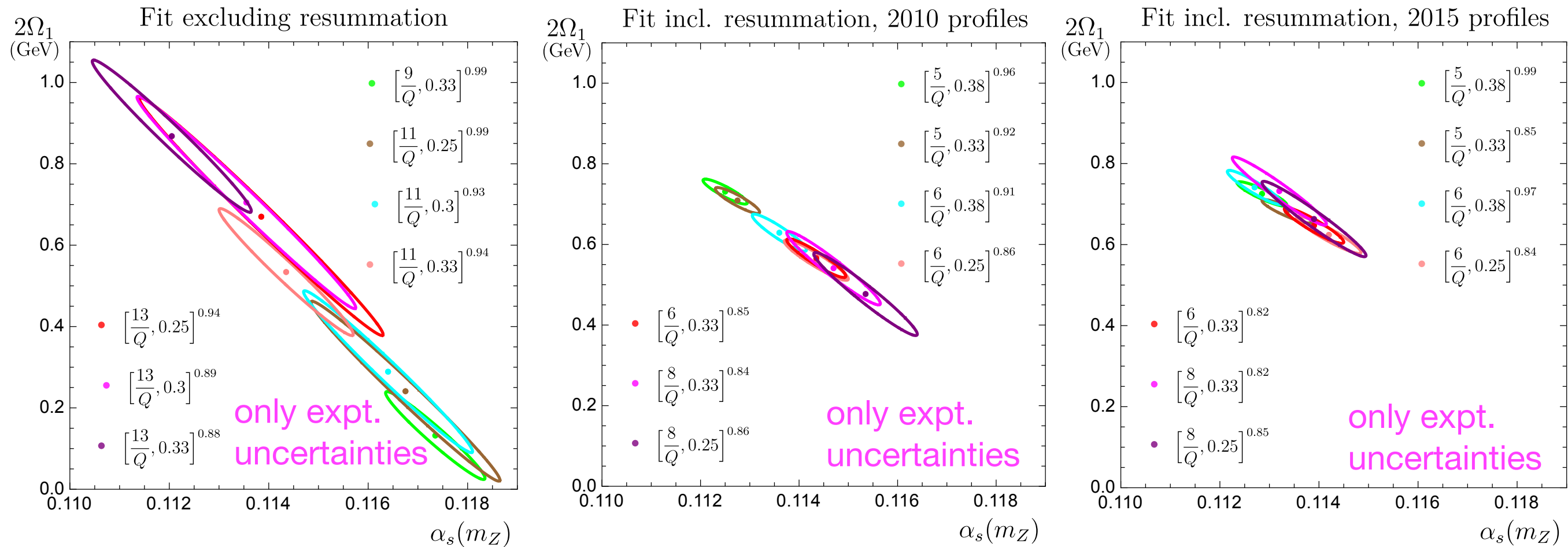
Note: Neglects resummation

Does not correlate theory across data bins



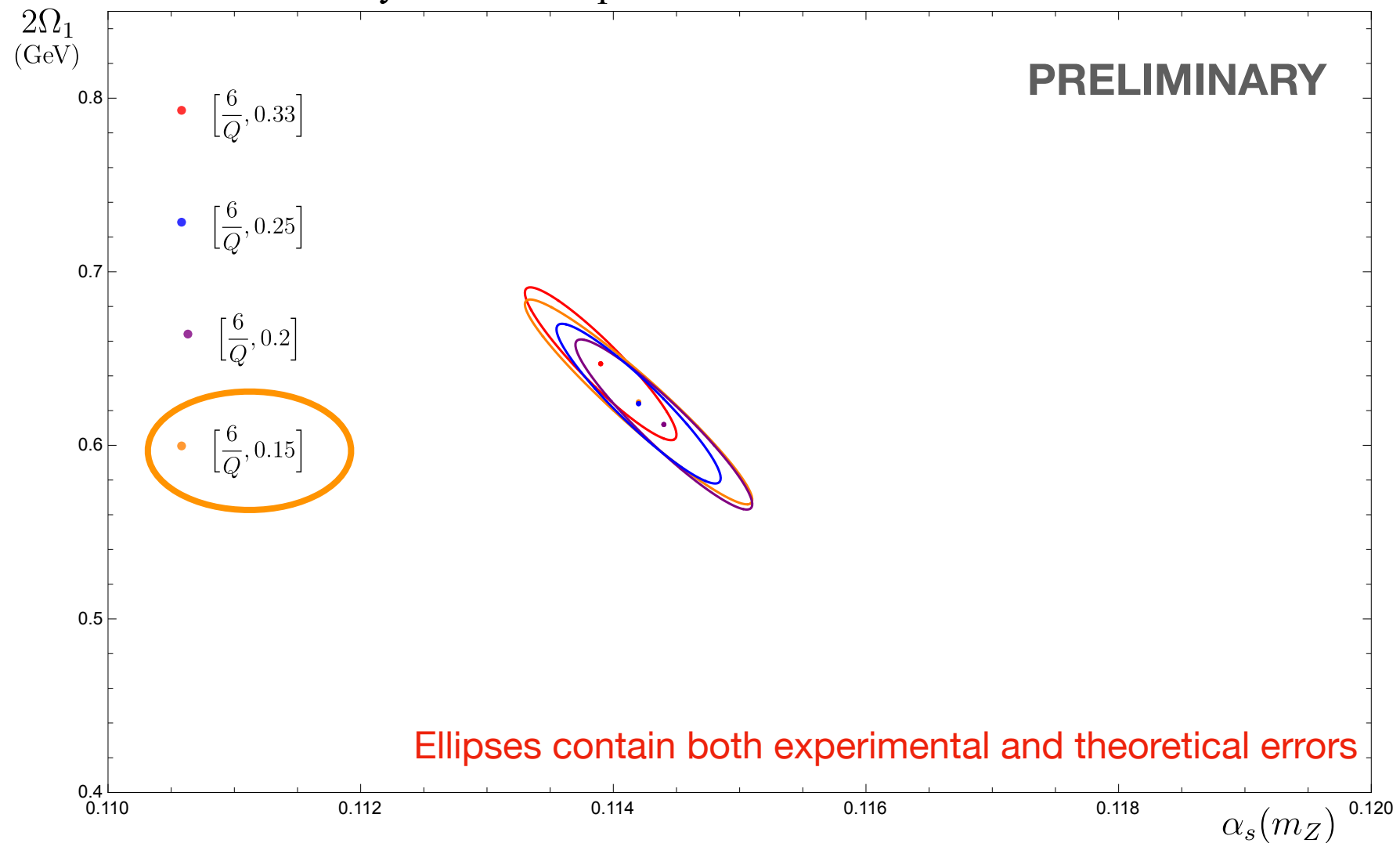
Issue 2: Fit range

Benitez-Rathgeb, Hoang, Mateu, IS, Vita
(work in progress)



$[\tau_{\min}, \tau_{\max}]^{\chi^2/dof}$ use of improved 2015 theory leads to stability

Opens up the possibility of a fit that focuses on region that is clearly dijet: $\tau \in [6/Q, .15]$



- Increased experimental uncertainty due to smaller fit region $[6/Q, .15]$:

$$\alpha_s = 0.1142 \pm 0.0006_{\text{pert}} \pm 0.0009_{\text{exp}+\Omega_1} \pm 0.0004_{3\text{jet}} = 0.1142 \pm 0.0012_{\text{tot}} \quad \langle \chi^2 \rangle / \text{dof} = 0.86$$

Compare to 2010 thrust fit ($[6/Q, .33]$, without QED corrections):

$$\alpha_s = 0.1140 \pm 0.0008_{\text{pert}} \pm 0.0005_{\text{exp}+\Omega_1} \pm 0.0001_{\text{had}} = 0.1140 \pm 0.0010_{\text{tot}}$$

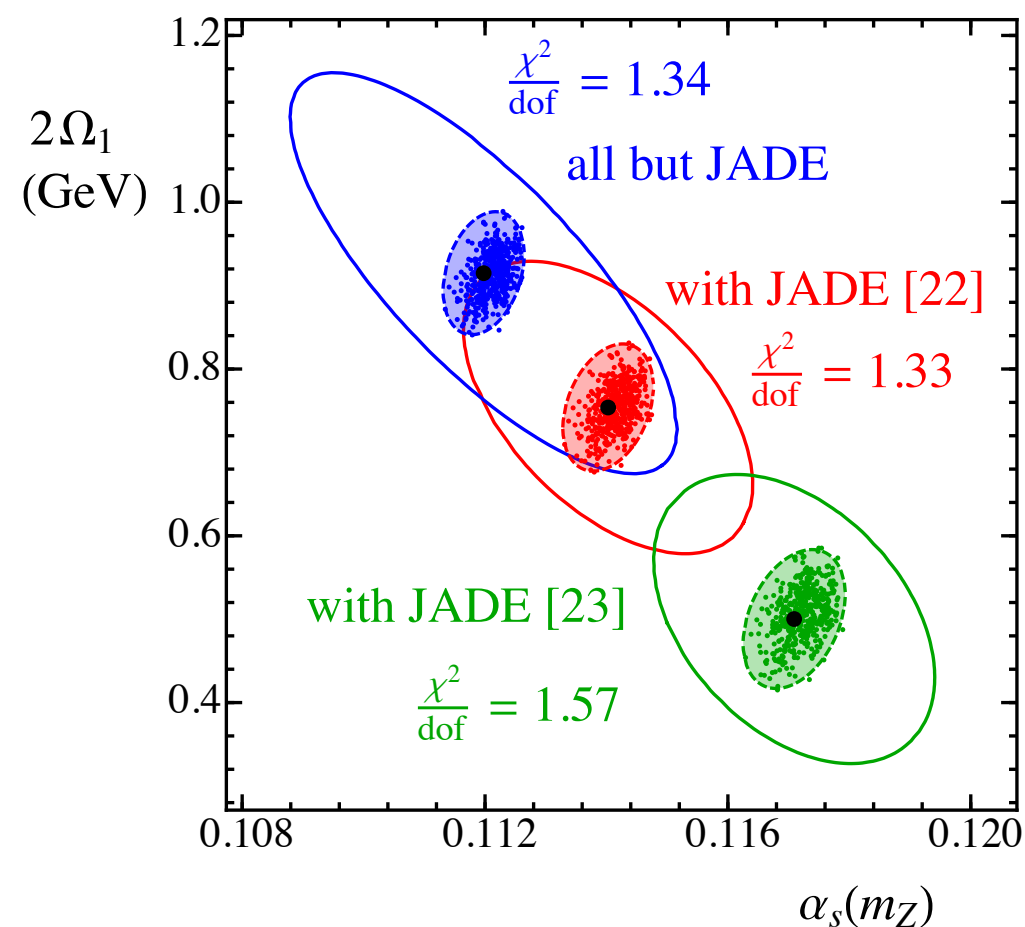
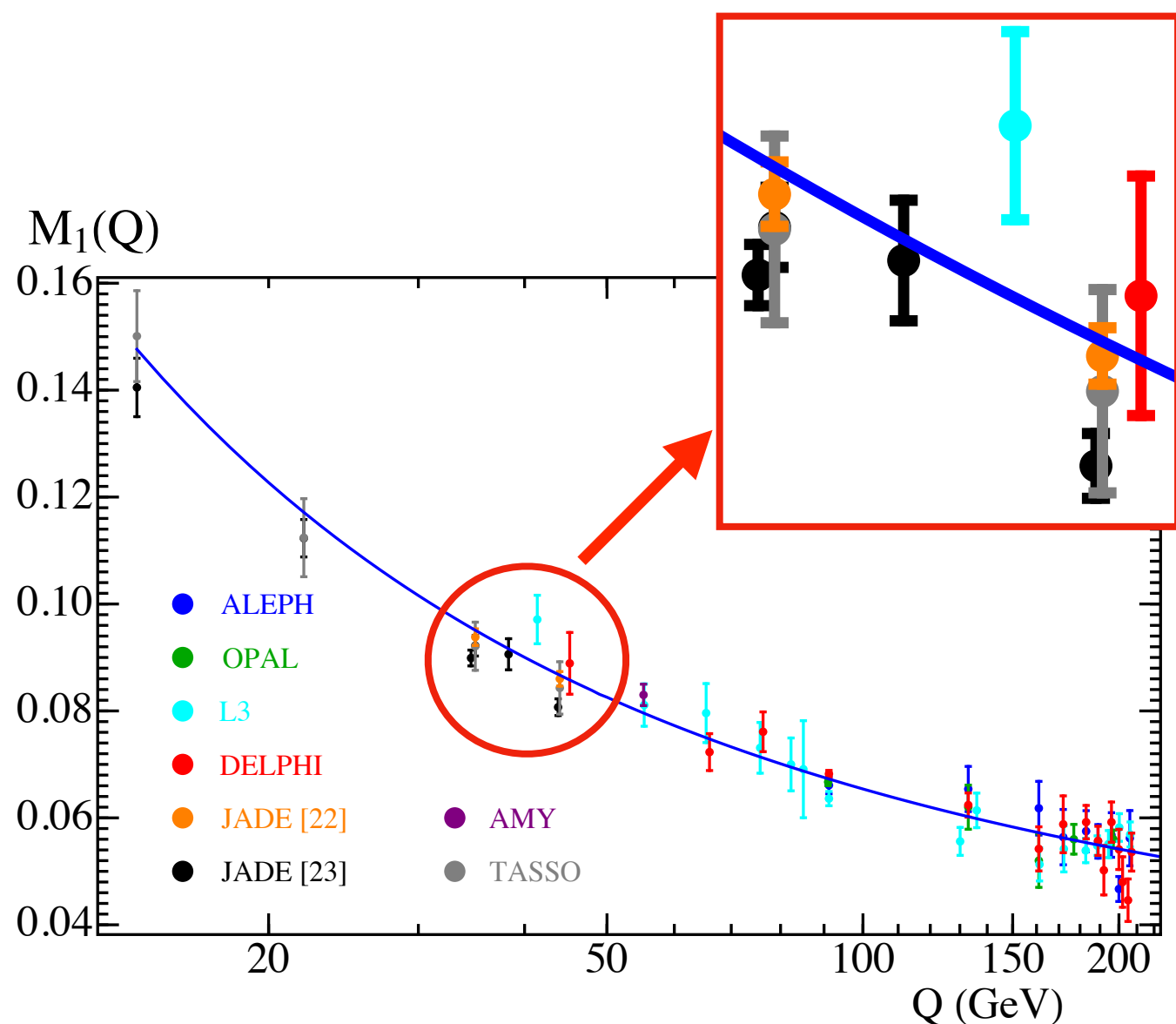
- Estimated size for 3-jet power corrections here $\simeq 0.0004$

Puzzle remains

Issue 3: Worth reconsidering systematics.

At the time of LEP, MC generators (used for calibration, determining acceptances, resolution ...) were much less sophisticated than those we have now.

An analysis that updated this was done by JADE for thrust moments in 2009 [23] (updating original 1998 [22])

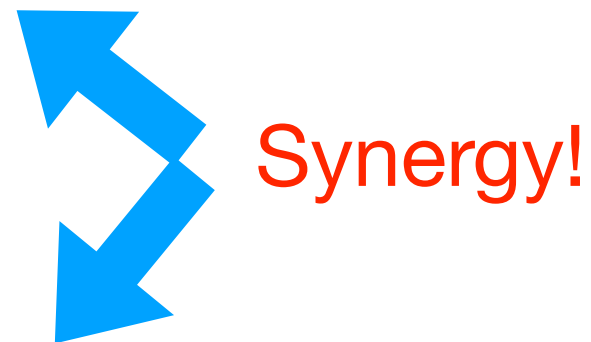


Summary & Outlook

$\alpha_s(m_Z)$ from e^+e^- jets, looking forward to the next e^+e^- collider

Theory wish list

- More rigorous treatment of 3-jet power corrections is crucial (including transition to 2-jet)
- Analyses with more observables (heavy-jet mass, EECs, ...), and combined observables while including all theory correlations
- Subleading power resummation, extension to N⁴LL, ...



Experimental wish list

- More kinematic info: multi-differential distributions, jet substructure, ...
- Full correlation matrices
- Impact of using modern Monte Carlo generators when comparing to LEP